



CHORUS

This is the accepted manuscript made available via CHORUS. The article has been published as:

Blue-green emission from epitaxial yet cation-disordered $\text{ZnGeN}_{2-x}\text{O}_x$

Celeste L. Melamed, M. Brooks Tellekamp, John S. Mangum, John D. Perkins, Patricia Dippo, Eric S. Toberer, and Adele C. Tamboli

Phys. Rev. Materials **3**, 051602 — Published 17 May 2019

DOI: [10.1103/PhysRevMaterials.3.051602](https://doi.org/10.1103/PhysRevMaterials.3.051602)

Blue-Green Emission from Epitaxial Yet Cation-Disordered $\text{ZnGeN}_{2-x}\text{O}_x$

C. L. Melamed^{1,2,*}, M. B. Tellekamp², J. S. Mangum¹, J. D. Perkins², P. Dippo², E. S. Toberer^{1,2}, and A. C. Tamboli^{2,1}

¹ *Physics Department, Colorado School of Mines*

² *National Renewable Energy Laboratory*

(Dated: April 19, 2019)

ZnGeN_2 offers a low-cost alternative to InGaN with the potential for bandgap tuning to span the green gap using cation site ordering. The addition of oxygen on the anion site creates an additional degree of electronic tunability. Here, we investigate the structure and optoelectronic properties of an epitaxial $\text{ZnGeN}_{2-x}\text{O}_x$ thin film library grown by combinatorial co-sputtering on $c\text{-Al}_2\text{O}_3$. Samples exhibit X-ray diffraction patterns and X-ray pole figures characteristic of a wurtzite (cation-disordered) structure with the expected 6-fold in-plane symmetry. Transmission electron microscopy reveals a semi-coherent interface with periodic dislocations that relieve strain from the large lattice mismatch, and confirms the in-plane and out-of-plane crystallographic orientation. Room-temperature photoluminescence exhibits peaks between 2.4 and 2.8 eV which are consistent with a sharp absorption onset observed by UV-vis spectroscopy. These results demonstrate low-cost synthesis of optically active yet cation disordered $\text{ZnGeN}_{2-x}\text{O}_x$, indicating a path toward application as a blue-green emitter.

I. INTRODUCTION

The III-V family of semiconductors revolutionized optoelectronics, from high-efficiency GaAs solar cells to the GaN-based LEDs which won the 2014 Nobel Prize for their record-breaking brightness and energy efficiency.^{1,2} However, there are serious limitations to the III-Vs, including (1) the requirement for low defect density which necessitates expensive epitaxial growth, and (2) the miscibility gap of InGaN which makes spanning the visible light spectrum difficult.^{3,4} A prospective solution to both issues is found in the II-IV- V_2 family of semiconductors, which are analogs to the III-Vs produced by heterovalent cation mutation (e.g. GaN becomes ZnGeN_2 by replacing each pair of Ga atoms with elements with one fewer and one more electron, maintaining the total charge). In general, adding elements into a compound offers more tunability than binaries, and nitride and oxynitride phase spaces are relatively unexplored.^{5,6} Additionally, the II-IV- V_2 phases crystallize in either a cation ordered or disordered structure, which is predicted to tune the bandgap at constant lattice parameter, enabling control of optoelectronic properties.⁷

ZnGeN_2 , which is analogous and lattice-matched to GaN, was first synthesized in 1970 and is theorized to have a bandgap of 3.5 eV in its cation-ordered structure.⁸⁻¹⁰ Like the other II-IV- V_2 materials, which are predicted to have ~ 1.0 eV of bandgap tunability, the bandgap of ZnGeN_2 may be tunable down to ~ 2.5 eV by introducing cation disorder. This range spans most of the visible spectrum, which could enable lattice-matched but bandgap-tunable optoelectronics that integrate with GaN. Control of cation ordering with synthesis temperature has been demonstrated for this system, but has not yet been correlated with optical properties.¹¹ While high-quality crystalline and epitaxial thin films have been synthesized by sputtering, molecular beam epitaxy (MBE) and metal-organic chemical vapor deposi-

tion (MOCVD), a thorough investigation of the impact of ordering on optoelectronic properties is lacking.¹²⁻¹⁶ Photoluminescence has been demonstrated both at room temperature¹⁴ and colder^{15,17-19} for cation ordered material grown by MOVPE and other vapor growth methods. Only once has photoluminescence been demonstrated for cation disordered material, in this case grown by MOCVD.^{16,20} Though bandgap tunability with ordering is an appealing goal, the optical properties of high-quality disordered material have not yet been fully explored.

Meanwhile, the past 15 years have seen a surge in interest in the $\text{ZnGeN}_2\text{-ZnO}$ alloy system for an entirely different application: photocatalytic processes.²¹⁻²⁴ The majority of this experimental work has been performed by either synthesizing precursor powder of Zn_2GeO_4 or mixing Zn and Ge oxide precursors, and exposing this to ammonia while heating, creating a compound of the form $(\text{Zn}_{1+x}\text{Ge})(\text{N}_2\text{O}_x)$.^{21,23-25} The structure of this alloy is reported to be wurtzite-like, with no observed peak splitting which would indicate orthorhombic ordering. Indeed, many of the reported studies of this system find that the reaction terminates in an orthorhombic (cation-ordered) ZnGeN_2 structure, but the orthorhombic superstructure peaks do not appear in XRD until the (nominally oxygen-free) end of the reaction.^{21,24} While the bandgap of ZnO is 3.37 eV, the alloy system with ZnGeN_2 exhibits bandgap bowing that lowers the bandgap to 2.0-3.0 eV.^{22,23} This is similar to the solid solution between ZnO and GaN.²⁶ Since the lattice constants of ZnO and ZnGeN_2 are very similar (less than 2% mismatch), this system offers another opportunity for bandgap tunability without significantly altering the lattice constant. No reports exist of thin films of this alloy, and the impact of cation disorder on the properties of this system has not been explored.

