



# CHORUS

This is the accepted manuscript made available via CHORUS. The article has been published as:

## First Results from CUORE: A Search for Lepton Number Violation via $0\nu\beta\beta$ Decay of $^{130}\text{Te}$

C. Alduino *et al.* (CUORE Collaboration)

Phys. Rev. Lett. **120**, 132501 — Published 26 March 2018

DOI: [10.1103/PhysRevLett.120.132501](https://doi.org/10.1103/PhysRevLett.120.132501)

**First Results from CUORE:**  
**A Search for Lepton Number Violation via  $0\nu\beta\beta$  Decay of  $^{130}\text{Te}$**

C. Alduino,<sup>1</sup> F. Alessandria,<sup>2</sup> K. Alfonso,<sup>3</sup> E. Andreotti,<sup>4,5, a</sup> C. Arnaboldi,<sup>6</sup> F. T. Avignone III,<sup>1</sup> O. Azzolini,<sup>7</sup> M. Balata,<sup>8</sup> I. Bandac,<sup>1</sup> T. I. Banks,<sup>9,10</sup> G. Bari,<sup>11</sup> M. Barucci,<sup>12,13, b</sup> J.W. Beeman,<sup>14</sup> F. Bellini,<sup>15,16</sup> G. Benato,<sup>9</sup> A. Bersani,<sup>17</sup> D. Biare,<sup>10</sup> M. Biassoni,<sup>5</sup> F. Bragazzi,<sup>17</sup> A. Branca,<sup>18</sup> C. Brofferio,<sup>6,5</sup> A. Bryant,<sup>10,9, c</sup> A. Buccheri,<sup>16</sup> C. Bucci,<sup>8</sup> C. Bulfon,<sup>16</sup> A. Camacho,<sup>7</sup> A. Caminata,<sup>17</sup> L. Canonica,<sup>19,8</sup> X. G. Cao,<sup>20</sup> S. Capelli,<sup>6,5</sup> M. Capodiferro,<sup>16</sup> L. Cappelli,<sup>9,10,8</sup> L. Cardani,<sup>16</sup> M. Cariello,<sup>17</sup> P. Carniti,<sup>6,5</sup> M. Carrettoni,<sup>6,5</sup> N. Casali,<sup>16</sup> L. Cassina,<sup>6,5</sup> R. Cereseto,<sup>17</sup> G. Ceruti,<sup>5</sup> A. Chiarini,<sup>11</sup> D. Chiesa,<sup>6,5</sup> N. Chott,<sup>1</sup> M. Clemenza,<sup>6,5</sup> D. Conventi,<sup>7</sup> S. Copello,<sup>21,17</sup> C. Cosmelli,<sup>15,16</sup> O. Cremonesi,<sup>5, d</sup> C. Crescentini,<sup>11</sup> R. J. Creswick,<sup>1</sup> J. S. Cushman,<sup>22</sup> A. D'Addabbo,<sup>8</sup> D. D'Aguzzo,<sup>8,23</sup> I. Dafinei,<sup>16</sup> V. Datskov,<sup>5</sup> C. J. Davis,<sup>22</sup> F. Del Corso,<sup>11</sup> S. Dell'Oro,<sup>24,8,25</sup> M. M. Deninno,<sup>11</sup> S. Di Domizio,<sup>21,17</sup> M. L. Di Vacri,<sup>8,26</sup> L. Di Paolo,<sup>10, e</sup> A. Drobizhev,<sup>9,10</sup> L. Ejzak,<sup>27, f</sup> R. Faccini,<sup>15,16</sup> D. Q. Fang,<sup>20</sup> M. Faverzani,<sup>6,5</sup> E. Ferri,<sup>5</sup> F. Ferroni,<sup>15,16</sup> E. Fiorini,<sup>5,6</sup> M. A. Franceschi,<sup>28</sup> S. J. Freedman,<sup>10,9, g</sup> B. K. Fujikawa,<sup>10</sup> R. Gaigher,<sup>5</sup> A. Giachero,<sup>6,5</sup> L. Gironi,<sup>6,5</sup> A. Giuliani,<sup>29</sup> L. Gladstone,<sup>19</sup> J. Goett,<sup>8, h</sup> P. Gorla,<sup>8</sup> C. Gotti,<sup>6,5</sup> C. Guandalini,<sup>