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# Guiding self-assembly of active colloids by temporal modulation of activity 

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#### Abstract

Self-organization phenomena in ensembles of self-propelled particles open pathways to the synthesis of new dynamic states not accessible by traditional equilibrium processes. The challenge is to develop a set of principles that facilitate the control and manipulation of emergent active states. Here, we report that dielectric rolling colloids energized by a pulsating electric field self-organize into alternating square lattices with a lattice constant controlled by the parameters of the field We combine experiments and simulations to examine spatiotemporal properties of the emergent collective patterns, and investigate the underlying dynamics of the self-organization. We reveal the resistance of the dynamic lattices to compression/expansion stresses leading to a hysteretic behavior of the lattice constant. The general mechanism of pattern synthesis and control in active ensembles via temporal modulation of activity can be applied to other active colloidal systems.


The subject of vast interest in the field of active matter physics is the emergence of self-organized collective behavior out of chaotic motions of individual particles as a result of inter-particle interactions [1-7]. Self-organization phenomena in biological active systems, such as flocks of birds and schools of fish, rely on communications, visual monitoring, sensing of individual positions and requires constant brain processing of collected information [8-11]. A collective motion in ensembles of simple organisms, such as swimming bacteria, may arise only out of inter-particle steric and hydrodynamic interactions [12-18] that makes them a popular model system for the investigation of collective behaviors and selforganization phenomena.

Systems of synthetic self-propelled particles, energized by chemical reactions or electromagnetic fields, provide even better control over individual and collective dynamics of active units and, as a result, complex out-ofequilibrium dynamics of such systems is the subject of extensive research [5, 6, 19-26]. Currently, the majority of collective phenomena are observed in the systems with a constant (or nearly constant) energy injection rate. For instance, a system of electrostatically driven rolling colloids, Quincke rollers, exhibits a large variety of dynamic phases ranging from an isotropic gas to polar bands, vortices, swarms, and rotating clusters [5, 19, 27, 28]. The temporal modulation of activity in that system was recently suggested to model behavior of living systems and control clustering of particles [6, 29, 30].

In this Letter, we report self-assembly of active Quincke rollers into dynamic square lattices when energized by a spatially uniform but modulated in time electric field. The formation of dynamic structures is triggered by a momentarily decoupling between dominant inter-particle interactions - hydrodynamic velocity alignment and near-field electrostatic repulsion as a result of a temporal cessation of the energy injection. Temporal cessation of the activity resets particles' interactions and redirects velocities according to the locally formed parti-
cle arrangements. In contrast to vertically vibrated granular matter [31-34], only square lattices are observed, and these lattices preserve their structure after termination of the activity

In our experiments, spherical polystyrene particles of the diameter $d=4.8 \mu \mathrm{~m}$ are dispersed in $0.15 \mathrm{~mol} \mathrm{~L}^{-1}$ AOT/hexadecane solution and sandwiched between two parallel ITO-coated glass slides spaced $45 \mu \mathrm{~m}$ apart, see Fig. 1(a). When a static (DC) electric field is applied, particles polarize, and above the critical amplitude of the field [19] they start to steadily rotate and roll on the bottom surface with a constant speed driven by electrohydrodynamic Quincke rotation phenomenon [36, 37]. The typical velocity of rollers in our experiments is $0.8 \mathrm{~mm} / \mathrm{s}$. At a particle area fraction above $\phi_{\mathrm{c}} \approx 0.002$, rollers form a steady vortex or a traveling band in confined systems [19, 27, 38]. The temporal profile of the mean square displacement (MSD) for individual rollers (red curve in Fig. 1(d)) reveals long ballistic regime of motion.