Here we present an epitaxial $\text{ZnGeN}_{2-x}\text{O}_x$ sample library on $c\text{-Al}_2\text{O}_3$ grown by combinatorial sputtering

which exhibits room-temperature photoluminescence between 2.4 and 2.8 eV. This result demonstrates that a cation-disordered $\text{ZnGeN}_{2-x}\text{O}_x$ thin film that is optically active in the green range can be grown with an inexpensive and scalable technique, indicating a path toward bandgap-tunable optoelectronic devices.

II. EXPERIMENTAL METHODS

$\text{ZnGeN}_{2-x}\text{O}_x$ thin film sample libraries were deposited by radio frequency co-sputtering onto stationary c-plane Al_2O_3 substrates with a 1° off-cut in the [100] direction heated to a set point of 750°C . Metallic zinc and germanium targets angled at 45° to the substrate normal created a gradient in cation flux during synthesis, generating an array of composition conditions for each deposition. The sputtering chamber was evacuated to a base pressure of $7 \cdot 10^{-7}$ Torr before deposition, and maintained at a working pressure of 15 mTorr during deposition with a gas flow of 10 sccm Ar and 20 sccm N_2 . Oxygen was not intentionally introduced, but was incorporated from background in the sputtering chamber.

Films were first characterized using the suite of spatially resolved characterization tools available at the National Renewable Energy Laboratory, which has been used to great success for previous Zn-IV- N_2 work.²⁷⁻²⁹ Cation composition, here reported as $\% \text{Zn}/(\text{Zn}+\text{Ge})$, was characterized using X-ray fluorescence (XRF) with Rh L-series excitation in energy-dispersive mode using Fischer XDV-SDD software. X-ray diffraction (XRD) was performed using a Bruker D8 Discover equipped with an area detector. Transmission and reflection spectra were collected in the UV-visible (UV-Vis) spectral ranges using a custom thin film optical spectroscopy system equipped with deuterium and tungsten/halogen light sources and Si detector arrays. The spectra were then used to calculate absorption coefficient α , using the relationship $\alpha = -\ln[T/(1-R)]/d$, where T , R , and d are transmission, reflection, and film thickness, respectively.

Specific points on the sample library were selected for further investigation. For characterization of epitaxial films, X-ray diffraction pole figure measurements were performed using an RU200 Rigaku DMAX-A instrument with a rotating $\text{Cu K}\alpha$ anode X-ray source operating at 40 kV/250 mA. A Rigaku SmartLab diffractometer was used to collect reciprocal space maps.

To investigate anion composition, Rutherford backscattering spectroscopy (RBS) was performed in a 130° backscattering geometry with a 2 MeV He+ beam energy using a model 3S-MR10 RBS system from National Electrostatics Corporation. A 2 MeV beam energy was used to avoid the non-Rutherford resonant scattering that occurs at 3 MeV for nitrogen and oxygen. The 130° scattering angle was chosen to provide well separated step edges in the scattered spectrum for, in order of increasing energy, substrate oxygen, film nitrogen, film oxygen and substrate aluminum. To obtain a

spectrum with sufficient signal-to-noise ratio to observe the scattering from oxygen, RBS data was collected until the total integrated charge delivered to the sample was $400 \mu\text{C}$. The film composition was determined by fitting using the RUMP analysis software³⁰ with the major cation ratio fixed at $\text{Zn}/(\text{Ge}+\text{Zn}) = 0.51$ as determined by XRF since Zn and Ge are too close in mass to be distinguished by RBS using a 2 MeV beam energy.

TEM micrographs were acquired with an FEI Co. Talos F200X transmission electron microscope with scanning capabilities operating at an accelerating voltage of 200 keV. Specimens for TEM were prepared from sputtered films via in-situ focused ion beam lift-out methods³¹ using an FEI Co. Helios Nanolab 600i SEM/FIB DualBeam workstation. Specimens were ion milled at 2 keV and 77 pA to remove Ga ion beam damage and achieve a final thickness of approximately 75 nm. Structural characterization was conducted by acquiring selected area electron diffraction (SAED) patterns on an FEI Co. Ceta 16M pixel CMOS camera at a camera length of 660 mm. The Al_2O_3 substrate was used to calibrate the camera constant, allowing SAED reflections to be accurately measured and indexed. Chemical mapping was performed in the TEM using the Super-X energy-dispersive X-ray spectroscopy (EDS) system.

Photoluminescence was acquired using a Newport MS260i imaging spectrometer, an InstaSpec X array detector with a 300/500 grating and an exposure time of 5 seconds. The PL was measured at room temperature with a 325 nm laser line at 4.5 mW. The CCD UV edge was at 400 nm and the long pass filter used was at 365 nm. Tapping mode atomic force microscopy (AFM) was performed on an Asylum Research system using a tip with a 150 N/m force constant to characterize surface morphology.

III. RESULTS

In this study, the substrate temperature and gun powers were fixed, resulting in a sample library with a linear composition gradient. This library consisted of 9 unique cation compositions from $\text{Zn}/(\text{Zn}+\text{Ge}) = 0.47$ to 0.51. After all spatially resolved characterization was performed, the library was cleaved into smaller pieces for single-point measurements (such as XRD pole figures, RBS and TEM).

