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Evidence for a New Component of High-Energy Solar Gamma-Ray Production

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The Sun is a bright source of multi-GeV γ-rays, with emission observed both from its halo — due to cosmic-rays electrons interacting with solar photons — and its disk — due to hadronic cosmic rays (mostly protons) interacting with solar gas. (Emission from solar particle acceleration is only bright during flares and has not been observed above 4 GeV [1–8].) Although the halo emission [9] agrees with theory [10–12], the disk emission does not, and hence is our focus.

Previously, the most extensive analysis of solar disk γ-ray emission was based on Fermi-LAT data from 2008–2014 [13] (for earlier work, see Refs. [9, 14]), and produced three results. First, the flux is bright, e.g., at 10 GeV, it exceeds the flux expected from Earth-directed cosmic rays interacting with the solar limb by a factor ≥50 [15]. Second, it continues to 100 GeV, requiring proton energies ∼1000 GeV. Third, the 1–10 GeV flux is anti-correlated with solar activity, and is ∼2.5× larger at solar minimum than maximum. The only theoretical model of disk emission is the 1991 paper of Seckel, Stanev, and Gaisser (SSG) [17], which proposes that magnetic flux tubes can reverse incoming protons deep within the solar atmosphere, where they have an appreciable probability of producing outgoing γ-rays. Though this enhances the γ-ray flux, the SSG prediction still falls a factor ~6 below the data at 10 GeV, and does not explain the time variation.

In this and a companion paper, we perform new analyses of Fermi-LAT data based on a longer exposure (2008–2017), better data quality (Pass 8), and improved methods. In Ref. [18], we focus on the 1–100 GeV spectrum and its time variability. Compared to our earlier work [13], the flux is robustly detected up to 100 GeV, and the anticorrelation of the flux with solar activity is detected up to ∼30 GeV. Intriguingly, we discover a spectral dip between 30–50 GeV. This dip is unexpected and its origin is unknown. Here we extend Refs. [13, 18] by going to higher energies, studying the time variation in a new way, and performing the first analysis of flux variations across the resolved solar disk. In the following, we detail our methodology, highlight key discoveries, and discuss possible theoretical implications.

The importance of this work is manifold. Because the disk γ-ray emission is brighter than expected, it motivates new searches with Fermi-LAT [19], the HAWC γ-ray experiment [20], and the IceCube neutrino observatory [21]. The results will yield valuable insights into the dynamic solar magnetic environment, from cosmic-ray modulation in the solar system to the fields deep within the photosphere. They will also advance searches for new physics [22–30]. Most generally, these searches provide the highest-energy data available to understand the Sun as an example of other stars.

Methodology.— We utilize Pass 8 Source events from August 4, 2008 to November 5, 2017 (MET: 239557417–531557417), employing standard cuts. We include events exceeding 10 GeV observed within 0.5° of the solar center (the Sun’s angular radius is 0.26°). The excellent angular resolution of >10 GeV γ-rays minimizes the flux lost from our Region of Interest (ROI). In Appendix A, we show that larger ROIs produce consistent results. We remove events observed when the Sun falls within 5° of the Galactic plane, due to the larger diffuse background. This cut is smaller than in previous work, but is sufficient due to the small ROI. We perform the first conversion of each γ-ray to Helioprojective coordinates utilizing sunpy [31] and astropy [32]. We ignore diffuse backgrounds, which we found in Ref. [18] to be negligible.

We calculate the Fermi-LAT exposure at the solar position in temporal bins of 5000 s (but use precise times for recorded events). Within this period, the Sun moves <0.1° in the Fermi coordinate system, and the Fermi-LAT effective exposure is approximately constant. We assume a single exposure over

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1 The “cycle 24” solar minimum approximately spanned 2007 through 2009. Fermi-LAT science observations (beginning in August, 2008), observed only the latter half of this period [16]
the full ROI in each time-bin, and bin the exposure into 32 logarithmic energy bins spanning 10 GeV to 1 TeV. Because the Sun occupies a unique position in instrumental ϕ-space, we calculate exposures calculated using 10 independent ϕ-bins. In Appendix B, we show that this ϕ-dependence does not affect our results.

Flux, Spectrum and Time Variation.—In Figure 1, we show the solar γ-ray flux before and after January 1, 2010, which roughly corresponds to the end of the Cycle 24 solar minimum. We note three key results.

- The γ-ray flux significantly exceeds the SSG prediction (based on a proton interaction probability of 0.5%), in fact approaching the maximum allowed solar disk flux (for a detailed calculation, see Appendix E).
- The 30–50 GeV spectral dip, which we will carefully examine in Ref. [18], is statistically significant both during and after solar minimum, though there is some evidence (2.5σ) that the dip deepens at solar minimum. Aside from the dip, the spectra in both time periods are significantly harder than predicted by SSG.
- The strongest time variation is observed between solar minimum (largest flux), and the remaining solar cycle. At low energies this variation is moderate [13, 14, 18]. However, the amplitude increases with energy above 50 GeV, reaching a factor ≥10 above 100 GeV.

