

Explaining the DAMPE e^+e^- excess using the Higgs triplet model with a vector dark matter

Chuan-Hung Chen,^{1,*} Cheng-Wei Chiang,^{2,3,†} and Takaaki Nomura^{4,‡}

¹*Department of Physics, National Cheng-Kung University, Tainan, Taiwan 70101, Republic of China*

²*Department of Physics, National Taiwan University, Taipei, Taiwan 10617, Republic of China*

³*Institute of Physics, Academia Sinica, Taipei, Taiwan 11529, Republic of China*

⁴*School of Physics, KIAS, Seoul 130-722, Korea*



(Received 4 January 2018; published 19 March 2018)

We explain the e^+e^- excess observed by the DAMPE Collaboration using a dark matter model based upon the Higgs triplet model and an additional hidden $SU(2)_X$ gauge symmetry. Two of the $SU(2)_X$ gauge bosons are stable due to a residual discrete symmetry and serve as the dark matter candidate. We search the parameter space for regions that can explain the observed relic abundance, and compute the flux of e^+e^- coming from a nearby dark matter subhalo. With the inclusion of background cosmic rays, we show that the model can render a good fit to the entire energy spectrum covering the AMS-02, Fermi-LAT, CALET and DAMPE data.

DOI: [10.1103/PhysRevD.97.061302](https://doi.org/10.1103/PhysRevD.97.061302)

I. INTRODUCTION

In a recent report, the Dark Matter Particle Explorer (DAMPE) Collaboration showed the high-energy cosmic-ray electron-and-positron (CRE) spectrum from 25 GeV to 4.6 TeV with unprecedented precision and power to discriminate between electrons and protons [1]. The overall data can be well fit with a smoothly broken power-law model in the range of 55 GeV to 2.63 TeV. However, a peak in the bin of ~ 1.5 TeV stands out of the continuum. Though statistically insignificant yet, such a sharp excess of electrons and positrons in the cosmic rays still leads people to wonder if they may come from the annihilation of dark matter (DM) in a nearby subhalo. In particular, quite a few works propose DM models involving a new leptophilic interaction [2–12]. A model-independent analysis regarding which type of DM model to fit data and constraints better is given in Ref. [13].

In this work, we revisit a DM model [14] proposed a few years ago. (See also similar models in Refs. [15,16].) In the model, we extend the Higgs triplet model (HTM) [17–19] with a hidden gauge symmetry of $SU(2)_X$ that is broken to its Z_3 subgroup by a quadruplet scalar field. Such a symmetry-breaking scheme renders the new gauge

bosons stable and good candidates for weakly interacting massive particle (WIMP) DM. By coupling the complex Higgs triplet field and the $SU(2)_X$ quadruplet scalar field, the vector DM bosons can annihilate through one component of the $SU(2)_X$ quadruplet into a pair of doubly charged Higgs bosons, each of which in turn decays into like-sign leptons provided the Higgs triplet vacuum expectation value (VEV) is sufficiently small. Assuming the like-sign electrons and positrons as the dominant decay modes of the doubly charged Higgs boson, we show that the required excess at ~ 1.5 TeV in the DAMPE CRE spectrum can be produced. Besides, the model provides a link between neutrino mass and dark matter phenomenology.

II. THE MODEL

Since the model has been detailed in Ref. [14], here we only review the relevant parts for explaining the DAMPE CRE excess. In addition to the SM gauge group, the model has an additional $SU(2)_X$ symmetry with the associated gauge field and coupling strength denoted by X_μ^a and g_X , respectively. The $SU(2)_X$ symmetry is broken by a quadruplet field $\Phi_4 = (\phi_{3/2}, \phi_{1/2}, -\phi_{-1/2}, \phi_{-3/2})^T/\sqrt{2}$ that does not carry SM gauge charges, where the subscript stands for the eigenvalue of the third generator (denoted by τ_3) for the field and we use the phase convention that $\phi_{-i} = \phi_i^*$. As in the HTM, we have a complex Higgs field Δ that is a triplet under the SM $SU(2)_L$ and carries hypercharge $Y = 1$, where we adopt the convention that the electric charge $Q = T_3 + Y$.