11</sup> M. Guerzoni,<sup>11</sup> T. D. Gutierrez,<sup>30</sup> E. E. Haller,<sup>14,31</sup> K. Han,<sup>32</sup> E. V. Hansen,<sup>19,3, i</sup> K. M. Heeger,<sup>22</sup> R. Hennings-Yeomans,<sup>9,10</sup> K. P. Hickerson,<sup>3</sup> H. Z. Huang,<sup>3</sup> M. Iannone,<sup>16</sup> L. Ioannucci,<sup>8</sup> R. Kadel,<sup>33</sup> G. Keppel,<sup>7</sup> L. Kogler,<sup>10,9</sup> Yu. G. Kolomensky,<sup>9,10</sup> A. Leder,<sup>19</sup> C. Ligi,<sup>28</sup> K. E. Lim,<sup>22</sup> X. Liu,<sup>3</sup> Y. G. Ma,<sup>20</sup> C. Maiano,<sup>6,5, j</sup> M. Maino,<sup>6,5</sup> L. Marini,<sup>21,17</sup> M. Martinez,<sup>15,16,34</sup> C. Martinez Amaya,<sup>1</sup> R. H. Maruyama,<sup>22</sup> Y. Mei,<sup>10</sup> N. Moggi,<sup>35,11</sup> S. Morganti,<sup>16</sup> P. J. Mosteiro,<sup>16</sup> S. S. Nagorny,<sup>8,25</sup> T. Napolitano,<sup>28</sup> M. Nastasi,<sup>6,5</sup> S. Nisi,<sup>8</sup> C. Nones,<sup>36</sup> E. B. Norman,<sup>37,38</sup> V. Novati,<sup>29</sup> A. Nucciotti,<sup>6,5</sup> I. Nutini,<sup>8,25</sup> T. O'Donnell,<sup>24</sup> M. Olcese,<sup>17, k</sup> E. Olivieri,<sup>12,13, l</sup> F. Orio,<sup>16</sup> D. Orlandi,<sup>8</sup> J. L. Ouellet,<sup>19</sup> C. E. Pagliarone,<sup>8,23</sup> M. Pallavicini,<sup>21,17</sup> V. Palmieri,<sup>7</sup> L. Pattavina,<sup>8</sup> M. Pavan,<sup>6,5</sup> M. Pedretti,<sup>37</sup> R. Pedrotta,<sup>18</sup> A. Pelosi,<sup>16</sup> G. Pessina,<sup>5</sup> V. Pettinacci,<sup>16</sup> G. Piperno,<sup>15,16, m</sup> C. Pira,<sup>7</sup> S. Pirro,<sup>8</sup> S. Pozzi,<sup>6,5</sup> E. Previtali,<sup>5</sup> F. Reindl,<sup>16</sup> F. Rimondi,<sup>35,11, g</sup> L. Risegari,<sup>12,13, n</sup> C. Rosenfeld,<sup>1</sup> C. Rossi,<sup>17</sup> C. Rusconi,<sup>1,8</sup> M. Sakai,<sup>3</sup> E. Sala,<sup>6,5, o</sup> C. Salvioni,<sup>4,5</sup> S. Sangiorgio,<sup>37</sup> D. Santone,<sup>8,26</sup> D. Schaeffer,<sup>6,5, p</sup> B. Schmidt,<sup>10</sup> J. Schmidt,<sup>3</sup> N. D. Scielzo,<sup>37</sup> V. Singh,<sup>9</sup> M. Sisti,<sup>6,5</sup> A. R. Smith,<sup>10</sup> F. Stivanello,<sup>7</sup> L. Taffarello,<sup>18</sup> L. Tatananni,<sup>8</sup> M. Tenconi,<sup>29</sup> F. Terranova,<sup>6,5</sup> M. Tessaro,<sup>18</sup> C. Tomei,<sup>16</sup> G. Ventura,<sup>12,13, q</sup> M. Vignati,<sup>16</sup> S. L. Wagaarachchi,<sup>9,10</sup> J. Wallig,<sup>39</sup> B. S. Wang,<sup>37,38</sup> H. W. Wang,<sup>20</sup> B. Welliver,<sup>10</sup> J. Wilson,<sup>1</sup> K. Wilson,<sup>1</sup> L. A. Winslow,<sup>19</sup> T. Wise,<sup>22,27</sup> L. Zanotti,<sup>6,5</sup> C. Zarra,<sup>8</sup> G. Q. Zhang,<sup>20</sup> B. X. Zhu,<sup>3, r</sup> S. Zimmermann,<sup>39</sup> and S. Zucchelli<sup>35,11</sup>