A. X-ray Diffraction

X-ray diffraction patterns are shown as a function of cation composition in Fig. 1a. At lower ($\sim 47\%$) Zn concentrations, three wurtzite peaks are present. No superstructure peaks or peak splitting are visible, indicating cation-disordered material has been synthesized.¹¹ The Al_2O_3 (001) substrate peak is also present, and no additional peaks indicating secondary phases are observed.

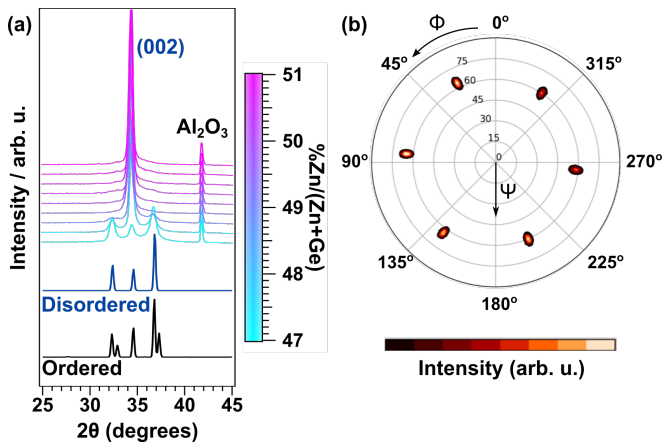


FIG. 1. (a) X-ray diffraction patterns as a function of cation composition (color axis). At lower Zn concentrations, three peaks indicating wurtzite $\text{ZnGeN}_{2-x}\text{O}_x$ are present, in addition to the Al_2O_3 peak. At stoichiometric and higher Zn concentrations, only the (002) wurtzite peak is present. (b) X-ray pole figures, which investigate in-plane orientation of a film, reveal 6-fold symmetry of the (101) peak indicating epitaxial alignment of the $\text{ZnGeN}_{2-x}\text{O}_x$ film with the Al_2O_3 substrate. The film is 30° rotated from the substrate (see Fig. S1 of the Supplemental Material³²), which is the same epitaxial relationship as GaN on *c*- Al_2O_3 .³³

As the Zn concentration increases, the wurtzite (002) peak increases in intensity and all other peaks are no longer observed, indicating a high degree of *c*-plane texturing in the film.

While the polycrystalline side of this library is easily identifiable as wurtzite from X-ray diffraction, the strongly textured side reveals a film peak position ranging from 34.3 to 34.5 degrees. This is slightly smaller than the theoretical values for the (002) peaks for the wurtzite and orthorhombic ZnGeN_2 structures, which are 34.56 degrees and 34.51 degrees respectively.^{8,9} Though slightly shifted, this peak position is consistent with reported values (002) peak of ZnGeN_2 , ZnO , and $\text{ZnGeN}_{2-x}\text{O}_x$, since these structures are all wurtzite and have a very similar lattice constant. In order to further investigate structure, X-ray pole figures were performed on the points where XRD indicated strong texturing. Peaks are present at the 2θ and Ψ positions that correspond to the wurtzite (101) peak, indicating that the film is indeed crystallized in the wurtzite structure. A representative scan of the (101) peak is shown in Fig. 1b. Six-fold symmetry indicates epitaxial alignment of the film with the Al_2O_3 substrate. The pole figures reveal a 30° in-plane rotation between the film and the substrate (Fig. S1), the same epitaxial relationship as between GaN and Al_2O_3 .³³

Reciprocal space maps (RSMs) reveal full-width at half-max (FWHM) values of 120 arcminutes for the (002) peak and 150 arcminutes for the (101) peak. In comparison, Daigo et. al. measured FWHM values of 100-140 arcminutes for a GaN film grown by sputter epitaxy on Si.³⁴ Additional RSMs were performed to determine

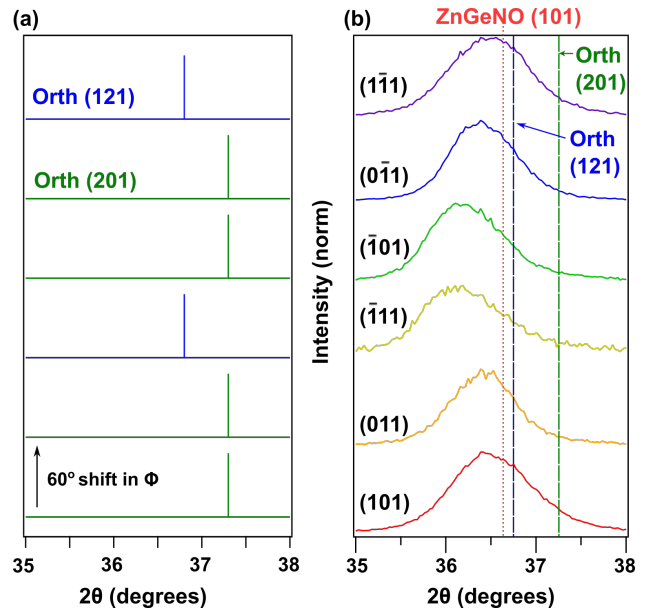


FIG. 2. (a) Simulated peak positions in 2θ for ordered (orthorhombic) ZnGeN_2 , each with a rotation of 60° in ϕ . The (201) family of reflections (blue) exhibit the same position in 2θ , while the two peaks in the (121) family (green) exhibit a distinct higher position in 2θ . (b) Line scans in 2θ from each RSM with each color at a rotation of $\sim 60^\circ$ in ϕ . The traces are indexed assuming a wurtzite structure. The slight shifts in 2θ are likely due to a small instrument misalignment, and no trend is seen indicating orthorhombic distortion when compared to (a). The red reference line is for the (101) peak of $\text{Zn}_{1.231}\text{Ge}_{0.689}\text{O}_{0.782}\text{N}_{1.218}$ reported in Ref. 25, while the blue and green reference lines are the orthorhombic ZnGeN_2 (121) and (201) positions shown in (a).