By requiring that the Φ_4 field spontaneously develops a VEV in the $\tau_3 = \pm 3/2$ component [20],

*physchen@mail.ncku.edu.tw

†chengwei@phys.ntu.edu.tw

‡nomura@kias.re.kr

Published by the American Physical Society under the terms of the [Creative Commons Attribution 4.0 International license](https://creativecommons.org/licenses/by/4.0/). Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Funded by SCOAP³.

$$\phi_{\pm 3/2} = \frac{1}{\sqrt{2}}(v_4 + \phi_r \pm i\xi), \quad (1)$$

with $v_4 \sim \mathcal{O}(10)$ TeV, the $SU(2)_X$ symmetry is broken and the gauge bosons $\chi_\mu(\bar{\chi}_\mu) = (X_\mu^1 \mp iX_\mu^2)/\sqrt{2}$ and X_μ^3 acquire mass, $m_\chi = \sqrt{3}g_X v_4/2$ and $m_{X^3} = \sqrt{3}m_\chi$, all at the TeV scale. The model still has a residual Z_3 symmetry, under which χ_μ and $\bar{\chi}_\mu$ carry nonzero charges. Such a discrete gauge symmetry ensures the stability of χ_μ and $\bar{\chi}_\mu$ to be good DM candidates. Here, ϕ_r plays the role of a messenger between the hidden sector and the visible sector through the gauge interaction given by [20]

$$I_{\chi\bar{\chi}\phi_r} = \sqrt{3}g_X m_\chi \phi_r \chi_\mu \bar{\chi}^\mu. \quad (2)$$

After the electroweak symmetry breaking as in the SM, the Higgs triplet is induced to develop a VEV, serving as a source of the Majorana mass for neutrinos. We parametrize the Higgs doublet and triplet fields as

$$\Phi = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + \phi + iG^0) \end{pmatrix} \quad \text{and} \quad \Delta = \begin{pmatrix} \delta^+/\sqrt{2} & \delta^{++} \\ (v_\Delta + \delta^0 + i\eta^0)/\sqrt{2} & -\delta^+/\sqrt{2} \end{pmatrix}, \quad (3)$$

where the triplet VEV v_Δ is constrained by the electroweak rho parameter to be less than a few GeV [19,21,22]. To produce the CRE excess given by the DAMPE experiment, we assume the dominant decay modes of the charged Higgs boson to be leptonic. In this case, v_Δ is required to be less than $\sim 10^{-4}$ GeV.¹

Because of the hierarchy of the VEVs among $\langle\Phi_4\rangle$, $\langle\Phi\rangle$, and $\langle\Delta\rangle$, $v_4 \gg v \gg v_\Delta$, there is little mixing among ϕ_r , δ^0 , and ϕ . We will, therefore, identify them as the physical Higgs bosons H , δ^0 and h , respectively, with the masses:

$$\begin{aligned} m_h &\approx m_\phi = \sqrt{2}\lambda v, \\ m_H &\approx m_{\phi_r} = \sqrt{2\lambda_\Phi} v_4, \\ m_{\delta^{\pm\pm}} &\approx m_{\delta^\pm} \approx m_{\delta^0} = m_{\eta^0} \equiv m_\Delta. \end{aligned} \quad (4)$$

In the limit of no mixing, H does not couple with the SM particles directly. Instead, it interacts with the visible sector via the interactions described by [20]

$$I_{H\Delta\bar{\Delta}} = v_4 \lambda_4 H \left[\delta^{++} \delta^{--} + \delta^+ \delta^- + \frac{1}{2}(\delta^{02} + \eta^{02}) \right]. \quad (5)$$

¹In general, the doubly charged (singly charged) Higgs boson can also decay to $W^\pm W^\pm$ ($W^\pm h$, $W^\pm Z$, tb) mode(s) whose widths are proportional to v_Δ^2 . The branching ratios of these modes are suppressed when $v_\Delta \lesssim 10^{-4}$ GeV, and the only dominant decay modes are leptonic [23].

