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Phys. Rev. Lett. 124, 112301 — Published 17 March 2020
DOI: 10.1103/PhysRevLett.124.112301
Multi-gluon correlations and evidence of saturation from dijet measurements at an Electron Ion Collider

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(Dated: February 6, 2020)

We study inclusive and diffractive dijet production in electron-proton and electron-nucleus collisions within the Color Glass Condensate effective field theory. We compute dijet cross sections differentially in both mean dijet transverse momentum |P| and recoil momentum |Δ|, as well as the anisotropy in the relative angle between |P| and |Δ|. Our results cover a much larger kinematic range than accessible in previous computations performed in the correlation limit approximation, where it is assumed that |P| ≫ |Δ|. We validate this approximation in its range of applicability and quantify its failure for |P| ≲ |Δ|. We also predict significant target-dependent deviations from the correlation limit approximation for |P| > |Δ| and |P| ≲ Qs, which offers a straightforward test of gluon saturation and access to multi-gluon distributions at a future electron ion collider.

Introduction. To gain a complete understanding of the complex multi-parton structure of nuclei at small x, measurements of a multitude of processes in high energy e + p(A) collisions over a wide range of kinematics are necessary. A future electron ion collider (EIC) [1–3] will provide an ideal tool for such an endeavor, with dijet production being one of the most important processes to access the structure of gluon fields and their non-linear dynamics inside protons and heavier nuclei.

While coherent diffractive dijet production allows to access the target’s spatial geometry [4–8], inclusive and incoherent diffractive dijet cross sections are sensitive to multi-gluon correlations in the target [9, 10] (see also [11–15]). In the back-to-back correlation limit approximation (CLA), where the mean dijet momentum is much larger than the recoil momentum, the inclusive dijet production cross section can be expressed in terms of the Weizsäcker-Williams transverse momentum dependent gluon distributions (TMDs), allowing experimental access to these fundamental quantities [9, 16–19]. We advocate going beyond this limit to allow for deeper insights into the multi-gluon structure of the nucleus. Inclusive and diffractive (incoherent) dijets are sensitive to the quadrupole and dipole-dipole correlators of light-like Wilson lines, respectively. These are among the fundamental objects describing the gluon structure at small x.

We present the first evaluation of inclusive and incoherent diffractive dijet cross sections and their azimuthal anisotropies for general small-x kinematics in the Color Glass Condensate (CGC) effective field theory (EFT) [20, 21], at leading order in αs, resumming all terms ∼ αs ln 1/x. For inclusive dijets, our results explicitly validate the CLA in the kinematic region |P| ≫ |Δ| and extend our knowledge of dijet production to the region |P| ≲ |Δ|, where deviations from the CLA turn out to be large. We further show that corrections to the CLA also become important when |P| ≲ Qs, even when |P| > |Δ| holds. These corrections, enhanced by the saturation scale Qs, probe genuine multi-gluon correlations [22, 23], and are not encompassed by the resummed kinematic twists of the improved TMD framework [24] (see also [25–30] for forward dijets in dilute-dense hadronic collisions, and experimental measurements from RHIC [31, 32]).

Calculations of the elliptic anisotropy employing multi-gluon correlators deviate strongly from the CLA for |P| ≲ |Δ|. In particular, for transverse polarization the calculated elliptic modulation is qualitatively different from that in the CLA, as a maximum appears both as a function of |P| and |Δ|.

For the first time within the CGC EFT (see also [33–35]), we predict the incoherent diffractive cross section. Our calculation predicts characteristic features of the cross section’s elliptic anisotropy as a function of |P| and |Δ|, involving sign changes and minima, which should be observable experimentally.

We compute the fraction of diffractive dijet events as a function of the mean dijet momentum. It increases with the mass number of the nucleus and decreases with Q2 at a slower rate than expected in the small dipole expansion, signaling gluon saturation [36].