The behavior of rollers becomes drastically different if activity of particles is modulated by the pulsating electric field shown in Fig. 1(b). For simplicity, we fix the magnitude of the electric field, but probe the rollers' dynamics by varying only the duration of pulses $\tau_{\text {on }}$ and the interval between the pulses $\tau_{\text {off }}$. This technique was recently suggested to model run-and-tumble behavior and Lévy walks via incomplete depolarization of rollers at relatively short resting time $\tau_{\text {off }}[6]$. The formation of lattices observed in our work requires full depolarization of particles after each cycle of the electric field which occurs if $\tau_{\text {off }}$ is several times larger than the Maxwell-Wagner polarization relaxation time $\tau_{\mathrm{MW}}\left(\tau_{\mathrm{MW}}=\left(\epsilon_{\mathrm{p}}+2 \epsilon_{\mathrm{f}}\right) /\left(\sigma_{\mathrm{p}}+2 \sigma_{\mathrm{f}}\right) \approx 1 \mathrm{~ms}\right)$, where $\epsilon_{\mathrm{p}, \mathrm{f}}$ and $\sigma_{\mathrm{p}, \mathrm{f}}$ are the permittivities and conductivities of particles (p) and fluids (f), respectively. Relaxation of large-scale flows and hydrodynamic interactions occur within viscous timescale that in the experimental cell of the thickness $d=45 \mu \mathrm{~m}$ is of the order of 1 ms . As a result, each roller in our experiments comes to a complete rest for $\tau_{\text {off }}>10 \mathrm{~ms}$. The Brownian motion


FIG. 1. Self-assembly of Quincke rollers into square lattices under a pulsating electric field. (a) A sketch of the experimental setup. (b) Temporal dependence of the electric field and the averaged particles' velocity. $T=\tau_{\text {on }}+\tau_{\text {off }}$ is the period of the signal. (c) A trajectory of an isolated roller energized by the pulsating electric field. (d) Mean square displacement curves of the rollers under periodic activity modulations. The bottom (purple) and the top (red) curves correspond to the motion of the rollers in Gas and Vortex phases respectively. $E=2.7 \mathrm{~V} \mathrm{\mu m}^{-1}$ and $T=100 \mathrm{~ms}$ for all curves except the case of a constant field. Slopes marked as 1 and 2 indicate diffusive $(\sim \tau)$ and ballistic ( $\sim \tau^{2}$ ) regimes respectively. (e) An experimental snapshot of a dynamic square lattice captured between electric field pulses. The blue squares marks four unit cells with a lattices constant $a$. The same area is also shown in (f-h). The particle area fraction $\phi=0.114 ; E=3.0 \mathrm{~V}^{-1} ; T=125 \mathrm{~ms} ; \tau_{\mathrm{on}}=112.5 \mathrm{~ms}$. The scale bar is 0.5 mm . See also Video S 1 [35]. (f) Overlayed square lattices formed after two consecutive cycles represented by blue and red circles. (g) Visualization of the particle trajectories over five periods of the field. For clarity only $10 \%$ of particle trajectories are shown. (h) Color-coded visualization of the time evolution of particles positions during one cycle. The color in each point of a particle trajectory corresponds to the time between 0 (purple) and $\tau_{\text {on }}$ (red) as shown by the colorbar.
has a negligible effect on particles' motion due to their large size and does not alter their positions when activity is terminated. At low concentrations ( $\phi<\phi_{c}$ ), each particle randomizes the velocity direction in each period (Fig. 1(c) and Fig. S9 [35]).

At a high particle concentration, the initial direction of motion upon field application is affected by an arrangement of neighbors. Rollers tend to move against the local gradient of particle density due to the electrostatic repulsion between polarized particles [30]. The spontaneous formation of a cluster is followed by its quick explosive decay and concentration of particles in previously depleted areas during the next cycle. While initially the positions of nuclei clusters are random, they slowly evolves into a well-defined stable structure, and in several hundred activity cycles particles' positions start to alternate between two sets of nearly perfect square lattices (A and B) with the same lattice constant $a$, see Fig. 1(e,f), Fig. S1, and Video S1 [35]. The structure of lattices A and B are geometrically identical but translated by the half of a diagonal of the unit cell $(a / 2, a / 2)$. A closer inspection of particles' motion shows that the average distance particles travel each cycle is noticeably smaller than the shift between the lattices $(\sqrt{2} a / 2)$, see Fig. $1(\mathrm{~g}, \mathrm{~h})$. The particles preserve their positions in one of two lattices as long as the field is off, but quickly form the alternative lattice upon the next pulse of the electric field. Therefore, the observed self-organization of particles is intrin-
sically different from activity-induced cluster formation in which active particles slow down in crowded areas, further increasing the local density and the size of the cluster. $[3,39,40]$. The reciprocating motion of rollers between the two lattices (period is $2 T$ ) effectively traps individual particles within unit cells for several cycles, as reflected by the oscillation of MSD curves, and particle trajectories (Fig. 1d,g). The motion of particles eventually becomes diffusive as particles migrate from cell to cell over the system. The diffusion constant $D$ increases with particle run time $\tau_{\text {on }}$, see Fig. S3(a) in [35].