None of these observations were anticipated by theory.

Morphology.—The large γ-ray flux suggests that a large fraction of the solar surface participates in γ-ray emission. To further elucidate the γ-ray generation mechanism(s), we resolve the γ-ray morphology across the solar surface. This reconstruction is possible at high (≥10 GeV) energies due to the excellent (∼0.1°) Fermi-LAT angular resolution.

In Figure 2, we show the location of γ-rays in our analysis, dividing the data into two temporal bins (before and after January 1, 2010; corresponding to the end of the solar minimum), and two energy bins (below and above 50 GeV; corresponding to the spectral dip discussed in Ref. [18]). We find that, contrary to the SSG model, the emission is neither isotropic nor time-invariant. Instead, it includes distinct polar and equatorial components, with separate time and energy dependences. In particular, it is apparent that γ-rays above 50 GeV are predominantly emitted near the solar equatorial plane during solar minimum, but are emitted from polar regions during the remaining cycle.

We utilize two separate methods to quantify the significance of this morphological shift. The first employs a Kolmogorov-Smirnov test to differentiate the distribution of γ-rays in observed helioprojective latitude (|$T_y$|) during and after solar minimum. This method is model independent, but loses sensitivity to convolving factors such as the instrumental PSF. Below 50 GeV, we find that the event morphologies are consistent to within 1.1σ. However, above 50 GeV, we reject the hypothesis that the latitude distributions during and after solar minimum are equivalent at 2.8σ. This provides reasonable evidence for a morphological shift.

Second, we define a two-component model of the solar surface, with equal-area equatorial and polar emission components (divided at |$T_y$| = ±0.108°). We fit the flux from each component, utilizing the angular reconstruction of each observed γ-ray (see Appendix F). This correctly accounts for the PSF, but provides results that depend on the assumed emission model. In Appendix G we show that different models produce similar results. This analysis provides two key results.

- At all energies, the γ-ray emission becomes more polar after solar minimum. However, the amplitude of this shift increases significantly at high energies.
- The morphological shift is produced by a significant decrease in the equatorial flux after solar minimum, while the polar flux remains relatively constant. Most significantly, at energies >50 GeV, the equatorial fluxes during and after solar minimum are inconsistent at 4.7σ.

In Figure 2, we also show the polar and equatorial spectra during and after solar minimum. While the polar emission spectrum remains relatively constant, the equatorial spectrum softens substantially after solar minimum. This significantly decreases the high-energy equatorial flux after solar minimum, despite the similar normalization of the equatorial component at low energies. Intriguingly, the equatorial γ-ray spectrum during solar minimum is extremely hard, and is consistent with dN/dE ~ E⁻² up to energies exceeding 100 GeV.
FIG. 2. (Top) The location and energy of solar $\gamma$-rays in Helioprojective coordinates. Data are cut into two temporal and two energy bins. The solid disk indicates the solar circle, dashed circle indicates the 0.5$^\circ$ ROI. The average 68% containment region of $\gamma$-rays in each bin is depicted in the top left. The histogram depicts the $T_y$ positions of photons compared to the expectation from isotropic solar emission smeared by the PSF (orange line). Events $>100$ GeV are marked with triangles, rather than circles. We stress that the exposure after solar minimum significant exceeds the exposure during solar minimum. Thus the observed number of counts does not indicate the relative flux. In each bin, we report the flux from the modeled polar and equatorial components, as described in the text. (Bottom) The energy spectrum of polar and equatorial emission, divided into regions during (left) and after (right) solar minimum. The polar emission is approximately constant, while the equatorial emission decreases drastically after solar minimum.

We note that we have combined high-energy spectral bins during solar minimum to improve the statistical separation of polar and equatorial components.