(CUORE Collaboration)

<sup>1</sup>*Department of Physics and Astronomy, University of South Carolina, Columbia, SC 29208, USA*

<sup>2</sup>*INFN – Sezione di Milano, Milano I-20133, Italy*

<sup>3</sup>*Department of Physics and Astronomy, University of California, Los Angeles, CA 90095, USA*

<sup>4</sup>*Dipartimento di Fisica e Matematica, Università dell'Insubria, Como I-22100, Italy*

<sup>5</sup>*INFN – Sezione di Milano Bicocca, Milano I-20126, Italy*

<sup>6</sup>*Dipartimento di Fisica, Università di Milano-Bicocca, Milano I-20126, Italy*

<sup>7</sup>*INFN – Laboratori Nazionali di Legnaro, Legnaro (Padova) I-35020, Italy*

<sup>8</sup>*INFN – Laboratori Nazionali del Gran Sasso, Assergi (L'Aquila) I-67100, Italy*

<sup>9</sup>*Department of Physics, University of California, Berkeley, CA 94720, USA*

<sup>10</sup>*Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

<sup>11</sup>*INFN – Sezione di Bologna, Bologna I-40127, Italy*

<sup>12</sup>*Dipartimento di Fisica, Università di Firenze, Firenze I-50125, Italy*

<sup>13</sup>*INFN – Sezione di Firenze, Firenze I-50125, Italy*

<sup>14</sup>*Materials Science Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

<sup>15</sup>*Dipartimento di Fisica, Sapienza Università di Roma, Roma I-00185, Italy*

<sup>16</sup>*INFN – Sezione di Roma, Roma I-00185, Italy*

<sup>17</sup>*INFN – Sezione di Genova, Genova I-16146, Italy*

<sup>18</sup>*INFN – Sezione di Padova, Padova I-35131, Italy*

<sup>19</sup>*Massachusetts Institute of Technology, Cambridge, MA 02139, USA*

<sup>20</sup>*Shanghai Institute of Applied Physics, Chinese Academy of Sciences, Shanghai 201800, China*

<sup>21</sup>*Dipartimento di Fisica, Università di Genova, Genova I-16146, Italy*

<sup>22</sup>*Wright Laboratory, Department of Physics, Yale University, New Haven, CT 06520, USA*

<sup>23</sup>*Dipartimento di Ingegneria Civile e Meccanica,*

*Università degli Studi di Cassino e del Lazio Meridionale, Cassino I-03043, Italy*

<sup>24</sup>*Center for Neutrino Physics, Virginia Polytechnic Institute and State University, Blacksburg, Virginia 24061, USA*

<sup>25</sup>*INFN – Gran Sasso Science Institute, L'Aquila I-67100, Italy*

<sup>26</sup> *Dipartimento di Scienze Fisiche e Chimiche, Università dell'Aquila, L'Aquila I-67100, Italy*

<sup>27</sup> *Department of Physics, University of Wisconsin, Madison, WI 53706, USA*

<sup>28</sup> *INFN – Laboratori Nazionali di Frascati, Frascati (Roma) I-00044, Italy*

<sup>29</sup> *CSNSM, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, 91405 Orsay, France*

<sup>30</sup> *Physics Department, California Polytechnic State University, San Luis Obispo, CA 93407, USA*

<sup>31</sup> *Department of Materials Science and Engineering,  
University of California, Berkeley, CA 94720, USA*