Dijet production in high energy DIS. In the dipole picture of high energy deeply inelastic scattering (DIS), the production of a forward q ¯q dijet can be seen as the splitting of a virtual photon γ∗ into a quark-antiquark dipole and its subsequent eikonal scattering off the target’s color field. We work in a frame in which the virtual photon and nucleon in the target have zero transverse momenta [37]. The photon has virtuality Q2 and four momentum qμ = (−Q2/2q−, q−, 0). Neglecting its mass, the nucleon has energy En and four momentum Pn = (√2En, 0, 0). The center of mass energy of the virtual photon-nucleon system is W. The transverse mo-
menta of the outgoing quark and antiquark are $p_1$ and $p_2$, their longitudinal momentum fractions are $z_1$ and $z_2$, with $z_i = p_i^+/q = 2E_i/p_i e^{-y_i}/W^2$, where $p_i^+$ and $y_i$ are the quark and antiquark longitudinal momenta and rapidities in this frame, respectively.

\[
\frac{d\sigma^{A\rightarrow q\bar{q}X}}{dy_1 dy_2 dq^2 P^4 \Delta} = 8\alpha_e Z_F^2 N_c S\frac{z_1^3 z_2^3 Q^2}{(2\pi)^6} \int e^{-iP\cdot(r-r')} e^{-i\Delta\cdot(b-b')} \mathcal{O}_{r,b,r',b'} K_0(\varepsilon_f |r|) K_0(\varepsilon_f |r'|),
\]  

(1)

\[
\frac{d\sigma^{A\rightarrow q\bar{q}X}}{dy_1 dy_2 dq^2 P^4 \Delta} = 2\alpha_e Z_F^2 N_c S\frac{z_1 z_2 (z_1^2 + z_2^2)}{2\pi^6} \int e^{-iP\cdot(r-r')} e^{-i\Delta\cdot(b-b')} \mathcal{O}_{r,b,r',b'} |r||r'| K_1(\varepsilon_f |r|) K_1(\varepsilon_f |r'|).
\]  

(2)

Here, $\alpha_e = e^2/(4\pi)$ is the electromagnetic coupling, $N_c = 3$ is the number of colors, $\delta_2 = \delta(1-z_1-z_2)$, $c_2^2 = z_1 z_2 Q^2$, and $\int d^2 \mathbf{x}$. We use $Z_F^2 = (\frac{2}{3})^2 + (-\frac{1}{3})^2 + (-\frac{1}{3})^2$, corresponding to $u$, $d$ and $s$ quarks. Assuming a homogeneous target, the cross section is proportional to the effective transverse area of the target $S_{\perp}$. The multi-gluon correlations are encoded in $O$, defined as

\[
O_{r,b,r',b'}^{(4)} = 1 - S_{x_1,x_2}^{(2)} - S_{x_1,x_2}^{(2)} + S_{x_1,x_2,x_1',x_1'}^{(4)}
\]  

(3)

for inclusive production, and

\[
O_{r,b,r',b'}^{(2,2)} = 1 - S_{x_1,x_2}^{(2)} - S_{x_1,x_2}^{(2)} + S_{x_1,x_2,x_1',x_1'}^{(2,2)}
\]  

(4)

for total diffractive (color singlet) production. The $x$ coordinates are related to $r$ and $b$ via $x_{1,2} = b \pm z_{2,1} r$ and $x_{1,2}' = b' \pm z_{2,1} r'$. The dipole, dipole-dipole, and quadrupole correlators of fundamental light-like Wilson lines $V$ are defined by $O_{[9,39]}$

\[
S_{x_1,x_2}^{(2)} = \frac{1}{N_c} \langle \text{tr} (V_{x_2} V_{x_1}) \rangle,
\]  

(5)

\[
S_{x_1,x_2,x_1',x_1'}^{(2,2)} = \frac{1}{N_c} \langle \text{tr} (V_{x_2} V_{x_1}) \text{tr} (V_{x_2} V_{x_1}) \rangle,
\]  

(6)

\[
S_{x_1,x_2,x_1',x_1'}^{(4)} = \frac{1}{N_c} \langle \text{tr} (V_{x_2} V_{x_1} V_{x_2} V_{x_1}) \rangle.
\]  

(7)

where the $\langle \cdot \rangle$ denote the average over static large $x$ color source configurations in the CGC EFT.