whether any cation ordering was occurring in the epitaxial sample. Orthorhombic distortion due to ordering would cause systematic peak shifting from the ideal wurtzite structure as a function of ϕ , as simulated in Fig. 2a. The (101) peak of the epitaxial film is observed to shift slightly to 2θ values lower than the ideal ZnGeN_2 structures, which is consistent with the incorporation of oxygen, as shown in Fig. 2b. The (101) peak displays small shifts with ϕ , but no systematic shifts that would indicate cation ordering (Fig. 2b). The observed peak shifts follow a sinusoid that is coincident with the measured alignment offset of the instrument. It is important to note that across the X-ray scan area, the superposition of cation-ordered orthorhombic grains with different in-plane orientations could yield a superposition of the orthorhombic (121) and (201) peaks shown in Fig. 2a. It is possible that the low signal and broad XRD peaks of this sputtered film are masking this type of spatially inhomogeneous ordering in Fig. 2b. However, the width of the asymmetric scan is the same order of magnitude as the symmetric scan rocking curve, which is not sensitive to ordering, suggesting that the broadening is due to sputtered material quality and not orthorhombic distortion. Additionally, considering that the ordered orthorhombic

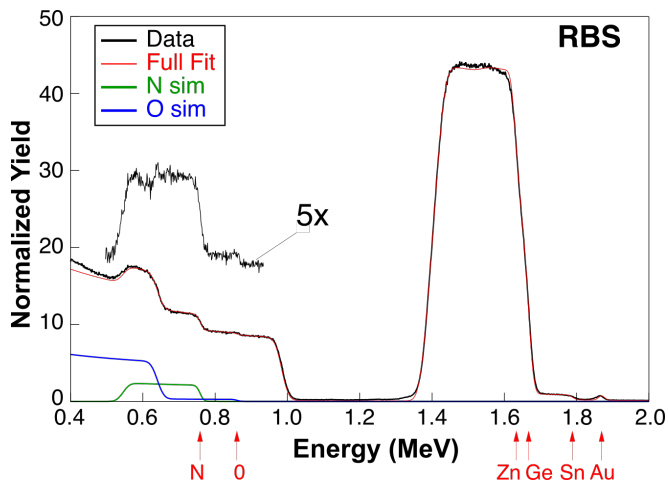


FIG. 3. Rutherford backscattering (RBS) analysis of one $\text{ZnGeN}_{2-x}\text{O}_x$ film on Al_2O_3 substrate with thin Au coating to avoid sample charging. Black: measured spectrum; Red: full fit to three layer model ($\text{Au/Zn-Ge-Sn-N-O/Al}_2\text{O}_3$); Green: N contribution to fit; Blue; O contribution to fit. Vertical red arrows indicate the front surface scattering energy for elements as labelled. The inset spectrum section shows the measured oxygen and nitrogen region with the modeled substrate signal subtracted expanded by 5X and offset for clarity.

structure of ZnGeN_2 has never been synthesized below 850°C ,¹¹ we hypothesize that this film is disordered.

B. Composition

Specific points on the sample library were selected for Rutherford backscattering (RBS) analysis to quantify oxygen and nitrogen content (Fig. 3). A thin gold coating was used to eliminate sample charging issues. Analysis was performed with different starting compositions in order to determine a standard error for the model. The Zn and Ge compositions were set by X-ray fluorescence to be 1.02 and 0.98 formula unit, respectively. The film was also found to contain 1 at.% of a heavy contaminant, here assumed to be Sn due to previous sputtering with this element in the synthesis chamber. This yields a compound composition of $\text{Zn}_{1.02}\text{Ge}_{0.98}\text{N}_{1.91}\text{O}_{0.18}$, with an anion-to-cation ratio of 1.05 and an oxygen composition of 8.8% O/(O+N). This is on the lower end of the reported oxygen content for $\text{ZnGeN}_{2-x}\text{O}_x$ material (discussed further in Section III E). Quantified EDS analysis shows uniform composition spatially throughout the investigated film, and confirms the composition numbers from RBS to an order of magnitude. EDS shows some carbon contamination as well, which mainly occurs at the substrate-film interface. It is promising that despite the Sn and C contaminants, the film is optically active and of relatively high crystalline quality.