<sup>32</sup> *INPAC and School of Physics and Astronomy,*

*Shanghai Jiao Tong University; Shanghai Laboratory for Particle Physics and Cosmology, Shanghai 200240, China*

<sup>33</sup> *Physics Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

<sup>34</sup> *Laboratorio de Física Nuclear y Astroparticulas,  
Universidad de Zaragoza, Zaragoza 50009, Spain*

<sup>35</sup> *Dipartimento di Fisica e Astronomia, Alma Mater Studiorum – Università di Bologna, Bologna I-40127, Italy*

<sup>36</sup> *Service de Physique des Particules, CEA / Saclay, 91191 Gif-sur-Yvette, France*

<sup>37</sup> *Lawrence Livermore National Laboratory, Livermore, CA 94550, USA*

<sup>38</sup> *Department of Nuclear Engineering, University of California, Berkeley, CA 94720, USA*

<sup>39</sup> *Engineering Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA*

(Dated: February 23, 2018)

The CUORE experiment, a ton-scale cryogenic bolometer array, recently began operation at the Laboratori Nazionali del Gran Sasso in Italy. The array represents a significant advancement in this technology, and in this work we apply it for the first time to a high-sensitivity search for a lepton-number–violating process:  $^{130}\text{Te}$  neutrinoless double-beta decay. Examining a total  $\text{TeO}_2$  exposure of 86.3 kg·yr, characterized by an effective energy resolution of  $(7.7 \pm 0.5)$  keV FWHM and a background in the region of interest of  $(0.014 \pm 0.002)$  counts/(keV·kg·yr), we find no evidence for neutrinoless double-beta decay. Including systematic uncertainties, we place a lower limit on the decay half-life of  $T_{1/2}^{0\nu}(^{130}\text{Te}) > 1.3 \times 10^{25}$  yr (90% C.L.); the median statistical sensitivity of this search is  $7.0 \times 10^{24}$  yr. Combining this result with those of two earlier experiments, Cuoricino and CUORE-0, we find  $T_{1/2}^{0\nu}(^{130}\text{Te}) > 1.5 \times 10^{25}$  yr (90% C.L.), which is the most stringent limit to date on this decay. Interpreting this result as a limit on the effective Majorana neutrino mass, we find  $m_{\beta\beta} < (110 - 520)$  meV, where the range reflects the nuclear matrix element estimates employed.

The existence of nonzero neutrino masses is well established by precision measurements of neutrino flavor oscillation [1]. This discovery has given renewed impetus to long-standing questions as to the Dirac or Majorana nature of the neutrino [2], the role of Majorana neutrinos in cosmological evolution [3], and the absolute neutrino mass. Neutrinoless double-beta ( $0\nu\beta\beta$ ) decay is a lepton-number–violating process that can occur only if neutrinos are Majorana fermions [4–7]. The discovery of this decay would unambiguously demonstrate that lepton number is not a symmetry of nature and that neutrinos are Majorana particles [8].

If it occurs,  $0\nu\beta\beta$  decay has a robust experimental signature: a peak in the summed energy spectrum of the final state electrons at the  $Q$ -value of the decay ( $Q_{\beta\beta}$ ). To maximize sensitivity to this signature, an experiment must have a low background rate near  $Q_{\beta\beta}$ , good energy resolution, and a large source mass. The Cryogenic Underground Observatory for Rare Events (CUORE) [9] is a new detector that applies the powerful macro-bolometer technique [10, 11] at an unprecedented scale to search for  $0\nu\beta\beta$  decay of tellurium isotopes. In this work, we focus on  $0\nu\beta\beta$  decay of  $^{130}\text{Te}$  to the ground state of  $^{130}\text{Xe}$ . Our sensitivity benefits from the high natural abundance of  $^{130}\text{Te}$ ,  $(34.167 \pm 0.002)\%$  [12], and large  $Q_{\beta\beta}$  of  $(2527.515 \pm 0.013)$  keV [13–15].

CUORE is composed of  $988 \ 5 \times 5 \times 5$  cm<sup>3</sup>  $\text{TeO}_2$  crys-

tals [16], each having a mass of 750 g, which we can cool to temperatures as low as 7 mK. When a crystal absorbs energy, we exploit the resulting temperature increase to measure that energy. Each crystal is instrumented with a thermistor [17] to record thermal pulses, and a heater [18, 19] for thermal gain stabilization.