Flux Above 100 GeV.— In Figure 1, we discovered a bright $\gamma$-ray flux above 100 GeV during solar minimum, but found no events in the remaining solar cycle. In Table I, we examine each $>100$ GeV event. We uncover no significant concerns regarding the event classes, or angular and energy reconstructions. All six events pass the UltraCleanVeto event cut, providing the highest confidence that they are $\gamma$-rays. We calculate the probability that each event has a non-solar origin by calculating the $\gamma$-ray flux above 100 GeV in each ROI during periods when Sun is not present. We find that diffuse contributions cannot explain these events. The diffuse $\gamma$-ray flux
<table>
<thead>
<tr>
<th>Date</th>
<th>E (GeV)</th>
<th>Distance</th>
<th>Event Class</th>
<th>P6</th>
<th>P7</th>
<th>BG Prob</th>
</tr>
</thead>
<tbody>
<tr>
<td>2008-11-09</td>
<td>212.8</td>
<td>0.068</td>
<td>UltraCleanVeto</td>
<td>✓</td>
<td>✓</td>
<td>0.00050</td>
</tr>
<tr>
<td>2008-12-13</td>
<td>139.3</td>
<td>0.126</td>
<td>UltraCleanVeto</td>
<td>X</td>
<td>X</td>
<td>0.00038</td>
</tr>
<tr>
<td>2008-12-13</td>
<td>103.3</td>
<td>0.399</td>
<td>UltraCleanVeto</td>
<td>X</td>
<td>X</td>
<td>0.00052</td>
</tr>
<tr>
<td>2009-03-22</td>
<td>117.2</td>
<td>0.255</td>
<td>UltraCleanVeto</td>
<td>✓</td>
<td>✓</td>
<td>0.00027</td>
</tr>
<tr>
<td>2009-08-15</td>
<td>138.5</td>
<td>0.261</td>
<td>UltraCleanVeto</td>
<td>✓</td>
<td>✓</td>
<td>0.00021</td>
</tr>
<tr>
<td>2009-11-20</td>
<td>112.6</td>
<td>0.288</td>
<td>UltraCleanVeto</td>
<td>X</td>
<td>X</td>
<td>0.00020</td>
</tr>
<tr>
<td>2008-12-24</td>
<td>226.9</td>
<td>0.069</td>
<td>UltraClean</td>
<td>X</td>
<td>X</td>
<td>0.00128</td>
</tr>
<tr>
<td>2009-12-20</td>
<td>467.7</td>
<td>0.338</td>
<td>UltraCleanVeto</td>
<td>X</td>
<td>X</td>
<td>0.00208</td>
</tr>
</tbody>
</table>

TABLE I. Event Information for P8R2_SOURCE_V2 events with recorded energies exceeding 100 GeV observed within 0.5° of the solar center. Checkmarks indicate events that were recorded as photons in previous Pass 6 and Pass 7 analyses, while “BG Prob” indicates the probability that diffuse emission produced the event. Events below the double-line did not pass our default selection criteria, as they were observed when the Sun was located within 5° of the Galactic plane. (Additional Event Information Provided in Appendix J).

above 100 GeV over the solar path produces ~0.3 background events over the full analysis period (see Appendix C).

Examining each event yields three insights. First, we observe several extremely high-energy events, including three events exceeding 200 GeV, and one event at 470 GeV. The average Fermi-LAT energy dispersion near 100 GeV is ~7%, indicating that these events have true energies well above 100 GeV. This suggests that multi-TeV protons can produce outgoing γ-rays through solar interactions, and that HAWC observations of the upcoming solar minimum may be illuminating. In Figure 5 of Ref. [18], we carefully compare the solar minimum flux against projected HAWC sensitivities.

Second, all six high-energy events were observed between November 2008 and November 2009, which is inconsistent with a steady-state hypothesis. We determine the significance of this variability by conducting a Kolmogorov-Smirnov test of the hypothesis that the data is Poissonian in solar exposure. We rule out the steady-state hypothesis at 4.2σ. Noting that the Sun moves through the Galactic plane during solar minimum and that the diffuse contribution along the plane is negligible above 100 GeV, we unblind the region b < 5°. This adds two new events above 100 GeV that were also observed at solar minimum and no additional events during the subsequent eight passages of the Galactic plane. Including these events increases the statistical significance for variability to 4.9σ.

Third, we note a peculiar “double-event” on December 13, 2008, when two >100 GeV γ-rays were observed within 3.5 hr. The probability that any two events are this closely correlated is inconsistent with Poissonian expectations at ~2.9σ.

Interpretation.—We have provided several lines of evidence that reveal two distinct γ-ray emission components on the solar disk. The first emits primarily from the Sun’s polar regions, has a constant amplitude over the solar cycle, and produces no observed flux above 100 GeV. The second emits primarily from the Sun’s equatorial plane, has an amplitude that decreases drastically after solar minimum, and has a hard spectrum at solar minimum that extends above 200 GeV.