To explain the DAMPE e^+e^- excess, we take $m_\chi \approx 3$ TeV, $m_\Delta/m_\chi = 0.995$, and $m_H \approx 6$ TeV. Such a parameter choice is consistent with the latest collider search of the doubly charged Higgs boson that gives a lower bound of $m_\Delta > 770\text{--}800$ GeV [24].

III. DARK MATTER RELIC ABUNDANCE AND ANNIHILATION SIGNAL

Here, we consider the scenario where the DM particles annihilate through the $\chi\bar{\chi}H$ gauge interaction given in Eq. (2) and the $H\Delta\bar{\Delta}$ interaction given in Eq. (5) into a pair of Higgs triplet bosons. The assumed masses above result in the Breit-Wigner resonance enhancement in the pair annihilation process. This in turn affects both relic abundance of DM [25] and positron/antiproton fluxes [26,27]. Subsequently, the Higgs triplet bosons decay dominantly into leptonic final states: $\delta^{\pm\pm} \rightarrow \ell^\pm \ell'^{\pm}$ and $\delta^\pm \rightarrow \ell^\pm \nu_{\ell'}$ with details depending on the values of v_Δ and the lepton Yukawa couplings with Δ . For definiteness, we assume that $\mathcal{B}(\delta^{\pm\pm} \rightarrow e^\pm e^\pm) \approx \mathcal{B}(\delta^\pm \rightarrow e^\pm \nu_e) \approx 100\%$.²

In the nonrelativistic limit, the DM annihilation cross section is given by [20]

$$\sigma v \simeq \frac{1}{192\pi} \left(\frac{\lambda_4}{m_\chi} \right)^2 \left[\left(\frac{v^2}{4} + 2\epsilon \right)^2 + \frac{\Gamma_H^2}{4m_\chi^2} \right]^{-1} \sqrt{1 - \frac{m_\Delta^2}{m_\chi^2}}, \quad (6)$$

where $\epsilon \equiv 1 - \frac{m_H}{2m_\chi}$ and Γ_H denotes the total width of H . Γ_H is found to be much smaller than m_χ and, therefore, its effect can be neglected. The average speed of DM v in units of the speed of light is typically $\sim 10^{-3}$ at the current Universe and ~ 0.3 at the freeze-out. After fixing $m_\chi = 3$ TeV and $m_\Delta = 0.995m_\chi$, σv is seen to depend only on two parameters: λ_4 and m_H .

Our numerical analysis is done by utilizing the micrOMEGAS 4.3.5 package [28] implemented with the model to solve the Boltzmann equation for the observed cold DM relic abundance, $\Omega_C h^2 = 0.1199 \pm 0.0027$ [29], at the 2σ confidential level (CL). We show in the left plot of Fig. 1 the thermally averaged DM cross section at the current Universe. It has a typical value around $10^{-27}\text{--}10^{-25}$ cm³/s around $m_H = 6$ TeV. The right plot gives the relation between λ_4 and m_H . The coupling λ_4 is seen to be perturbative within the displayed mass range of m_H . The peculiar behavior at $m_H = 6$ TeV in both plots is owing to the resonance effect in the s-channel DM annihilation process.

²This assumption is not crucial in our analysis. All that is required is that the two modes are dominant in the decays of singly and doubly charged Higgs bosons. This occurs when neutrinos have an inverted mass hierarchy and nonzero Majorana phases in mixing [23].