The crystals are arranged into 19 copper-framed towers, with each tower consisting of 13 floors of 4 crystals. The crystals are held in the tower frame by polytetrafluoroethylene supports. The towers are arranged in a close-packed array and thermally connected to the mixing chamber of a  $^3\text{He}/^4\text{He}$  dilution refrigerator [20], which is precooled by five two-stage ( $\sim 40$  K and  $\sim 4$  K) pulse tube cryocoolers [21] and a Joule-Thomson expansion valve.

To suppress external  $\gamma$ -ray backgrounds, two lead shields are integrated into the cryogenic volume: a 30-cm thick shield at  $\sim 50$  mK above the detectors and a 6-cm thick shield at  $\sim 4$  K around and below the detectors. The lateral and lower shields are made from ancient Roman lead with extremely low levels of radioactivity [22]. An external lead shield (25 cm thick) surrounded by borated polyethylene and boric acid (20 cm thick) provide additional shielding. More details on the experimental subsystems and shielding can be found in Refs. [23–27].

A prototype detector equivalent to a single CUORE tower, CUORE-0, operated at Laboratori Nazionali del

Gran Sasso from 2013 to 2015 and served to validate the materials and low-background assembly techniques used for CUORE [16, 28–31]. Before the current work, the strongest probe of  $\beta\beta$  decay of  $^{130}\text{Te}$  came from CUORE-0 [32–35].

The data presented here are from two month-long datasets collected from May to June (Dataset 1) and August to September (Dataset 2) of 2017. Between the two datasets, we improved the detector operating conditions; in particular, we implemented an active noise cancellation system on the cryocoolers [36] and improved the electrical grounding of the experiment. The detector operating temperature is a compromise between minimizing the heat capacity of the crystals, thus maximizing the thermal gain, and optimizing the signal bandwidth. To select the optimal operating temperature, we performed a temperature scan to study the energy resolution achieved by a representative subset of detectors. An operating temperature of approximately 15 mK was selected for both datasets.

Each dataset is bookended by periods devoted to energy calibration with  $^{232}\text{Th}$   $\gamma$ -ray sources [37]; the closing calibration is performed to verify the stability of the detector response over the dataset. We use the data collected between calibrations, which we refer to as *physics data*, for our  $0\nu\beta\beta$  decay search.

The voltage across each thermistor is amplified and filtered [26, 38–40] and continuously digitized with a sampling rate of 1 kHz [41–43]. A total of 984 of 988 channels are functioning. Thermal event pulses are identified by a software derivative trigger with channel-dependent thresholds ranging from 20 to a few hundred keV; we anticipate reducing these thresholds for future low-energy studies [44, 45]. The rise and fall times of thermal pulses are on the order of 100 ms and 400 ms, respectively. We analyze a 10-s window consisting of 3 s before and 7 s after each trigger. The pre-trigger voltage provides a proxy for the bolometer temperature before the event, while we determine the event energy from the pulse amplitude. The average event rate per detector is  $\sim 50$  mHz in calibration data and  $\sim 6$  mHz in physics data. In addition to triggered pulses, every few minutes each heater is injected with a stable voltage pulse ( $\sim 1$  ppm absolute stability) [46] to generate tagged reference events with fixed thermal energy. To monitor and characterize noise we also analyze waveforms with no discernible thermal pulses.

To improve the signal-to-noise ratio we use an optimal filter [47], which exploits the distinct frequency characteristics of particle-induced and noise waveforms. The pulse amplitude is determined from the maximum value attained by the filtered waveform. To monitor and correct for possible drifts in the energy-to-amplitude response of the detection chain (e.g., due to small drifts in operating temperature), which could otherwise spoil the energy resolution, we apply thermal gain stabiliza-