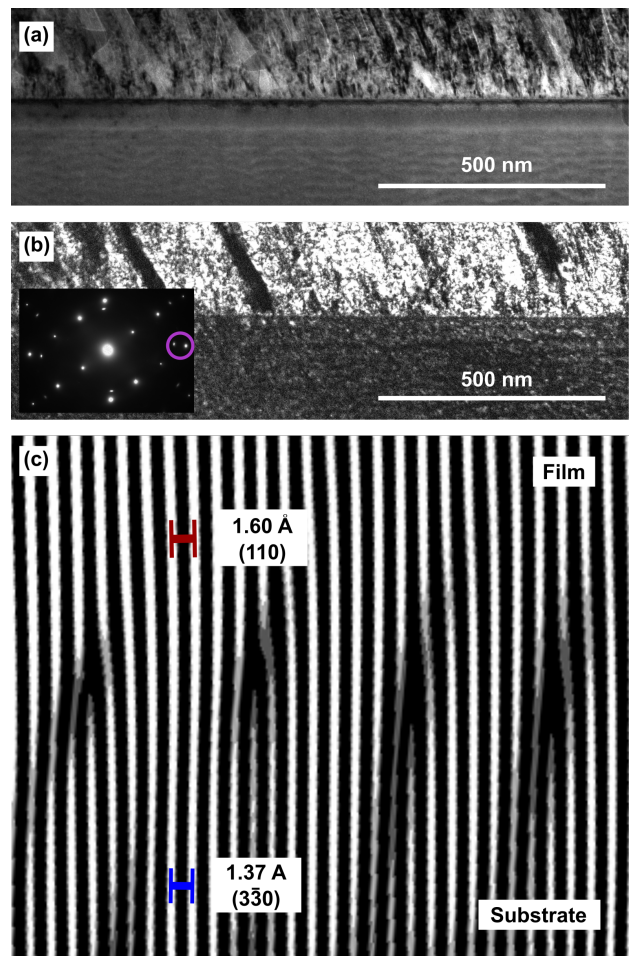


FIG. 4. (a) Bright-field TEM reveals a uniform, columnar film approximately 200 nm thick. (b) Dark-field TEM acquired using the film (110) and the substrate ($3\bar{3}0$) peaks (circled in selected area electron diffraction, inset). The film is non-uniform and shows a few dark regions indicating non-epitaxial grains. However, the majority of the film is aligned with the substrate and the same orientation. The substrate appears darker than the film due to the lower intensity of the substrate peak. (c) Fourier filtered high-resolution TEM micrograph showing a semi-coherent substrate-film interface with periodic dislocations approximately every 7-8 lattice planes to relieve strain due to large lattice mismatch between $\text{ZnGeN}_{2-x}\text{O}_x$ and Al_2O_3 . (The original lattice image is shown in Fig. S2 of the Supplemental Material³².)

C. Microscopy

Figure 4a-b shows TEM bright-field and dark-field micrographs of the epitaxial $\text{ZnGeN}_{2-x}\text{O}_x$ film on Al_2O_3 substrate. Selected area electron diffraction in the inset of (b) shows reflections corresponding to the $[110]$ zone axis of the Al_2O_3 substrate and the $[1\bar{1}0]$ zone axis of the $\text{ZnGeN}_{2-x}\text{O}_x$ film; the purple circle indicates the (110) film and ($3\bar{3}0$) substrate reflections used for dark-field imaging. As expected for a combinatorial sample grown by sputtering, the film is non-uniform and shows some

dark regions indicating non-epitaxial grains. However, the majority of the film is aligned with the substrate and a consistent crystallographic orientation. Fig. 4c shows a Fourier filtered high-resolution TEM micrograph of the substrate-film interface. Peaks in the Fourier transform of the as-acquired micrograph corresponding to the $\text{Al}_2\text{O}_3(3\bar{3}0)$ and $\text{ZnGeN}_{2-x}\text{O}_x(110)$ lattice planes were selectively filtered out to better view their crystallographic relationship. A semi-coherent interface is observed exhibiting periodic dislocations every 7-8 lattice planes which relieve strain due to lattice mismatch between the film and substrate.

The X-ray diffraction results from Section III A and the SAED results allow us to calculate the epitaxial alignment between substrate and film. Since only the (002) $\text{ZnGeN}_{2-x}\text{O}_x$ peak appears by X-ray diffraction, this suggests that the c-planes are aligned between the film and the substrate. This is confirmed by SAED, which demonstrates that $\text{ZnGeN}_{2-x}\text{O}_x(002)$ is parallel to $\text{Al}_2\text{O}_3(006)$. With regard to in-plane alignment, the X-ray pole figures reveal a 30° rotation between the film and the substrate. This is also confirmed by TEM: Fig. 4c demonstrates alignment between the $\text{ZnGeN}_{2-x}\text{O}_x(110)$ planes and the $\text{Al}_2\text{O}_3(3\bar{3}0)$ planes, which corresponds to a 30° rotation. This is the same relationship as reported previously for ZnGeN_2 on *c*- Al_2O_3 , and for GaN on *c*- Al_2O_3 .^{16,33} Based on these in-plane and out-of-plane relationships, the calculated in-plane lattice mismatch is 13.3%, which is reflected in the SAED and the high-resolution TEM (Fig. 4b-c).

Interestingly, the reciprocal space maps, pole figures and SAED all reveal a $\sim 3.6^\circ$ tilt between the film and the substrate. While Nagai tilt has been observed in GaN due to substrate off-cut, the tilt we observe is rotated 90° from the angle of the off-cut, inconsistent with the Nagai model.³⁵ The origin of this tilt is unknown and warrants further exploration.