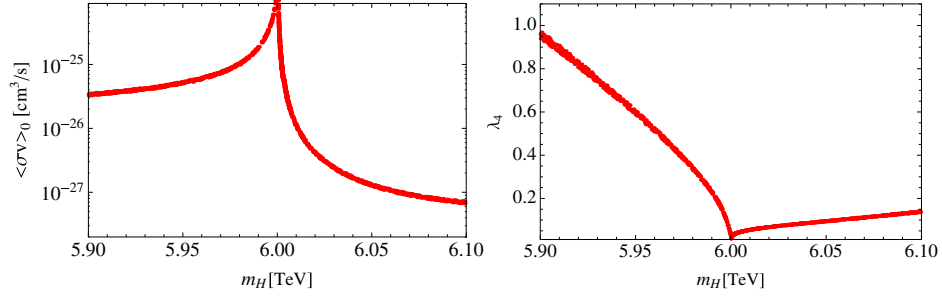


FIG. 1. Parameters required to fit the DM relic abundance in the vicinity of the Breit-Wigner enhancement region. Left plot: thermally averaged DM annihilation cross section as a function of m_H . Right plot: The quartic coupling λ_4 as a function of m_H . In these plots, we have fixed the DM mass $m_\chi = 3$ TeV.

IV. CRE SPECTRUM

As alluded to before, we assume that the CRE excess observed by the DAMPE experiment comes from the nearby subhalo, where the source term $q_{\text{DM}}^c(\vec{x}, E)$ in the diffusion-loss equation for a nearby DM clump at $x = x_c$ is given by [30]

$$q_{\text{DM}}^c(\vec{x}, E) = \frac{1}{4} \langle\sigma v\rangle_0 \frac{L}{m_\chi^2} \delta^3(\vec{x} - \vec{x}_c) \quad \text{with} \quad L = \int \rho^2 dV, \quad (7)$$

where $\langle\sigma v\rangle_0$ is the current thermally averaged cross section, L denotes the subhalo luminosity, and ρ is the DM density profile. In the following, we consider the expected CRE flux by taking into account propagation effects. The differential flux of e^\pm , defined as $d\Phi_{e^\pm}/dE = cF/4\pi$ in any point of our Galaxy, is given by [31]

$$\frac{d\Phi_{e^\pm}}{dE} = \frac{c}{4\pi b(E, \vec{x})} \left[L \int_E^{m_\chi} dE_s \sum_f \left(\kappa_f \frac{dN_f^\pm}{dE}(E_s) \right) \times \frac{I(E, E_s, \vec{x})}{(4\pi\lambda(E, E_s, \vec{x}))^{3/2}} \right], \quad (8)$$

where the speed of e^\pm is approximated by the speed of light, $v_{e^\pm} = c$, the effect through an annihilation channel f is expressed as $\kappa_f = \langle\sigma v\rangle_0^f/(4m_\chi)$, E_s denotes the e^\pm energy at the production, E is the observed e^\pm energy, $b(E, \vec{x})$ is the e^\pm energy loss coefficient function, and $I(E, E_s, \vec{x})$ is the generalized halo function. The detailed definition of parameter $\lambda(E, E_s, \vec{x})$ can be found in Ref. [32], where it is not only related to the diffusion coefficient function and $b(E, \vec{x})$, but also to the propagation distance of e^\pm . Representing the Green function of the diffusion-loss equation with the electrons and positrons produced by the DM annihilation as the source, the generalized halo function is independent of the DM model.

In addition to the peak at around $E \sim 1.5$ TeV, to obtain the e^+e^- spectrum from GeV to TeV, we also need the

fluxes of the cosmic-ray electron and positron backgrounds, which arise from various astrophysical sources. We follow the parametrizations in Ref. [33], in which the results were used to fit the AMS-02 data [34]. For positrons, we write the flux as

$$\begin{aligned} \Phi_{e^+}^{\text{sec,IS}}(E) &= C_{e^+} E^{-\gamma_{e^+}}, \\ \Phi_{e^+}^{\text{source,IS}}(E) &= C_s E^{-\gamma_s} \exp(-E/E_s), \end{aligned} \quad (9)$$

where $\Phi_{e^+}^{\text{sec,IS}}$ and $\Phi_{e^+}^{\text{source,IS}}$ denote the secondary and primary interstellar (IS) positrons, respectively, and the free parameters $C_{e^+,s}$, $\gamma_{e^+,s}$, and E_s are determined by the cosmic-ray data. Thus, the spectrum at the top of the atmosphere (TOA) is given by [35]