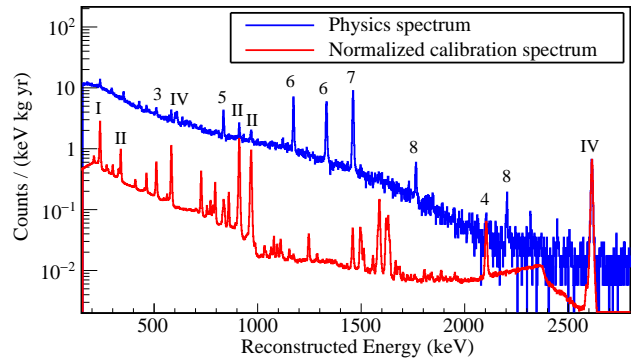


FIG. 1. Reconstructed energy spectra of physics (blue) and calibration (red) data. The calibration spectrum is normalized to the physics data at the 2615-keV line. The sources of the labeled peaks are identified as: (I)  $^{212}\text{Pb}$ , (II)  $^{228}\text{Ac}$ , (3)  $e^+e^-$  annihilation, (4/IV)  $^{208}\text{Tl}$ , (5)  $^{54}\text{Mn}$ , (6)  $^{60}\text{Co}$ , (7)  $^{40}\text{K}$ , (8)  $^{214}\text{Bi}$ . Roman numerals indicate the spectral lines used for calibration.

tion (TGS) to each event amplitude. We apply one of two methods: the first uses monoenergetic heater pulses (heater-TGS), and the second uses pulses induced by  $\gamma$  rays from the 2615-keV  $^{208}\text{Tl}$  calibration line (calibration-TGS). Both methods were developed and used in CUORE-0 [33]. Heater-TGS is our default algorithm, while we use calibration-TGS for the  $\sim 3\%$  of bolometers without functioning heaters and for channels in which calibration-TGS yields a statistically significant improvement in sensitivity compared to heater-TGS. In total, 96.6% of our exposure utilizes heater-TGS while the remainder uses calibration-TGS.

To calibrate the detectors, we use six  $\gamma$  lines from the  $^{232}\text{Th}$  calibration sources ranging from 239 keV to 2615 keV. We estimate the mean stabilized amplitude of each line and create a calibration function for each bolometer in each dataset (each *bolometer-dataset*), which maps stabilized pulse amplitudes to physical energies. We find that the calibration functions of each bolometer-dataset are well described by a second-order polynomial with zero constant term throughout the calibrated energy range. After calibrating, to blind the region near  $Q_{\beta\beta}$ , we take events that reconstruct within 20 keV of the 2615  $^{208}\text{Tl}$  line in physics data and move a blinded fraction of them down by 87 keV; this procedure produces an artificial peak at  $Q_{\beta\beta}$  [33] and is later reversed once the  $0\nu\beta\beta$  search analysis is finalized. The calibration and unblinded physics spectra are shown in Figure 1.

To select  $0\nu\beta\beta$  decay candidates in the physics data, we apply the following selection criteria. Firstly, we discard periods of noisy data caused, for example, by activity in the laboratory. This reduces the exposure by 1%. Next, we impose basic pulse quality requirements

to each event, requiring a single pulse-like feature in the event window and a stable pre-trigger voltage. We then require the shape of each waveform to be consistent with that of a true signal-like event. We build a signal-like event sample in physics data from events that reconstruct within 10 keV of the  $\gamma$  lines from  $^{40}\text{K}$  at 1461 keV and  $^{60}\text{Co}$  at 1173 keV and 1332 keV. We characterize event waveforms with six pulse-shape parameters and represent each event with a point in this 6-dimensional space. We calculate the Mahalanobis distance  $D_M$  [48] for each event from the mean position of the signal sample. We choose the upper limit on  $D_M$  that maximizes the discovery sensitivity [49]. Throughout this optimization, data from the region of interest for  $0\nu\beta\beta$  decay (ROI) are not used. In calculating the figure of merit for a given  $D_M$  cutoff, we estimate the signal selection efficiency from  $^{40}\text{K}$  events near 1461 keV and the background selection efficiency from events with energy between 2700 – 3900 keV. Events in this latter energy range are dominated by partially contained alpha particles and are representative of the dominant background in the ROI. Once the optimal  $D_M$  cutoff is chosen, we evaluate the efficiency of the pulse shape selection using events belonging to the  $^{208}\text{Tl}$  2615-keV line.