Our minimalistic phenomenological model (see [35] for details) that involves only isotropic inter-particle repulsion and velocity alignment mechanisms accurately reproduces the formation of patterns observed in experiments (see Fig. S2, S3 in [35]) as well as main dynamic properties of the lattices discussed below. The results of our simulations highlight the role of repulsive forces immediately upon the system activation. In the model, while the system is active all particles move with nearly the same constant speed (small fluctuations due to noise or interactions do not alter the results). The initial direction of motion for each particle at the beginning of each cycle is defined by the net repulsive force from the neighbors. If the initial direction were chosen randomly, particles would not form any stable structure. These findings provide a hint on the mechanism of lattice formation. Upon the system activation, particles start to experience


FIG. 2. (a) Dynamic phases formed by rollers for different $\tau_{\text {on }}$ and $\tau_{\text {off }}$. $\phi=0.120 ; E=2.7 \mathrm{~V} \mathrm{~mm}^{-1}$. The system is confined by a cylindrical well with a diameter $D=1 \mathrm{~cm}$. The phases marked as "*" have been previously reported in Ref. [6]. See Video S3 in [35] for more details. (b) Lattices and Gas formed at different $\tau_{\text {on }}$ and $\phi\left(\tau_{\text {off }} \geq 40 \mathrm{~ms}\right)$. The dashed line is a fit of the boundary between phases by a function $\phi=A / \tau_{\text {on }}$, see text for details. Inset: dynamic phase diagram obtained from simulations. $N$ is the number of particles in the system. (c-d) Dependence of the lattice constant on $\tau_{\text {on }}$ for different area fractions in experiments (c) and simulations (d). $D=$ 1.5 mm . The dashed lines are linear fits for the lowest and highest area fractions. Insets in (c) and (d) experimental and simulation data with the lattice constant re-scaled by a $\sqrt{\phi}$.
repulsion from the neighbors (due to a field-induced polarization) and roll away from dense clusters. A temporal cessation of the activity and subsequent re-energizing of the system lead to rapid reorientation of particles' velocities against the local density gradients, and results in a reciprocating motion of particles between the two sets of lattices.

In order to gain additional insights into the observed self-organization, we explore the response of the system to changes in the pulse duration, $\tau_{\text {on }}$, and distance between the pulses, $\tau_{\text {off }}$, see Fig. 2a and Video S3 in [35]. The lower limit of $\tau_{\text {on }} \approx 20 \mathrm{~ms}$ for the lattices phase comes from a minimum time and distance rollers must travel to interact with each other as well as to overcome the intrinsic positional noise in the system. This limit naturally depends on the particles density and may be significantly longer for low densities as discussed below. For $\tau_{\text {on }}>1 \mathrm{~s}$ the particles' velocities at the end of each cycle become uncorrelated with initial orientations and, therefore, rollers are not able to alternate between fixed stable patterns. The vortex phase in Fig. 2(a) corresponds to the formation of a continuous single vortex at $\tau_{\text {off }}<10 \mathrm{~ms}$. At this phase, the termination of energy injection is so brief that particles do not fully stop. In the pulsed flocks phase the formation of the global vortex is interrupted every period of the signal. At small $\tau_{\text {on }}$
rollers' dynamics resembles a gas motion with a tendency to form clusters as $\tau_{\text {on }}$ further decreases, see Fig. 2(a). The Lattices can be distinguished from the Gas by periodic oscillations of the density profile, see Fig. S4 and S6 in [35]. The crossover between those two dynamic states is continuous and smooth (see Fig. S4). Complex behavior of the system at small $\tau_{\text {off }}$ and $\tau_{\text {on }}$ (marked gray in Fig. 2(a)), corresponds to rollers with incomplete depolarization studied in [6].