The correlators $O_{r,b,r',b'}$ contain both the elastic $[40]$ and inelastic parts. In this work we neglect the impact parameter dependence of the target such that the elastic cross section vanishes at non-zero $\Delta$. This amounts to the replacements $O_{r,b,r',b'}^{(4)} \rightarrow S_{x_1,x_2,x_1',x_1'}^{(4)}$, and $O_{r,b,r',b'}^{(2,2)} \rightarrow S_{x_1,x_2,x_1',x_1'}^{(2,2)}$, which restrict the cross sections to the inelastic part.

The correlators above are evaluated at $x = (Q^2+|\Delta|^2+M_0^2)/\sqrt{W}$, which follows from kinematics and energy-momentum conservation [9, 19], where the invariant mass of the dijet is given by $M^2_{q\bar{q}} = |P|^2/(z_1 z_2)$.

Expressed using the momenta $P = z_2 p_1 - z_1 p_2$ and $\Delta = \pmb{p}_1 + \pmb{p}_2$, at leading order in $\alpha_s$, the cross sections for dijet production of massless quarks for longitudinal ($L$) and transverse ($T$) photon polarization read $[9, 16, 38]$

To reduce the computational cost of our calculation, we employ the nonlinear Gaussian approximation $[9, 39, 41–44]$, which allows one to express any $n$-point correlator of light-like Wilson lines as a non-linear function of the dipole correlator in Eq. (5), and was shown to approximate the full quadrupole operator very well $[41, 45]$, even after JIMWLK small $x$ evolution $[46–52]$ for many units in rapidity. Gaussian approximation expressions for the dipole-dipole and quadrupole correlators are given in $[9, 39]$ and are summarized in $[53]$. For the dipole correlator we use the solution of the running coupling Balitsky-Kovchegov equation $[54–56]$ with McLerran-Venugopalan model initial conditions $[57]$. The parameters were obtained by fitting HERA deep inelastic scattering data $[58]$ (see $[59]$ and $[53]$ for details).

**Cross section and elliptic anisotropy.** We present results for the angle averaged cross section and elliptic anisotropy for inclusive and diffractive dijet production in the scattering of longitudinally and transversely polarized photons with virtuality $Q^2 = 10$ GeV$^2$ off nuclear targets and center of mass energy of the photon-nucleon system $W = 90$ GeV. These are defined as follows $[60]$

\[
\frac{d\sigma^{A\rightarrow q\bar{q}X}_{L/T}}{d\Pi} = \int \frac{d\theta_P}{2\pi} \frac{d\Delta}{2\pi} \frac{d\sigma^{A\rightarrow q\bar{q}X}_{L/T}}{dy_1 dy_2 dq^2 P^4 \Delta},
\]  

(8)

and

\[
\frac{d\sigma^{A\rightarrow q\bar{q}X}_{2,L/T}}{dy_1 dy_2 dq^2 P^4 \Delta} = \int \frac{d\theta_P}{2\pi} \frac{d\Delta}{2\pi} \frac{d\sigma^{A\rightarrow q\bar{q}X}_{2,L/T}}{dy_1 dy_2 dq^2 P^4 \Delta}.
\]  

(9)

We study proton and gold targets and in the inclusive case compare to the CLA. Additionally, we predict the ratio of diffractive to inclusive events as a function of dijet momentum for different targets and $Q^2$. All results are for fixed $z_1 = z_2 = 0.5$.