To reduce backgrounds from decays depositing energy in multiple crystals (e.g.,  $\alpha$  particles on crystal surfaces or multiple Compton scatters of  $\gamma$  rays), we reject events that occur within 10 ms of an event in a different bolometer (anti-coincidence selection). The width of the coincidence window is chosen after correcting for differences in detector rise times and trigger configurations that can affect the timestamp assigned to an event. The inter-bolometer timestamp differences are determined using physically coincident multi-detector events, such as pair-production events occurring in calibration data. The energy threshold for coincident events in the current analysis is set to 150 keV. The anti-coincidence selection efficiency has two components: the probability for a  $0\nu\beta\beta$  decay to be fully contained in a single crystal and the probability to not accidentally coincide with another event. We estimate the former from simulation [50, 51] and the latter we determine using the 1461-keV  $\gamma$  ray from  $^{40}\text{K}$  electron capture, which is a single-event decay that is not expected to produce physical coincidences.

We evaluate the trigger efficiency as the fraction of tagged heater pulses that produce an event trigger. The heater pulse amplitude is scanned to study the energy dependence of the trigger efficiency. We also exploit heater events to measure the basic pulse quality selection efficiency mentioned above and the energy reconstruction efficiency (i.e., the probability that a monoenergetic pulse reconstructs correctly). The combined trigger, basic pulse quality, and reconstruction efficiency, denoted by *base efficiency*, is averaged over all channels with functioning heaters and applied to all channels. In cases where a step in the event reconstruction procedure

TABLE I. Number of channels studied, event selection efficiencies, and performance parameters for the two datasets analyzed in this work. The effective resolution and background parameters are given at  $Q_{\beta\beta}$ . The uncertainty on exposure is negligible.

	Dataset 1	Dataset 2
<b>Selection Efficiency (%)</b>		
Base	$95.63 \pm 0.01$	$96.69 \pm 0.01$
Pulse shape ( $D_M$ )	$91.1 \pm 3.6$	$98.2 \pm 3.0$
Anti-coincidence (accidental)	$99.4 \pm 0.5$	$100.0 \pm 0.4$
Anti-coincidence ( $\beta\beta$ containment)	$88.35 \pm 0.09$	
Total (excl. $\beta\beta$ containment)	$85.7 \pm 3.4$	$94.0 \pm 2.9$
<b>Performance Parameters</b>		
Channels used	876	935
TeO <sub>2</sub> exposure (kg·yr)	37.6	48.7
Effective resolution (keV)	$8.3 \pm 0.4$	$7.4 \pm 0.7$
Background ( $10^{-2}$ c/(keV·kg·yr))	$1.49 \pm 0.18$	$1.35 \pm 0.19$

fails for a channel, we remove that channel from the subsequent analysis. The selection efficiencies are summarized in Table I.

We establish the detector response to a monoenergetic event near  $Q_{\beta\beta}$  using the high-statistics  $^{208}\text{Tl}$  2615-keV  $\gamma$  line from calibration data. The CUORE detectors exhibit a slightly non-Gaussian line shape, as was observed in CUORE-0 [33] and Cuoricino [52, 53]. The origin of this structure is under investigation; however, we model it empirically with a primary Gaussian component centered at 2615 keV and two additional Gaussian components, one on the right and one on the left of the main peak. We find this model provides a better description of the data compared to other models considered, for example, a single- or double-Gaussian photopeak. The choice of line shape is treated as a systematic uncertainty. The three Gaussian components are parametrized with the same bolometer–dataset dependent width. The normalized line shape function of each bolometer–dataset thus has 6 parameters: the means of the main peak and two subpeaks, the relative intensities of the subpeaks, and the common peak width. We estimate the line shape parameters for each bolometer–dataset with a simultaneous, unbinned extended maximum likelihood (UEML) fit performed on each tower in the energy range 2530–2720 keV. The simultaneous fit over a tower helps constrain common nuisance parameters such as relative intensity of x-ray escape peaks and continuum background. A simultaneous fit over the full array was not performed due to the computational demands. A comparison of the fit results with the calibration data and a breakdown of the fit model are shown in Figure 2.