Finally, AFM was performed at multiple points on the film to investigate surface morphology and growth microstructure (see Fig. S3 of the Supplemental Material³²). Scans revealed a dense, round microstructure with features ranging 75-250 nm diameter, with RMS roughness of 5-6 nm. Four-point electrical measurements were also performed, and the film was found to be insulating. This is expected considering the insulating properties of ZnGeN_2 predicted by DFT.³⁶

D. Optical Properties

In order to determine the feasibility of incorporation of $\text{ZnGeN}_{2-x}\text{O}_x$ into optoelectronic devices, optical properties of these films were studied using UV-visible spectroscopy and room-temperature photoluminescence. Figure 5a shows the absorption coefficient from UV-vis spectroscopy as a function of energy, with different colors representing different points on the sample library. The absorption onset is ~ 1.0 eV lower than the theoretical

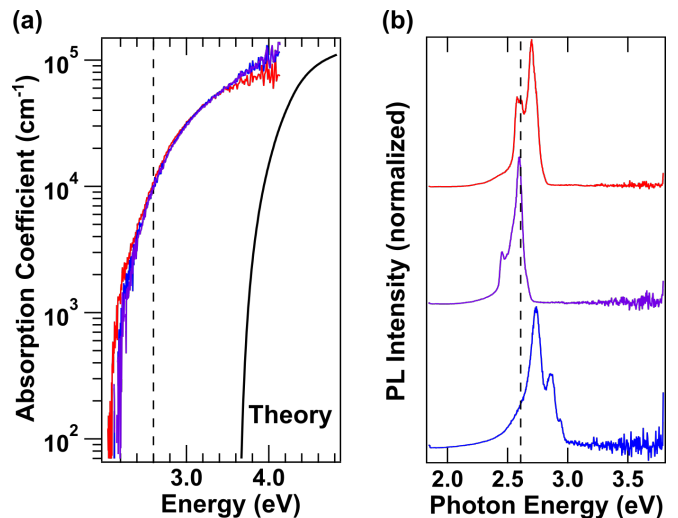


FIG. 5. (a) UV-Vis Spectroscopy reveals an absorption onset at approximately 2.6 eV for all samples. This is ~ 1.0 eV lower than the theoretical absorption onset for ordered (orthorhombic) ZnGeN_2 , shown as a black curve, consistent with bandgap tuning due to oxygen incorporation. The dashed line at 2.6 eV represents the energy at an absorption coefficient of 10^4 cm^{-1} . Colors are added to distinguish individual samples used in both UV-vis and PL. DFT calculations reproduced from Ref.⁷. (b) Room-temperature photoluminescence of $\text{ZnGeN}_{2-x}\text{O}_x$ films reveal peaks from 2.4 to 2.8 eV. Different points on the sample yield a slightly different set of peak positions and intensities. No trend is observed with cation composition variation. These values align with the observed 10^4 absorption coefficient value from the UV-vis spectroscopy (dashed line).

absorption onset for ordered ZnGeN_2 (shown as a solid black line), but remains a similar shape to theory, without band tailing that would indicate a high defect concentration. The energy value at an absorption coefficient of $10^4/\text{cm}$ is 2.6 eV, shown as a dashed line. Little change is observed with cation composition. Normalized room-temperature photoluminescence for different points on the sample library is shown in Fig. 5b. Though the PL signal was spatially varying throughout each film, bright spots were identified, revealing peaks from 2.4 to 2.8 eV. The peak variation does not correlate with changes in cation composition. These values are consistent with the UV-vis optical absorption onset of 2.6 eV (dashed line). Multiple varying-energy PL peaks is characteristic of a material exhibiting potential fluctuations near the band edge, which is expected in materials with structural disorder.^{28,37}

Due to the alignment of our PL results with the measured absorption data, we hypothesize that our PL emission is due to band-like recombination in the $\text{ZnGeN}_{2-x}\text{O}_x$ thin film. Though we could be observing “yellow band” defect states, to which other groups have attributed PL emission between 2 and 3 eV for cation-ordered ZnGeN_2 , the traditional yellow band is ~ 0.5 eV wide.¹⁷ The PL peaks reported here exhibit an average

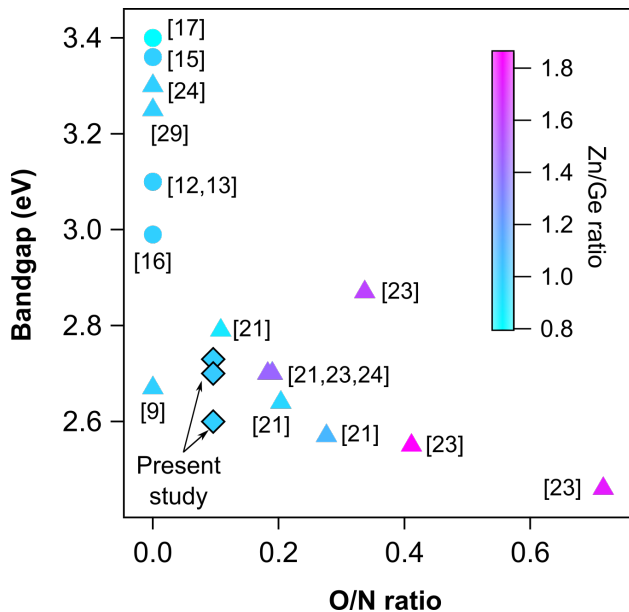


FIG. 6. Plot of reported bandgap versus O/N ratio from the literature, with triangular traces denoting a reaction process using oxide precursors. Circular traces denote a reaction process from metallic and/or nitride precursors. Diamonds represent the PL peak positions reported in this work (Fig. 5b). Zn/Ge ratio is plotted on the color axis.

FWHM value of 0.16 eV when each spectrum is fit with one Gaussian envelope. When each individual peak is fit, the average FWHM value is 0.08 eV. These widths are narrow compared to traditional defect band luminescence, and since the PL energies match the sharp UV-vis absorption onset and the high-energy tails follow a Maxwell-Boltzman distribution, this suggests that the optical signal originates from band-like luminescence.

