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Visualizing Mie resonances in low-index dielectric nanoparticles

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ABSTRACT: Resonant light scattering by metallic and high-index dielectric 13 nanoparticles has received enormous attention and found many great applications. 14 15 However, low-index dielectric nanoparticles typically do not show resonant scattering behaviors due to poor light confinement caused by small index contrast. This letter 16 17 describes a simple and effective approach to drastically enhance the resonance effect of the low-index particles by partial metal dressing. Mie resonances of low-index 18 19 nanoparticles can now be easily visualized by scattered light. This scattering peak 20 depends on sphere size and has a reasonable line width. A size difference as small as 21 8 nm was resolved by peak shift or even by color change. The scattering peak is attributed to the enhanced TE_{11} Mie resonance of the low-index nanospheres. The 22 23 metal dress not only provides high-reflection boundary, but also functions as an 24 antenna to couple the confined light power to the far field, leading to scattering 25 maxima in the spectra. Additionally, the enhanced TE_{11} Mie resonance in low-index 26 nanoparticles features considerable magnetic response due to the strong circulating displacement currents induced by the intensified *E*-field instead of a high permittivity 27 for index particles. The enhanced Mie resonances could be used to sense minute 28 changes in size or refractive index of low-index nanoparticles and benefit a wide 29 30 range of applications.

31 TEXT:

32 Light scattering by small particles is a fundamental topic in electromagnetics and 33 attracts much attention throughout modern history [1]. Among all the phenomena, resonant scattering is the most outstanding behavior because it can be easily detected 34 by spectroscopy and features a strong local field, enabling various applications of 35 nanoparticles [2-8]. Metallic nanoparticles with localized surface plasmon resonances 36 have been extensively studied and practically utilized in surface enhanced 37 spectroscopy and metamaterials [9-11]. Recently, high-index dielectric nanoparticles 38 with strong Mie resonances drew considerable interest. Efficient field localization and 39 40 negligible material absorption facilitate prominent Mie resonances, leading to optical magnetic response and scattering manipulation [7,8,12-16]. For low-index dielectric 41 nanoparticles, the Mie resonances are weak and broad due to poor light confinement 42 caused by low reflection at particle boundaries, so that they are typically not visible in 43 light scattering spectra. However, low-index dielectric nanoparticles, such as SiO₂ and 44 polystyrene nanoparticles, have more applications than their high-index or even 45 metallic counterparts. If pronounced Mie resonances can be realized, many potential 46 47 applications can be envisioned. In this work, we introduce a metal dressing approach to enhance Mie resonances of low-index nanoparticles so that they can be easily 48 visualized by scattered light. The metal dress improves light confinement ability of 49 the low-index nanoparticles by enhancing boundary reflection. In addition, the open 50 configuration of the metal dress turns dark modes bright. We experimentally and 51 theoretically investigated the metal-dressed SiO₂ nanospheres. While strong 52 plasmonic resonance associated with the crescent shape structures [17,18] occur in 53 sub-100nm size particles, they become very weak in the NIR range for large particles. 54 In this work, we found significantly enhanced TE_{11} Mie mode in the visible range. 55 Low order Mie resonances scales almost linearly with the particle size, allowing us to 56 distinguish an 8 nm size difference. At this resonance, the strong circulating 57 displacement currents induced by the intensified *E*-field rather than high permittivity 58 could result in considerable magnetic response. The enhanced Mie resonances could 59 be used to sense minute changes in size or refractive index and thus benefit a lot of 60

61 applications utilizing the low-index nanoparticles.

62 The metal-dressed SiO₂ nanospheres were created by a two-step metal deposition process by electron beam evaporation. First, a thin metal film was deposited on a 63 transparent glass substrate. Then, a dilute solution of nanospheres was spin-casted on 64 the metal film to create well-separated particle distribution. Finally, another thin metal 65 film was deposited on top of the particle and also on the first metal layer. With this 66 process, each of the SiO₂ nanospheres is partially covered by a metal cap and located 67 in a shallow well in the metal film (Fig. 1(a)). In our case, the substrate is fused silica 68 and the deposited metal films are both Au, each with 30 nm thickness. According to 69 70 SEM inspection, the sizes of the SiO_2 nanospheres range from 400 nm to 500 nm. In the microscopic dark field transmission view, the metal-dressed nanospheres exhibit 71 72 different colors (Fig. 1(b)). The NA of the dark field condenser is 0.80-0.95 and that of the objective is 0.75. In comparison, bare SiO_2 nanospheres could be identified (Fig. 73 1(c)), but shows no distinct color difference despite slightly different sizes because of 74 their flat scattering spectra. At middle stage of the two-step metal deposition process, 75 where bare SiO_2 nanospheres are located on a thin Au film, there is also no color 76 77 difference over individual nanospheres (Fig. 1(d)). The greenish color comes from the transmission through Au film that allows a higher transmission of green light. 78