**Inclusive dijets.** In Fig. 1 we show the angular averaged cross sections, comparing results using full multi-particle correlators Eqs. (1) and (2) (solid lines) to the
CLA [18, 19, 53] (dashed lines). The former are valid for any value of \( \Delta \), while the latter are expected to be valid only for \( |P| \gg |\Delta| \). The expected agreement between the CLA and the more general result at \( |P| \gg |\Delta| \) is clearly confirmed in all cases. Deviations from the CLA become large when extrapolated to the regime \( |\Delta| > |P| \).

Importantly, we observe significant deviations from the CLA at \( |\Delta| < |P| < 1.5 \) GeV for the gold target, and much milder deviations for the proton. This difference is explained by saturation effects: The cross sections beyond the CLA receive genuine saturation corrections of order \( Q_2^2/|P|^2 \) and \( Q_2^2/Q^2 \), in addition to kinematic corrections of order \( |\Delta|^2/|P|^2 \) [22, 23]. This observation demonstrates that inclusive dijet production in e+A collisions at a future EIC can provide direct access to gluon saturation.

In Fig. 2 we present the elliptic modulation of the cross section in the angle between \( \vec{P} \) and \( \Delta \). Again, the CLA provides a good estimate in the region \( |P| \gg |\Delta| \), and deviations become large for \( |\Delta| \sim |P| \). We predict a minimum \( \nu_{TT} \sim -30\% \) for proton targets in the range \( |P| \sim |\Delta| \sim 1.8 \) GeV, and \( \nu_{TT} \sim -20\% \) for gold for \( |P| \sim |\Delta| \sim 2.2 \) GeV, unlike the CLA, which predicts decreasing values of \( \nu_{TT} \) as \( \Delta \) increases [53]. To probe these, and the aforementioned saturation effects, experiments should focus on the kinematics \( |P| \sim |\Delta| \). We further confirm the large elliptic modulation for the longitudinally polarized photon, which was obtained previously in the CLA [18, 19, 53].

**Diffractive dijets.** We show results of diffractive dijet cross sections and elliptic anisotropies for virtual photon off gold scattering in Fig. 3. Although our results contain only incoherent diffraction, these are the most dominant at momentum transfer \( \Delta > 1/R_A \) (\( \sim 0.2 \) GeV for gold), such that the result is approximately equal to the total diffractive cross section.

Comparing the inclusive (Fig. 1) and diffractive cross sections (Fig. 3), we observe a strong suppression of diffractive events and a different \( |P| \)-dependence for the longitudinal and transverse cases. Theoretically, this can be directly related to the properties of multi-gluon correlators in the target. The only difference between the inclusive and diffractive cross sections are the different color structures of the correlators \( \mathcal{O} \). A small dipole expansion explains the effect of this difference: The first non-vanishing term in the expansion occurs at linear...

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**FIG. 1.** Angle averaged inclusive dijet cross section for proton (upper) and gold (lower) targets. Solid lines: full multiparticle correlator result. Dashed lines: CLA. Contour labels represent powers of 10. Panels on the left show a vertical section of the contour plots at fixed \( |\Delta| = 1 \) GeV.
order for the inclusive case and at quadratic order for diffractive production, because diffractive events require at least two gluons exchanged in the amplitude to ensure color neutrality.

The elliptic modulation of the incoherent diffractive cross section is shown in the middle panel of Fig. 3. For both polarizations it exhibits a sign change as a function of $|P|$, similar to that observed in coherent diffractive dijet production [6, 7]. The transverse case also shows a sign change in $|\Delta|$ for $|P| \gtrsim 2$ GeV. Importantly, the elliptic modulation reaches large values (tens of percent) in the studied kinematic range.

In the right panels of Fig. 3 we show the ratio of diffractive to inclusive events as a function of $|P|$ for fixed $|\Delta| = 1.5$ GeV. The fraction of diffractive events increases with the target saturation scale $Q_s$ from proton to gold, and decreases with increasing photon virtuality $Q^2$. An expansion in small dipoles predicts the fraction of diffractive events to increase as $Q_s^2$. We expect a factor of 2.6 increase in the considered kinematics after BK evolution from proton to gold; however, we find a smaller increase of 1.9 (2.3) for transversely (longitudinally) polarized photons at $|P| \approx 1$ GeV and $Q^2 = 4$ GeV$^2$, with a mild increase towards the expected value of 2.6 with growing $|P|$. This behavior indicates effects of gluon saturation, which are stronger in larger nuclei. We argue that this ratio is a key measurement at a future EIC, allowing to quantify gluon saturation (differentially in $|P|$ and $Q^2$).

