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Comment on “Inverse Square Lévy Walks are not Optimal Search Strategies for math
 $\frac{d}{m} \geq \frac{2}{n}$ ”

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Comment on “Inverse Square Lévy Walks are not Optimal Search Strategies for $d \geq 2$ ”

It is widely accepted that “inverse square Lévy walks are optimal search strategies because they maximize the encounter rate with sparse, randomly distributed, replenishable targets” [1], when the search restarts in the vicinity of the previously visited target, which becomes revisitable again with high probability, i.e., non-destructive foraging [2]. Three objecting claims are raised in Ref. [1] for $d \geq 2$: (i) the capture rate η has linear dependence on the target density ρ for all values of the Lévy index α ; (ii) “the gain η_{\max}/η achieved by varying α is bounded even in the limit $\rho \rightarrow 0$ ” so that “tuning α can only yield a marginal gain”; (iii) depending on the values of the radius of detection a , the restarting distance l_c and the scale parameter s , the optimum is realized for a range of α .

We agree with claim (i), but as we will see, it is not relevant in $d \geq 2$ to whether or not inverse square Lévy walk searches are optimal for non-destructive foraging. Claim (iii) is also correct, however this claim was made already in Refs. [2–5]. In particular Ref. [2] showed that $\alpha = 1$ is optimal only in the limit $l_c \rightarrow a$, which is the main condition of non-destructive foraging, with the quantity l_c in Ref. [1] being none other than the distance r_o in Ref. [2]. Otherwise for large l_c the optimal strategy in the limit $\rho \rightarrow 0$ is to go along straight lines, i.e. $\alpha \rightarrow 0$. Moreover, it is known since 2003 that a range of α can be optimal (see Fig. 1 of Ref. [3], Figs. 2–3 of Ref. [4] and Figs. 1 and S1 of Ref. [6], none of which are cited in Ref. [1]). Crucially, claims (i) and (iii) do not *per se* contradict the main finding of Ref. [2] that $\alpha = 1$ is optimal under the specific conditions of non-destructive foraging (or of destructive foraging in patchy landscapes) [2–10].

To test claim (ii), we have numerically simulated the identical model proposed in Ref. [1] (see Fig. 1). The scaling for η with ρ proposed in Ref. [2] and proved in Ref. [8] for $d = 1$ does not hold in $d = 2$, in agreement with Ref. [1]. However, we find, for small enough $\delta = l_c/a - 1$, that η develops a maximum at $\alpha = 1$ with an arbitrarily large gain relative to the ballistic ($\alpha \rightarrow 0$) and Brownian ($\alpha = 2$) limits, contradicting claim (ii) about “marginal gain” in Ref. [1].

The main problem with Ref. [1] is that Eq. (3) fails in the limit $l_c \rightarrow a$ of non-destructing foraging. Eq. (3) yields a gain $K_d \sim 1/[A(a^\beta - B l_c^\beta)]$ in Eq. (5), with $\beta = -1$ for $\alpha < 1$ and $\beta = \alpha - 2$ for $\alpha > 1$. This gain, which agrees with claim (ii), is wrong in the limit $l_c \rightarrow a$.

Finally, we present a heuristic argument for the correct scaling of K_d for $d = 2$ when $l_c \rightarrow a$. Note that l_c is the distance at which the target stops hiding. The limit $\delta \rightarrow 0$ has biological relevance in this “hide-and-seek” model [10]. Let $\sigma = s/a$ and $\eta_0(\alpha, \delta, \rho, \sigma) =$

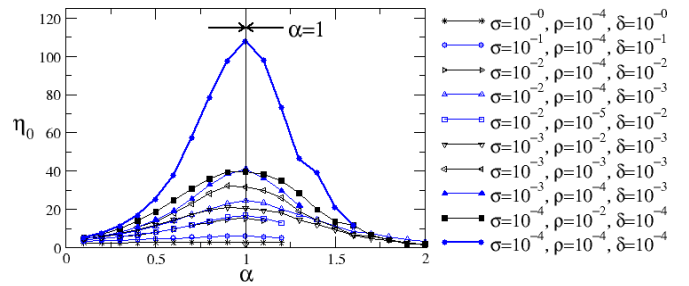


Figure 1. $\eta_0 = \eta/(\rho a)$ vs. α , for $N = 10^6$ Poisson distributed targets on a square of size $\sqrt{N/\rho}$ with periodic boundary conditions, averaged over 10^5 targets found.

$\eta/(\rho a)$. When $\delta \rightarrow 0$, the (radial) motion of the forager near the border of the detection circle is essentially one dimensional, hence the rigorous theory of the Riesz operator [8] on the interval of length L with absorbing ends becomes applicable. For $\sigma > \delta$ the efficiency increases when σ decreases because there are fewer large jumps away from the previous target that make re-encountering it difficult. When $\sigma \approx \delta$ the efficiency reaches its maximum. In the limit $\sigma \approx \delta \rightarrow 0$ we expect the same scaling behavior as in $d = 1$: $\eta_0 \sim \delta^{-\alpha/2}$ for $\alpha < 1$ and $\eta_0 \sim \delta^{-1+\alpha/2}$ for $\alpha > 1$. Hence η_0 has an arbitrarily strong maximum at $\alpha = 1$ when $\sigma \approx \delta \rightarrow 0$, in agreement with Fig. 1, and in disagreement with the title and claim (ii) of Ref. [1], restoring thus the original result for non-destructive foraging in Ref. [2] of the optimality of inverse square Lévy flights.

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- [1] N. Levernier, J. Textor, O. Bénichou and R. Voituriez, Phys. Rev. Lett. **124**, 080601 (2020).
 - [2] G. M. Viswanathan, S. V. Buldyrev, S. Havlin, M. G. E. da Luz, E. P. Raposo and H. E. Stanley, Nature **401**, 911 (1999).
 - [3] E. P. Raposo, S. V. Buldyrev, M. G. E. da Luz, M. C. Santos, H. E. Stanley and G. M. Viswanathan, Phys. Rev. Lett. **91**, 240601 (2003).
 - [4] M. C. Santos, E. P. Raposo, G. M. Viswanathan and M. G. E. da Luz, Europhys. Lett. **67**, 734 (2004).
 - [5] S. A. Sotelo-López, M. C. Santos, E. P. Raposo, G. M. Viswanathan and M. G. E. da Luz, Phys. Rev. E **86**, 031133 (2012).
 - [6] F. Bartumeus and S. A. Levin, Proc. Natl. Acad. Sci. U.S.A. **105**, 19072 (2008).
 - [7] S. V. Buldyrev, S. Havlin, A. Ya. Kazakov, M. G. E. da Luz, E. P. Raposo, H. E. Stanley and G. M. Viswanathan, Phys. Rev. E **64**, 041108 (2001).
 - [8] S. V. Buldyrev, M. Gitterman, S. Havlin, A. Ya. Kazakov, M. G. E. da Luz, E. P. Raposo, H. E. Stanley and G. M. Viswanathan, Physica A **302**, 148 (2001).
 - [9] M. G. E. da Luz, E. P. Raposo and G. M. Viswanathan, Phys. Life Rev. **14**, 94 (2015).
 - [10] A. M. Reynolds and F. Bartumeus, J. Theor. Bio. **260**, 98 (2009).

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