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Fig. 1 (a) SEM image of a metal-dressed SiO₂ nanosphere. (b)-(d) Microscopic dark field transmission views of metal-dressed SiO₂ nanospheres, SiO₂ nanospheres on a fused silica substrate and SiO₂ nanospheres on a 30 nm Au film over a fused silica substrate. (e) Measured forward scattering spectrum of a 486 nm nanosphere (black) together with simulation (red), and that at a blank area without nanospheres (blue). The simulated spectrum is normalized to the incident power. The incident *E*-field is along the *x*-axis.

87 With a multimode fiber (core diameter ~200 nm), the dark field transmission of a single nanosphere could be selectively collected and analyzed by a spectrometer. 88 89 Since the background transmission is significantly suppressed, the local property of 90 any single metal-dressed nanosphere could be revealed by the spectrum. For example, 91 the dark field transmission at a nanopshere with a diameter around 486 nm (size 92 determined by SEM) is plotted in Fig. 1(e) (black curve). The peak around 510 nm is 93 due to the interband transition in Au. This can be confirmed by analyzing the stray 94 light through a blank area without any nanospheres but only the 60 nm Au film; the

peak around 510 nm in the spectrum (blue curve in Fig. 1(e)) is attributed to the 95 96 material property of Au film only. The peak around 630 nm is a resonance of the metal-dressed SiO₂ nanosphere and responsible for the distinctive colors. Simulations 97 by finite element method were also performed for comparison. The forward scattering 98 was simulated by integrating the transmitted power over a solid angle corresponding 99 100 to the NA of the objective lens. The simulated spectrum (dashed red curve in Fig. 1(e)) agrees with the measurement (black curve) in terms of peak positions and profiles. 101 The simulations were performed in 3D and the geometry followed the real sample. An 102 oblique incidence with an angle 57° was used to mimic the dark field illumination. 103

104 Five nanospheres with size difference from 20 nm down to single digit nm are selected for spectral characterization (Fig. 2(a)). The diameter of each nanosphere is 105 106 obtained through SEM inspection of the same particle at a high magnification. In order to reduce the influence of the stray-light transmission peak from the thin Au 107 108 film itself, the spectrum measured at a blank area is subtracted as a background from that measured at each nanosphere (Fig. 2(b)). For example, the curve for nanosphere 109 No. 5 is obtained by subtracting the blue curve from the black curve in Fig. 1(e), thus 110 111 revealing the main peak associated with the resonance of the metal-dressed SiO₂ 112 nanosphere. With the sphere size increasing from 421 nm to 486 nm, this peak is 113 red-shifted accordingly as confirmed by simulations (red dashed curves). It is worth 114 noting that a size difference as small as 8 nm between nanosphere No. 1 and 2 could 115 be clearly identified by the peak shift in spectrum, while they appear the same without the metal dressing. Since the Mie resonance is within the visible range, the peak shift 116 117 due to size variation could also be distinguished by the color difference. As this resonance is red-shifted with the increasing sphere size, the color of the nanospheres 118 119 in the dark field transmission view varies from green to orange (Fig. 2(c)). By 120 calibrating colors with spectra, sizes of nanospheres could be read out directly through their colors with some accuracy [19]. In comparison, the SEM image (Fig. 121 2(d)) with the same field of view and the same magnification of the optical 122 microscope image (Fig. 2(c)) fails to tell the size difference; only zoom-in to 123 individual particles can tell its exact size. In addition, the enhanced Mie resonance is 124

- robust against small variations in the metal dressing. A 5 nm deviation in the thickness
- of each metal deposition step does not affect the resonant frequency [19].



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Fig. 2 (a) SEM images of five selected metal-dressed SiO_2 nanospheres with different sizes. (b) Background subtracted forward scattering spectra (black lines) of the five nanospheres and the corresponding simulations (dashed red lines). (c) Photopicture in the dark field transmission view for nanospheres including the five selected ones with size dependent colors. (d) SEM image of the same nanospheres in (c) with the same scope and the same magnification.

134 The pronounced peak in the forward scattering spectra is revealed as an enhanced Mie resonance by comparing the metal-dressed SiO₂ nanospheres with 135 complete-metal-shell ones. Fig. 3(a) presents the scattering efficiency $Q_{\rm sc}$ of a 136 complete-Au-shell SiO₂ nanosphere in air as a function of wavelength and core size 137 theory [1,24]. The shell thickness 30 138 based on Mie remains nm.