Similar to other active systems, the formation of globally correlated states (like lattices) can only be observed at particle densities above a certain threshold. We investigate in experiment and simulations the behavior of the critical particle number density, necessary to facilitate the formation of lattices, as a function of the run time, $\tau_{\text {on }}$. This transition is accessible only at $\tau_{\text {off }}$ and $\tau_{\text {on }}$ time larger than the polarization relaxation time $\tau_{\text {MW }}$. At low densities particles need longer time to reach and interact with the neighbors, and as a result, a longer run time is required to support the dynamic lattice. Correspondingly, as the density of the rollers increases, the minimum run time to form the lattice decreases, see Fig. 2(b). The boundary between the Gas and Lattices follows a simple relation connecting the particle number density, $\phi$, and the run time $\tau_{\text {on }}: \phi=A / \tau_{\text {on }}$. The above scaling is valid for the whole range of scanned parameters ( 10 ms $<\tau_{\text {on }}<400 \mathrm{~ms}$ and $0.015<\phi<0.17$ ), and $A$ here is a constant that depends on the properties of the system (particle size, activity, and liquid media). $A=5.5 \mathrm{~ms}$ for our experimental systems.

The lattice constant $a$ growths with $\tau_{\text {on }}$, see Fig. 2c. Nevertheless, all the lines obtained at different densities collapse into a single line when re-scaled by a $\sqrt{\phi}$, see Fig. 2c inset. Such scaling possibly comes from the similarity of the lattice density profiles at different $\phi$ that scale with the size of the unit cell, see Fig. S12 in [35]. Correspondingly, the evolution of the lattice constant $a$ with the run time $\tau_{\text {on }}$ can be written as $a=k d \tau_{\text {on }} \sqrt{\phi}$, where $k$ is the system dependent constant, the slope of the re-scaled line ( $k=1.63 \mathrm{~ms}^{-1}$ for our system). The simulation results capture the similar trend, see Fig. 2(d). The above scalings can be used to estimate a number of particles per unit cell of the lattice at the transition point between the gas and lattice phases. As the particles' distribution over the system is homogeneous and, therefore, the number of particles per unit cell of the square lattice can be estimated as $n=4 a^{2} \phi /\left(\pi d^{2}\right)$. Substituting $a$ and $\phi$ by its dependence on $\tau_{\text {on }}$ at the transition boundary between the gas and lattice phases, one obtains that the number of rollers per unit cell at the transition is $n_{\mathrm{c}}=4 k^{2} A^{2} / \pi$ (see more details in Note S2) [35]. Since $n_{c}$ is $\phi$ and $\tau_{o n}$ independent, the number of particles per unit cell is the same at any transition point between the gas and lattice phases (the dashed line in Fig. 2b) regardless of the size of the unit cell.