E. Discussion

To place our results in historical context, data of reported bandgap versus O/N ratio was collected from the literature and is shown in Fig. 6. Triangular traces are plotted for material grown using a reaction process starting from oxygen-containing precursors, while circular traces denote a reaction from metallic and/or nitride precursors. Diamond traces represent values reported in this study. A clear decrease in bandgap is observed with increasing oxygen content that follows the energy bandgap bowing trend. The cation composition (color axis in Fig. 6) follows the expected trend as well, since an alloy between ZnGeN_2 and ZnO implies a trend of increasing Zn content with increasing O content. Many of the data points that report pure ZnGeN_2 material (corresponding to an O/N ratio of 0) do not report characterization of oxygen, suggesting that some unintentional oxygen content may be decreasing the bandgap and contributing to the large bandgap spread at O/N=0.

Though all reported $\text{ZnGeN}_{2-x}\text{O}_x$ material exhibits the disordered wurtzite structure, one could conceive of a cation-ordered $\text{ZnGeN}_{2-x}\text{O}_x$ structure of the same orthorhombic superstructure as the ground state of ZnGeN_2 . This raises a few questions about the $\text{ZnGeN}_{2-x}\text{O}_x$ system: Could cation ordering be used to tune the optical properties of $\text{ZnGeN}_{2-x}\text{O}_x$ at a fixed composition and lattice parameter? On the other hand, could a small amount of oxygen be used to trap ZnGeN_2 in a cation-disordered state? Control of cation ordering has been a sought-after research goal in the ZnGeN_2 system for optoelectronic researchers, but the presence of oxygen (intentional or not) may help or hinder this prospect. From a physical vapor deposition perspective, oxygen often incorporates unintentionally into thin films due to background contamination in vacuum systems, causing deleterious defects in material systems that require high purity. This work demonstrates that harnessing oxygen incorporation could provide an additional knob with which to tune properties of the $\text{ZnGeN}_{2-x}\text{O}_x$ system. Since this alloy exhibits only small changes in lattice parameter, this provides a tunable bandgap material from 2.4 eV to 3.4 eV that is lattice matched to GaN.

It is also important to note the diversity of characterization methods used to collect the bandgaps shown in this plot; the studies on ZnGeN_2 thin films typically report bandgaps from PL^{15,17,19} or Tauc plots from UV-Vis data,^{12,20} while most $\text{ZnGeN}_{2-x}\text{O}_x$ studies report Kubelka-Munk absorbance curves.^{21,23,24} We have plotted PL peak positions (2.60, 2.70, and 2.73 eV). None of these optical characterization methods are perfect, and it is evident that more data is necessary to deconvolve these results. Additionally, while Fig. 6 shows the impact of Zn/Ge and O/N on the bandgap, some of the variation in the data may be due to order parameter, which has not been taken into account. Further optical and structural characterization is key to deconvolve the impacts of cation composition, anion composition, and disorder on optical properties of $\text{ZnGeN}_{2-x}\text{O}_x$.

IV. CONCLUSION

In this work, we have reported an epitaxial $\text{ZnGeN}_{2-x}\text{O}_x$ sample library grown on c-plane Al_2O_3 by sputtering. X-ray diffraction and TEM confirm epitaxial alignment with the substrate despite 13.3% in-plane lattice mismatch. The sputter-deposited samples are optically active with room-temperature PL peaks that are narrow compared to traditional defect band luminescence. These peaks align with a sharp optical absorption onset at 2.6 eV, and we hypothesize this optical signal originates from band-related luminescence in $\text{ZnGeN}_{2-x}\text{O}_x$. To the authors' knowledge, this result is the first reported $\text{ZnGeN}_{2-x}\text{O}_x$ thin film, and sets the stage for low-cost light-emitting devices which integrate with GaN based on the ZnGeN_2 - ZnO material system.

V. ACKNOWLEDGMENTS

This work was authored in part by Alliance for Sustainable Energy, LLC, the manager and operator of the National Renewable Energy Laboratory for the U.S. Department of Energy (DOE) under Contract No. DE-AC36-08GO28308. This work was supported by the U.S. Department of Energy, Office of Science, Basic Energy Sciences, Materials Sciences and Engineering Division.