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$$Q_{sc} = \frac{2}{(k_0 r)^2} \sum_{n=1}^{\infty} (2n+1) (|a_n|^2 + |b_n|^2)$$
, where k_0 is free space wave vector; r is the core

radius; a_n and b_n are coefficients of the scattered field decomposed in the basis of 140 141 vector spherical harmonics [1,24]. Due to the tight light confinement provided by the 142 metal shell, the Mie resonances of the SiO_2 nanosphere are significantly enhanced but become less radiant at the same time. Through interaction, the Mie resonances show 143 144 up as Fano dips (cutting lines in Fig. 3(a)) in the background of the metal sphere scattering [25,26], leading to resonant scattering suppression. Fig. 3 (b) shows Q_{sc} at 145 the core size of 550 nm and the decomposed contributions from the first three 146 modes: $Q_{a1} = 6|a_1|^2/(k_0 r)^2$, $Q_{b1} = 6|b_1|^2/(k_0 r)^2$ Mie 147 fundamental and $Q_{a2} = 10 |a_2|^2 / (k_0 r)^2$. a_1 , b_1 and a_2 correspond to the TM₁₁, TE₁₁ and TM₁₂ mode, 148 respectively. TM/TE stands for transvers magnetic/electric modes with no radial 149 150 magnetic/electric field. The subscripts denote different orders. For a better 151 comparison with our experiment, the complete-Au-shell nanosphere is located on a 152 Au film and illuminated under dark field scheme (Fig. 3(c) inset). The Au film has a 30 nm indentation below the nanosphere and is 60 nm thick elsewhere. The shell 153 154 merges with the Au plane at the bottom. As shown in Fig. 3(c), the Mie resonances are 155 observed in the simulated forward scattering spectrum at the same wavelengths as 156 those in Fig. 3(a) and (b). In this case, the three resonances are still less radiant and 157 thus show Fano shapes in the background of the transmission through the Au film. 158 The light field distribution at the TE_{11} resonance (Fig. 3(d)) reveals the mode pattern as a an inclining magnetic dipole (MD) with electric field vortex circulating around it, 159 160 as diagrammatically sketched by the inset of Fig. 3(c).



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Fig. 3 (a) Scattering efficiency (Q_{sca}) spectrum at different core sizes of the 162 complete-shell nanosphere (shell thickness 30 nm) in air. (b) Q_{sca} spectrum at the core 163 size of 550 nm (indicated by the horizontal dashed line in (a)) and the decomposed 164 165 contributions from the first three fundamental Mie modes (TM_{11} , TE_{11} and TM_{12}). (c) Forward scattering spectrum of the complete-shell nanosphere (core size 550 nm) 166 167 located on a Au film under dark field illumination. (d) Distributions of scattered electric field x-component $|E_{sc-x}|$ and scattered in-plane magnetic field vector (H_{sc-y} , 168 169 H_{sc-z}) at the TE₁₁ resonance on the y-z plane at x=0. (e) and (f) Forward scattering 170 spectrum and field distribution of the metal-dressed nanosphere (486 nm) at the TE_{11} 171 resonance.

Fig. 3(e) presents the simulated forward scattering spectrum of a metal-dressed SiO_2 172 173 nanosphere with a diameter around 486 nm. This spectrum is the same as the dashed 174 red line in Fig. 1 (e) but extended to longer wavelengths. Fig. 3(f) presents the field 175 distribution at the resonance (~ 630 nm). A direct comparison between Fig. 3(d) and (f) reveals that the pronounced peak we observed in experiment corresponds to the 176 177 enhanced TE_{11} Mie resonance of the SiO₂ nanosphere. In another model, the enhanced 178 TE_{11} Mie mode could be considered as a Fabry-Perot (F-P) resonance of a waveguide 179 mode propagating at a tilted angle (Fig. 3 (f)). Compared to the mode pattern of the metal-sandwiched cylindrical waveguide, the light field inside the metal-dressed 180 181 sphere could be regarded as a standing wave of the waveguide mode due to multiple 182 reflections [19]. The waveguide mode F-P resonance in our configuration is a localized mode confined in the metal-dressed nanoparticle, and therefore insensitive 183

to the incident angle. Although a complete metal shell provides a better light 184 185 confinement than the metal dress as revealed by the higher field concentration inside 186 the SiO_2 nanosphere (comparison between Fig. 3(d) and (e)), the metal dress is much more favorable to fabrication. Core-shell structures below 100 nm could be created by 187 chemical synthesis [27]. However, making larger shells (several hundred nm to 1 μ m) 188 189 is much tougher than making a metal dress by a two-step deposition. Further, the Mie 190 resonances in the metal-dressed nanosphere appear as peaks in scattering spectra (Fig. 191 1(e) and Fig. 3(e)), while the Mie resonances in the complete-shell nanosphere appear 192 as Fano dips (Fig. 3(a)-(c)). The metal dress not only enhances the Mie resonances by 193 providing high reflecting boundaries but also couples the Mie mode into the far field 194 by serving as an antenna.