Conclusions. We computed inclusive and (incoherent) diffractive dijet production cross sections in e+p and e+A collisions at a future EIC within the CGC EFT. These cross sections are sensitive probes of multi-gluon correlations inside nuclear targets at small $x$ and allow to quantitatively probe gluon saturation experimentally.

Our approach is not restricted to $|P| \gg |\Delta|$ and significantly increases the theoretically accessible kinematic range. We employed the non-linear Gaussian approximation, using dipole correlators obtained from rcBK fits to HERA data. We validated the CLA in inclusive dijet production for most $|P| \gg |\Delta|$, but found significant target dependent corrections for $|P| \lesssim |\Delta|$ or $|P| \lesssim Q_s$, the latter being caused by gluon saturation effects. We thus argue that the regime of moderate $|P| \sim Q_s$ of the target is particularly interesting when studying dijet production at a future EIC. Differential measurements in $P$
and $\Delta$ within a range that includes $Q_s$, will allow to reveal the complex multi parton structure of nuclei and uncover saturation.

We presented the first calculation of diffractive dijet cross sections and their elliptic modulation within the CGC EFT. We studied the nuclear modification of the ratio between the differential inclusive and diffractive dijet cross sections by comparing gold to proton targets at different values of $Q^2$. The dependence of the ratio between the cross sections on the target’s saturation momentum indicates that saturation effects are significant in the studied kinematic regime.

In future work, we plan to include parton showers, hadronization, and full jet reconstruction. Based on results in [19], we expect the $v_2$ of the produced $q\bar{q}$ pair presented here to be a good estimator of the observable dijet $v_2$. It will also be important to include next-to-leading order (NLO) corrections, both in small-$x$ evolution equations: NLO BK [61–64] or NLO JIMWLK [65, 66], and the NLO impact factor [67–71], and to consider the effects of soft gluon radiation of the final state jets that is not captured by the jet algorithm [8].

Detailed extraction of multi-gluon correlators in nuclei and experimental confirmation of gluon saturation will likely require complex global fits to a wide variety of experimental data. We have demonstrated that inclusive and diffractive dijet production are two of the most important processes to consider.

**Acknowledgments.** We thank Elke-Caroline Aschenauer, Renaud Boussarie, Kaushik Roy, Sören Schlichting, Vladimir Skokov, Alba Soto-Ontoso, and Raju Venugopalan for useful discussions. H.M. is supported by the Academy of Finland project 314764. N.M., F.S., and B.P.S. are supported under DOE Contract No. DE-SC0012704. N.M. is funded by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) - project 404640738. This research used resources of the National Energy Research Scientific Computing Center, which is supported by the Office of Science of the U.S. Department of Energy under Contract No. DE-AC02-05CH11231.

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[34] V. Guzey and M. Klasen, “A fresh look at factorization breaking in diffractive photoproduction of dijets at
The elastic (coherent) production of dijets is given by Eqs. (1) and (2) with \( \mathcal{O}_{\nu, b, r', \nu'} = 1 - S_{a_1, a_2}^{(2)} S_{a_1', a_2'}^{(2)} + S_{a_1, a_2}^{(2)} S_{a_1', a_2'}^{(2)} \).

[60] The differential production of dijets is given by Eqs. (1) and (2) with \( \mathcal{O}_{\nu, b, r', \nu'} = 1 - S_{a_1, a_2}^{(2)} S_{a_1', a_2'}^{(2)} + S_{a_1, a_2}^{(2)} S_{a_1', a_2'}^{(2)} \).