$$\begin{aligned} \Phi_{e^+}^{\text{TOA}}(E) &= \frac{E^2}{(E + \phi_{e^+})^2} [\Phi_{e^+}^{\text{sec,IS}}(E + \phi_{e^+}) \\ &\quad + \Phi_{e^+}^{\text{source,IS}}(E + \phi_{e^+})], \end{aligned} \quad (10)$$

where the typical value of the parameter ϕ_{e^+} varies between 0.5 and 1.3 GV. The electron flux, on the other hand, is parametrized as

$$\begin{aligned} \Phi_e^{\text{TOA}}(E) &= \frac{E^2}{(E + \phi_e)^2} \left[C_1 (E + \phi_e)^{-\gamma_1} \right. \\ &\quad \left. + C_2 (E + \phi_e)^{-\gamma_2} \left(1 + \frac{E}{E_s} \right)^{-\gamma_3} \right]. \end{aligned} \quad (11)$$

We note that the parametrization of the electron flux in Ref. [33] is only suitable for the range of $E < 500$ GeV and the spectrum is not suppressed when $E > 1$ TeV. To solve

TABLE I. Values of parameters for the interstellar positron and electron fluxes, where the units of various parameters are $[C_{e^+,s,1,2}] = \text{s}^{-1} \text{sr}^{-1} \text{m}^{-2} \text{GeV}^{-1}$, $[\phi_{e^\pm}] = \text{GV}$, and $[E_s] = \text{GeV}$.

C_{e^+}	γ_{e^+}	C_s	γ_s	ϕ_{e^+}	E_s
50	4.0	2.4	2.757	1.3	1100
C_1	γ_1	C_2	γ_2	ϕ_{e^-}	γ_3
3340	3.9	18	2.575	1.35	1.98

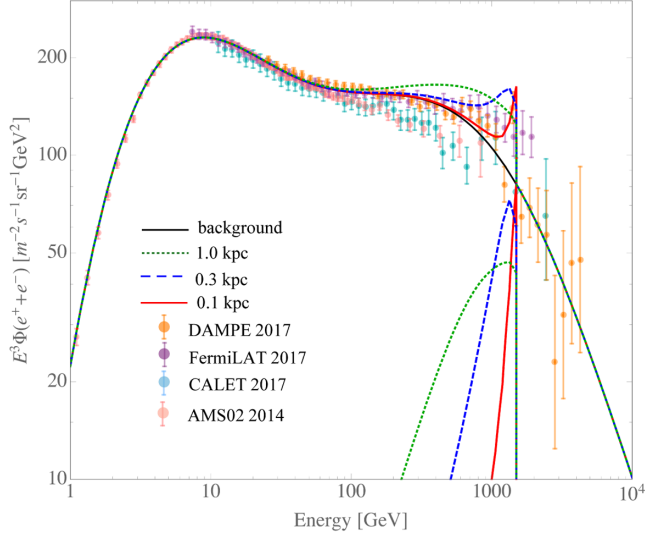


FIG. 2. The CRE spectrum (scaled by E^3) in the range of $1-10^4$ GeV. The AMS-02 data are drawn in pink, the Fermi-LAT data in purple, the CALET data in cyan, and the DAMPE data in orange. The black curve is the background contribution. The excess from DM annihilation in our model is shown for three scenarios: the red curve for a subhalo at a distance of 0.1 kpc; the blue dashed curve for 0.3 kpc; and the green dotted curve for 1 kpc.

the problem, we have slightly modified the parametrization so that the DAMPE data at $E > 2$ TeV can be accommodated. Values of the parameters used for estimating the background cosmic-ray electron and positron fluxes in this work are given in Table I.