These results are not explained by the SSG model. The bright γ-rays flux across the solar surface does support the SSG mechanism of cosmic-ray reversal deep within the photosphere. However, the flux, spectrum, time-variation, morphological shift, and spectral dip of solar γ-rays are unexplained. We can qualitatively parameterize the solar γ-ray flux as:

\[ \Phi_{\odot}(E) = \frac{R^2_{\odot}}{4D^2} I_{CR}(E_{CR}) C(E_{\gamma}, E_{CR}) f_{\text{surf}} f_{\text{turn}} f_{\text{int}} \]  

where \( \Phi_{\odot} \) is the disk γ-ray flux, \( I_{CR} \) is the cosmic-ray flux at the solar surface, \( C \) describes the γ-ray flux at energy \( E_{\gamma} \) produced by a hadronic interaction at energy \( E_{CR} \) (see Appendix E), \( f_{\text{surf}} \) is the fraction of the solar surface that produces γ-rays, \( f_{\text{turn}} \) is the fraction of incoming cosmic rays that are reversed by magnetic fields within the solar photosphere, and \( f_{\text{int}} \) is the fraction of these cosmic rays that undergo a hadronic interaction and produce outgoing γ-rays before leaving the surface. SSG found solar modulation to be negligible, implying that \( I_{CR} \) is similar to the interstellar cosmic-ray flux. SSG assumes that each efficiency term is energy, position, and time independent. In particular, SSG set \( f_{\text{surf}} \) and \( f_{\text{turn}} \) to unity, and calculated \( f_{\text{int}} \sim 0.5% \).

Our observations instead indicate that these parameters strongly depend on the cosmic-ray energy, solar cycle, and solar latitude. At solar minimum, these shifts are remarkable for four reasons. The large flux, within a factor of ~4 of the maximal value, implies that all efficiency parameters are near unity. The hard spectrum, significantly exceeding the \( E^{-2.7} \) interstellar cosmic-ray spectrum, indicates that these efficiencies rise quickly with energy. The equatorial morphology indicates that polar regions are not emitting efficiently, implying \( f_{\text{surf}} \lesssim 0.5 \). Finally, symmetry constrains \( f_{\text{int}} \sim 0.5 \), as cosmic rays should undergo equal interactions while entering and exiting the photosphere. These considerations produce significant tension with any SSG-like model.

This tension motivates us to consider scenarios that violate the assumptions of Eq. (1) and allow for larger γ-ray fluxes. In Appendix K, we discuss the potential for effects such as magnetic focusing, cosmic-ray trapping, or anisotropic emission to boost the observed γ-ray flux. However, we find that each possibility is either observationally excluded or theoretically unmotivated.

One potential insight stems from the two >100 GeV γ-rays observed on December 13, 2008, which may be connected to a contemporaneous Earth-bound coronal mass ejection (CME) that began on December 12, 2008 and encountered Earth on December 17 [33–35]. However, none of the remaining >100 GeV γ-rays correspond to significant CMEs. In Appendix K we discuss the potential for CMEs, helmet streamers, and coronal holes to explain the peculiar morphology and temporal variability of >100 GeV events.

A New Event—While finalizing this letter, we found a new >100 GeV event. Observed on February 13, 2018 at 17:49:15 UTC, it has an energy of 162 GeV, is located 0.36° from the solar center, passes the UltraCleanVeto event selection, and belongs to the PSF0 and EDISP3 event classes. As we re-enter solar minimum, this is the first >100 GeV event recorded within 0.5° of the sun since 2009. It may be
connected to a Earth-bound CME observed on February 12, 2018. Preliminary work indicates that this event increases the significance of the >100 GeV time variability above 5σ, and supports the hypothesis that the upcoming solar minimum will provide a substantial flux of high-energy events.

**Future Outlook.**—We have discovered statistically significant temporal variations in the intensity, spectrum and morphology of solar γ-ray emission. These variations strongly suggest that two distinct components substantially contribute to the total solar γ-ray flux, including (1) a polar component that varies moderately in time and has a γ-ray spectrum that falls sharply around 100 GeV, and (2) an equatorial component with an extremely hard γ-ray spectrum that continues above 200 GeV, but is dominant only during solar minimum. These observations provide important new clues about the mechanisms behind solar disk γ-ray emission, which remains mysterious.

This mystery is deepened by the high intensity and hard spectrum of disk emission. In particular, the solar minimum flux appears to be in tension with the most optimistic predictions from the class of models that convert the interstellar cosmic-ray flux into a time-invariant and isotropic γ-ray flux. If future observations detect emission at even moderately higher energies, a new mechanism will be necessary to explain the highest-energy solar emission.

Observations of the upcoming Cycle 25 solar minimum by both the Fermi-LAT and HAWC will provide valuable information. Preliminary estimates indicate that the Cycle 25 minimum will be even quieter than the Cycle 24 minimum [36]. Moreover, the Fermi-LAT satellite was operational only during the second half of the 2007–2009 cycle 24 minimum [16]. The observation of >100 GeV γ-rays during this period will significantly enhance our understanding of high-energy solar disk γ-ray emission. One new event has recently been detected. With improved statistics, it will soon become possible to correlate γ-rays with solar observables, shining light on the processes responsible for the high-energy γ-ray flux.

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3 For additional information regarding the analyses and interpretation of these results, please see the supplemental material, which includes references [37–49].