195 The enhanced TE₁₁ Mie resonance is associated with a MD resonance since the 196 field distribution (Fig. 3(f)) shows circulating displacement currents and an intensified 197 magnetic field at the center. Actually, the ideal TE_{11} Mie resonance of a nanosphere could just be regarded as a MD resonance [1]. The magnetic response of high-index 198 199 dielectric nanoparticles is mainly due to that and have attracted much attention [7,8,12]. Since the circulating displacement current density $(J_D = (-i\omega)\varepsilon_0\varepsilon_r E)$ as well 200 201 as the induced MD increases with permittivity, prominent magnetic responses are 202 observed. From another point of view, one need not to rely solely on the high 203 refractive index to produce strong magnetic resonances if a strong E-field can be 204 obtained [12]. Typically intensified *E*-field can be obtained in an optical cavity with 205 high Q-factor modes. Since low-index particles provide very weak light confinement 206 and therefore low-Q Mie modes, they are typically are not suited for inducing 207 magnetic response. However, with help of the metal dressing, the Mie modes of 208 low-index nanoparticles are greatly enhanced. At the TE_{11} mode (Fig. 3(f)), the strong 209 circulating electric field leads to strong circulating displacement currents and thus a 210 strong MD. This result implies that considerable magnetic response could also be induced in low-index nanoparticles by a simple metal dressing process. Nevertheless, 211 due to the induced currents in the metal dress, the optical response at the TE₁₁ Mie 212

resonance is not a pure MD response. Other components could be identified by a multipole decomposition calculation based on the simulated light field distributions [19]. In the view of method of images, the tilted MD on a Au film could be considered as the MD with its image and the appearance of high order multipoles is reasonable [19].

218 The enhanced Mie resonance is distinguished from plasmon modes. The Mie mode is concentrated within the SiO_2 nanosphere as an optical cavity while plasmon 219 modes are typically at metal/dielectric interfaces [19]. Moreover, the Mie mode is also 220 different from split ring like resonances, associated with highly concentrated field 221 222 within the split gaps [28], although the metal dress shares some similarities with a 223 split shell. The metal dressing process has also been applied to semiconductor 224 nanorods with high refractive index for plasmonic cloaking [29,30] and nanoscale 225 color filtering [31]. For deep-subwavelength sizes, when the induced electric dipole 226 moment in the core and that in the shell are equal in magnitude but opposite in sign, scattering cancellation is realized. For spheres with larger sizes in our case, phase 227 228 retardation plays a role and Mie modes become the key factor of scattering properties. 229 The enhanced Mie resonances as local modes could be utilized for sensing the 230 individual low-index nanospheres. Low-index nanoparticles such as SiO₂ and polystyrene nanospheres have a wide range of important applications [32-38]. Sensing 231 232 the local properties of individual nanoparticles, such as size and refractive index, in a 233 proper manner would greatly benefit all those applications. The enhanced Mie 234 resonances as peaks in scattering spectra could work for sensing. Since the resonance 235 occurs in the visible range, the minute size difference could also be sensed by color 236 difference. After mapping the particle size with color, we could obtain the sizes of all 237 the nanospheres in one micrograph by analyzing their colors. Compared to SEM 238 inspection and dynamic light scattering, this colorimetric approach could lead to a 239 new particle sizing method with both efficiency and accuracy [19]. When particle size 240 is fixed, a minute change in refractive index of the particle could be sensed by shift of the enhanced Mie resonance [19]. 241

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In summary, the metal dress created by a two-step deposition process

prominently enhances the Mie resonances of low-index nanoparticles. The metal dress 243 244 improves light confinement ability of low-index nanoparticles by enhancing boundary 245 reflection and thus enhances the Mie resonances. Beyond a high reflection cavity, the metal dress also functions as a broadband antenna to couple the concentrated light 246 247 power to the far field. The resonance occurring in the visible range leads to size dependent colors. At the enhanced TE_{11} Mie resonance, the intensified *E*-field instead 248 249 of high permittivity induces strong circulating displacement currents and thus 250 considerable magnetic response. The enhanced Mie resonances could be used to sense 251 minute changes in size or refractive index of low-index nanoparticles and benefit 252 many applications.

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