We show in Fig. 2 our result of the E^3 -scaled CRE spectrum in the energy range of 1 to 10 TeV. The background contribution is given by the black curve, which is seen to fit well the AMS-02 data (pink) [36] for $E \lesssim 60$ GeV, the Fermi-LAT data (purple) [37] for $10 \text{ GeV} \lesssim E \lesssim 500$ GeV, the CALET data (cyan) [38] for $10 \text{ GeV} \lesssim E \lesssim 3$ TeV, and the DAMPE data (orange) for $E \gtrsim 120$ GeV except around the 1.5-TeV peak.³ The red curve shows the result after including the DM contribution from a subhalo at a distance of 0.1 kpc from the Earth. The blue dashed (green dotted) curve shows how the e^+e^- energy spectrum gets smeared out if the subhalo is located 0.3 kpc (1 kpc) away.

³It is noted that for the energy range of $70 \text{ GeV} \lesssim E \lesssim 500$ GeV, the DAMPE data are consistently higher than the AMS-02 data. Such a discrepancy may be partially attributed to the absolute energy scale calibration in both experiments [1].

V. SUMMARY

In view of the peak structure in the cosmic-ray electron-and-positron spectrum around 1.5 TeV as reported recently by the DAMPE Collaboration, we revisit a dark matter model as an extension of the Higgs triplet model. The dark matter candidate with a mass of about 3 TeV is the gauge boson associated with a hidden $SU(2)_X$ symmetry that is broken to its Z_3 subgroup by a quadruplet scalar field. The stability of the dark matter is ensured by the discrete gauge symmetry. The coupling between the $SU(2)_X$ quadruplet and the Higgs triplet facilitates the pair annihilation of dark matter particles into the charged Higgs bosons with mass slightly less than the dark matter candidate. The annihilation cross section enjoys a Breit-Wigner enhancement when we take the mediator mass to be about 6 TeV. We show the parameter space that can explain the observed dark matter relic abundance at 2σ level.

For a sufficiently small triplet vacuum expectation value induced by that of the standard model Higgs doublet, the charged Higgs bosons preferentially decay into lepton pairs. Using 100% branching ratios to the $e^\pm e^\pm$ and $e^\pm \nu$ modes, respectively, for the doubly charged and singly charged Higgs bosons, we show that it is possible to explain the 1.5-TeV peak as a result of the charged Higgs boson decays. Moreover, we consider that the putative signal comes from a nearby dark matter subhalo. The e^+e^- flux spectrum is evaluated with both background and propagation effects taken into account. Our result agrees well with the AMS-02 data in the lower energy regime and the DAMPE in the higher energy regime. The subhalo is preferred to locate at a distance of ~ 0.1 kpc away from us.

Finally, we note that because of the assumed little mixing between the 125-GeV Higgs boson and 6-TeV dark matter annihilation mediator, the scattering cross section between dark matter and nucleons is negligibly small. Hence, the model can readily evade the constraints from direct searches.

ACKNOWLEDGMENTS

This research was supported in part by the Ministry of Science and Technology of Taiwan under Grants No. MOST-106-2112-M-006-010-MY2 and No. MOST-104-2628-M-002-014-MY4.

Note added.—After finishing this work, a similar analysis also appeared [39]. In that work, the authors also employ leptonic decays of the Higgs triplet fields to explain the DAMPE e^+e^- excess but has a scalar dark matter candidate.