The dynamic square lattices do not have an ideal or-


FIG. 3. Probability distribution functions (PDFs) of the local order parameters $\Psi_{4}$ in a stable lattice formed at different $\tau_{\text {on }}$ in experiments (a) and simulations (b). Insets: Temporal evolution of $\left\langle\Psi_{4}\right\rangle . E=3.0 \mathrm{~V} \mathrm{um}^{-1} ; \tau_{\text {on }}=112.5 \mathrm{~ms} ; \tau_{\text {off }}=$ $12.5 \mathrm{~ms} ; \phi=0.114 ; D=1 \mathrm{~cm}$. Also see Video S4-S5 in [35].
der at macroscopic distances and often develop a number of defects, see Fig. S2 [35]. To characterize the order of self-organized lattices we calculate the local degree of 4 -fold symmetry for each vertex $j$ in a lattice: $\Psi_{4}=1 / 4\left|\sum_{k=1}^{4} \exp \left(4 i \theta_{k}\right)\right|$, where $\theta_{k}$ are polar angles of four closest neighbors with the origin at $j$ vertex. The probability distribution functions ( PDFs ) of $\Psi_{4}$ for stable lattices at different $\tau_{\text {on }}$ show peaks at $\Psi_{4}=1$, indicating a presence of a well defined 4 -fold symmetry in the order, see Fig. 3. The peaks get suppressed with the increase of the run time $\tau_{\text {on }}$, reflecting the emergence of imperfections due to fluctuations of particles' velocities and the development of a global collective motion (vortex). The formation of the local order may be further quantified by a time evolution of the order parameter $\left\langle\Psi_{4}\right\rangle$ defined as $\Psi_{4}$ averaged over the whole system, see Fig. 3(a-b) insets. The temporal evolution of $\left\langle\Psi_{4}\right\rangle$ suggests that dynamic lattices emerge within a hundred cycles, however, they never become geometrically perfect as $\left\langle\Psi_{4}\right\rangle$ reaches the plateau $(\approx 0.8)$. The system does not fully anneal all the defects even after many thousands of activity cycles.

In Quincke rollers system, the mechanism of selforganization with two alternating complementary lattices regulated by a density gradient limits the types of stable lattices to square lattices for any combination of excitation field parameters, while granular vibrated systems often demonstrate transitions between stripes, square, and hexagonal lattices with a change of the driving parameters[31-34, 41, 42]. The fixed geometry of the lattices at a wide range of the driving field parameters enables the study of the collective memory effects [43] in the system. Ramping the control parameter $\tau_{\text {on }}$ up and down at a fixed rate reveals a dynamic hysteresis on $a$ vs $\tau_{\text {on }}$, see Fig. 4(a-b). The size of the hysteresis loop increases with the ramping rate. Interestingly, the lattices demonstrate asymmetry in the resistances against the expansion and compression stresses at a certain range of stresses, see Fig. 4(c-d). The lattices smoothly transition to a smaller size unit cells if $\tau_{\text {on }}$ is decreased by $10 \%$, but break down and reassemble into a larger lattices when $\tau_{\text {on }}$ is increased by the same amount. This is


FIG. 4. Hysteresis loops for different ramping rates of $\tau_{\text {on }}$ in experiments (a) and simulations (b). Arrows indicate the direction of the hysteresis loop. The starting point is $\tau_{\text {on }}=$ 80 ms (experiments) and 30 (simulations). See also video S6 and S7 [35]. (c-d) The evolution of the lattice constant (c) and the local order parameter $\left\langle\Psi_{4}\right\rangle$ of the square lattices (d) in response to the increase (solid symbols) or decrease (open symbols) of the $\tau_{\text {on }}$ relative to the initial value of $\tau_{\text {on }}=120$ ms . The data are averaged over 5 experimental realizations for each curve. Inset in (c) shows the absolute change of $a$ compared to the initial value $a_{0}$.
manifested by a temporal drop of the order parameter, see Fig 4(d). If the relative change of $\tau_{\text {on }}$ is larger than $15 \%$ then lattices prefer to completely disintegrate for both compression and expansion stresses. In addition, the response of the lattices to the compression occurs faster than to expansion if $\tau_{\text {on }} \leq 10 \%$, see Fig 4 (c).

In conclusion, the temporal modulation of activity in the system of Quincke rollers provides a robust technique for accessing and controlling dynamic selforganized states that are not available upon continuous energy injection. A combination of experiments and numerical simulations has been used to investigate the physical mechanism that guides the formation of the reported patterns. A dominant role of electrostatic repulsion over hydrodynamic velocity alignment interactions immediately upon system reactivation results in a reciprocating motion of rollers between two stable square lattices. The lattices are re-configurable with the control of the characteristic lattice constant by the particles' run time. The minimal number of particles per unit cell of the lattice required to facilitate the lattice formation is independent on the lattice constant and the average particle number density. Our results provide new insights into the collective behavior and control of active colloidal ensembles by means of a temporal modulation of activity. The reported mechanism should, in principle, be applicable
to other active systems where the collective behavior is governed by the interplay of isotropic repulsion and hydrodynamic velocity alignment interactions.

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