* cmelamed@mymail.mines.edu

- ¹ S. Essig, C. Allebé, T. Remo, J. F. Geisz, M. A. Steiner, K. Horowitz, L. Barraud, J. S. Ward, M. Schnabel, A. Descoedres, D. Young, M. Woodhouse, M. Despeisse, C. Balif, and A. Tamboli, *Nature Energy* **2**, 17144 (2017).
- ² NobelPrize.org, Nobel Media AB 2018 (2014).
- ³ S. Strite, *Journal of Vacuum Science & Technology B: Microelectronics and Nanometer Structures* **10**, 1237 (1992).
- ⁴ T. D. Moustakas, E. Iliopoulos, A. V. Sampath, H. M. Ng, D. Doppalapudi, M. Misra, D. Korakakis, and R. Singh, *Journal of Crystal Growth Proceeding of the Eleventh International Conference on Molecular Beam Epitaxy*, **227-228**, 13 (2001).
- ⁵ A. Zakutayev, *Journal of Materials Chemistry A* **4**, 6742 (2016).
- ⁶ W. Sun, C. Bartel, E. Arca, S. Bauers, B. Matthews, B. Orvañanos, B.-R. Chen, M. F. Toney, L. T. Schelhas, W. Tumas, J. Tate, A. Zakutayev, S. Lany, A. Holder, and G. Ceder, arXiv:1809.09202 [cond-mat] (2018), arXiv:1809.09202.
- ⁷ A. D. Martinez, A. N. Fioretti, E. S. Toberer, and A. C. Tamboli, *Journal of Materials Chemistry A* **5**, 11418 (2017).
- ⁸ M. Maunaye and J. Lang, *Materials Research Bulletin* **5**, 793 (1970).
- ⁹ W. L. Larson, H. P. Maruska, and D. A. Stevenson, *Journal of The Electrochemical Society* **121**, 1673 (1974).
- ¹⁰ A. Punya, W. R. L. Lambrecht, and M. van Schilfgaarde, *Physical Review B* **84** (2011), 10.1103/PhysRevB.84.165204.
- ¹¹ E. W. Blanton, K. He, J. Shan, and K. Kash, *Journal of Crystal Growth* **461**, 38 (2017).
- ¹² P. Narang, S. Chen, N. C. Coronel, S. Gul, J. Yano, L. W. Wang, N. S. Lewis, and H. A. Atwater, *Advanced Materials* **26** (2014), 10.1002/adma.201304473.
- ¹³ S. Kikkawa and H. Morisaka, *Solid State Communications* **112**, 513 (1999).
- ¹⁴ Q. H. Zhang, J. Wang, C. W. Yeh, W. C. Ke, R. S. Liu, J. K. Tang, M. B. Xie, H. B. Liang, and Q. Su, *Acta Materialia* **58**, 6728 (2010).
- ¹⁵ T. Misaki, A. Wakahara, H. Okada, and A. Yoshida, *Journal of Crystal Growth* **260**, 125 (2004).
- ¹⁶ L. Zhu, P. Maruska, P. Norris, P. Yip, and L. Bouthillette, *MRS Proceedings* **537** (1998), 10.1557/PROC-537-G3.8.
- ¹⁷ K. Du, C. Bekele, C. C. Hayman, J. C. Angus, P. Pirouz, and K. Kash, *Journal of Crystal Growth* **310**, 1057 (2008).
- ¹⁸ T. Misaki, K. Tsuchiya, D. Sakai, A. Wakahara, H. Okada, and A. Yoshida, *Physica Status Solidi (c)* **0**, 188 (2003).
- ¹⁹ R. Viennois, T. Taliercio, V. Potin, A. Errebah, B. Gil, S. Charar, A. Haidoux, and J.-C. Tdenac, *Materials Science and Engineering: B* **82**, 45 (2001).
- ²⁰ L. Zhu, P. Norris, and L. Bouthillette, *MRS Proceedings* **607**, 291 (1999).
- ²¹ N. Zhang, S. Ouyang, T. Kako, and J. Ye, *Chem. Commun.* **48**, 1269 (2012).
- ²² X. Wang, K. Maeda, Y. Lee, and K. Domen, *Chemical Physics Letters* **457**, 134 (2008).
- ²³ F. Tessier, P. Maillard, Y. Lee, C. Bleugat, and K. Domen, *The Journal of Physical Chemistry C* **113**, 8526 (2009).
- ²⁴ Y. Lee, K. Teramura, M. Hara, and K. Domen, *Chemistry of Materials* **19**, 2120 (2007).
- ²⁵ P. Bacher, G. Roult, M. Ghers, O. Merdrignac, J. Guyader, and Y. Laurent, *Materials Chemistry and Physics* **21**, 223 (1989).
- ²⁶ R.-J. Xie and H. T. Bert Hintzen, *Journal of the American Ceramic Society* **96**, 665 (2013).
- ²⁷ A. N. Fioretti, A. Zakutayev, H. Moutinho, C. Melamed, J. D. Perkins, A. G. Norman, M. Al-Jassim, E. S. Toberer, and A. C. Tamboli, *J. Mater. Chem. C* **3**, 11017 (2015).
- ²⁸ A. N. Fioretti, J. Pan, B. R. Ortiz, C. Melamed, P. C. Dippo, L. T. Schelhas, J. D. Perkins, D. Kuciauskas, S. Lany, A. Zakutayev, E. S. Toberer, and A. C. Tamboli, *Materials Horizons* **5**, 823 (2018).
- ²⁹ E. Arca, S. Lany, J. D. Perkins, C. Bartel, J. Mangum, W. Sun, A. Holder, G. Ceder, B. Gorman, G. Teeter, W. Tumas, and A. Zakutayev, *Journal of the American Chemical Society* **140**, 4293 (2018).
- ³⁰ N. Barradas, K. Arstila, G. Battistig, M. Bianconi, N. Dytlewski, C. Jeynes, E. Kótai, G. Lulli, M. Mayer, E. Rauhala, E. Szilágyi, and M. Thompson, *Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms* **266**, 1338 (2008).
- ³¹ L. Giannuzzi and F. Stevie, *Micron* **30**, 197 (1999).
- ³² "See supplemental materials at [url] for additional x-ray pole figures, unprocessed tem data, and atomic force microscopy."
- ³³ N. Itoh, J. C. Rhee, T. Kawabata, and S. Koike, *Journal of Applied Physics* **58**, 1828 (1985).
- ³⁴ Y. Daigo and N. Mutsukura, *Thin Solid Films* **483**, 38 (2005).
- ³⁵ X. R. Huang, J. Bai, M. Dudley, B. Wagner, R. F. Davis, and Y. Zhu, *Physical Review Letters* **95** (2005), 10.1103/PhysRevLett.95.086101.
- ³⁶ D. Skachkov, A. Punya Jaroenjittichai, L.-y. Huang, and W. R. L. Lambrecht, *Physical Review B* **93**, 155202 (2016).

- ³⁷ J. J. S. Scragg, J. K. Larsen, M. Kumar, C. Persson, J. Sandler, S. Siebentritt, and C. Platzer Björkman, *physica status solidi (b)* **253**, 247 (2016).