- [1] G. Ambrosi *et al.* (DAMPE Collaboration), *Nature (London)* **552**, 63 (2017).
- [2] Q. Yuan *et al.*, [arXiv:1711.10989](#).
- [3] Y. Z. Fan, W. C. Huang, M. Spinrath, Y. L. S. Tsai, and Q. Yuan, [arXiv:1711.10995](#).
- [4] P. H. Gu and X. G. He, *Phys. Lett. B* **778**, 292 (2018).
- [5] G. H. Duan, L. Feng, F. Wang, L. Wu, J. M. Yang, and R. Zheng, *J. High Energy Phys.* **02** (2018) 107.
- [6] L. Zu, C. Zhang, L. Feng, Q. Yuan, and Y. Z. Fan, [arXiv:1711.11052](#).
- [7] Y. L. Tang, L. Wu, M. Zhang, and R. Zheng, [arXiv:1711.11058](#).
- [8] W. Chao and Q. Yuan, [arXiv:1711.11182](#).
- [9] P. H. Gu, [arXiv:1711.11333](#).
- [10] J. Cao, L. Feng, X. Guo, L. Shang, F. Wang, and P. Wu, [arXiv:1711.11452](#).
- [11] G. H. Duan, X. G. He, L. Wu, and J. M. Yang, [arXiv:1711.11563](#).
- [12] X. Liu and Z. Liu, [arXiv:1711.11579](#).
- [13] P. Athron, C. Balazs, A. Fowlie, and Y. Zhang, *J. High Energy Phys.* **02** (2018) 121.
- [14] C. H. Chen, C. W. Chiang, and T. Nomura, *Phys. Lett. B* **747**, 495 (2015).
- [15] P. S. B. Dev, D. K. Ghosh, N. Okada, and I. Saha, *Phys. Rev. D* **89**, 095001 (2014).
- [16] C. H. Chen and T. Nomura, *J. High Energy Phys.* **09** (2014) 120.
- [17] W. Konetschny and W. Kummer, *Phys. Lett.* **70B**, 433 (1977).
- [18] J. Schechter and J. W. F. Valle, *Phys. Rev. D* **22**, 2227 (1980).
- [19] T. P. Cheng and L. F. Li, *Phys. Rev. D* **22**, 2860 (1980).
- [20] C. H. Chen and T. Nomura, *Phys. Lett. B* **746**, 351 (2015).
- [21] J. F. Gunion, R. Vega, and J. Wudka, *Phys. Rev. D* **43**, 2322 (1991).
- [22] C. Patrignani *et al.* (Particle Data Group), *Chin. Phys. C* **40**, 100001 (2016).
- [23] P. Fileviez Perez, T. Han, G. y. Huang, T. Li, and K. Wang, *Phys. Rev. D* **78**, 015018 (2008).
- [24] The ATLAS Collaboration (ATLAS Collaboration), Report No. ATLAS-CONF-2017-053.
- [25] P. Gondolo and G. Gelmini, *Nucl. Phys.* **B360**, 145 (1991).
- [26] D. Feldman, Z. Liu, and P. Nath, *Phys. Rev. D* **79**, 063509 (2009).
- [27] M. Ibe, H. Murayama, and T. T. Yanagida, *Phys. Rev. D* **79**, 095009 (2009).
- [28] G. Bélanger, F. Boudjema, A. Pukhov, and A. Semenov, *Comput. Phys. Commun.* **192**, 322 (2015).
- [29] P. A. R. Ade *et al.* (Planck Collaboration), *Astron. Astrophys.* **571**, A16 (2014).
- [30] P. Brun, T. Delahaye, J. Diemand, S. Profumo, and P. Salati, *Phys. Rev. D* **80**, 035023 (2009).
- [31] M. Cirelli, G. Corcella, A. Hektor, G. Hütsi, M. Kadastik, P. Panci, M. Raidal, F. Sala, and A. Strumia, *J. Cosmol. Astropart. Phys.* **03** (2011) 051; **10** (2012) E01.
- [32] M. Kuhlen and D. Malyshev, *Phys. Rev. D* **79**, 123517 (2009).
- [33] A. Ibarra, A. S. Lamperstorfer, and J. Silk, *Phys. Rev. D* **89**, 063539 (2014).
- [34] M. Aguilar *et al.* (AMS Collaboration), *Phys. Rev. Lett.* **110**, 141102 (2013).
- [35] J. S. Perko, *Astron. Astrophys.* **184**, 119 (1987).
- [36] M. Aguilar *et al.* (AMS Collaboration), *Phys. Rev. Lett.* **113**, 221102 (2014).
- [37] S. Abdollahi *et al.* (Fermi-LAT Collaboration), *Phys. Rev. D* **95**, 082007 (2017).
- [38] O. Adriani *et al.* (CALET Collaboration), *Phys. Rev. Lett.* **119**, 181101 (2017).
- [39] T. Li, N. Okada, and Q. Shafi, *Phys. Lett. B* **779**, 130 (2018).