To characterize possible differences in the detector response between physics and calibration data we fit

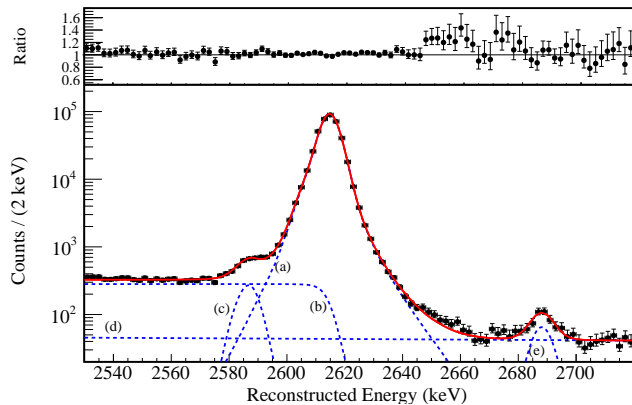


FIG. 2. Bottom: Sum of the results of the 19 tower-dependent UEML fits we use to estimate the line shape parameters of each bolometer–dataset in calibration data. The solid red line is the sum of the best-fit line shape model of each bolometer–dataset; the components of this summed best-fit model are shown by the blue dashed lines. We identify (a) the multi-Gaussian photopeak that describes the detector response function, (b) a multiscatter Compton contribution, (c) multiple peaks due to 27–31 keV Te X-ray escape following an incident 2615-keV  $\gamma$  ray, (d) a linear continuum background due to coincident events, and (e) a line due to coincident absorption of 2615-keV and 583-keV  $\gamma$  rays from the  $^{232}\text{Th}$  decay chain followed by escape of a 511-keV annihilation  $\gamma$  from pair production. Top: Ratio between calibration data and line shape model.

prominent background peaks in the physics data, with known energies between 800 and 2615 keV, using the best-fit line shape parameters determined above for each bolometer–dataset. At each energy this fit includes a dataset-dependent (i.e., channel independent) energy offset variable to parametrize energy misreconstruction. In addition, as the calibration line shape study was performed near 2615 keV, each fit includes a dataset-dependent (channel independent) energy resolution scaling variable to parameterize energy dependence of the resolution or a difference between background and calibration resolution. We find the energy misreconstruction is less than 0.5 keV over the calibrated energy range. The best-fit resolution scaling parameters at 2615 keV are  $0.95 \pm 0.07$  and  $1.01 \pm 0.06$  for the first and second dataset, respectively. To parametrize the energy dependence of the resolution scaling, we fit the set of scaling parameters determined at each peak energy studied with a quadratic function. The resulting best-fit function is then used to estimate the resolution scaling at  $Q_{\beta\beta}$ . The exposure-weighted harmonic mean energy resolution of the detectors (denoted *effective resolution*) in physics data, extrapolated to  $Q_{\beta\beta}$ , is given for each dataset in Table I; to quote a single characteristic energy resolution for our entire exposure, we combine these, finding  $(7.7 \pm 0.5)$  keV FWHM.

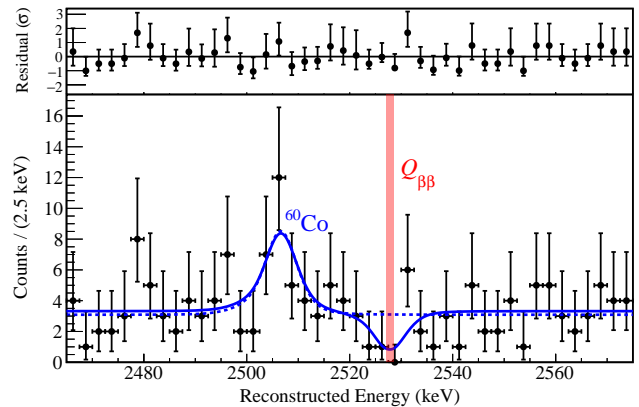


FIG. 3. Bottom: Best-fit model from the UEML fit (solid blue line) overlaid on the spectrum of  $0\nu\beta\beta$  decay candidates observed in CUORE. The peak near 2506 keV is attributed to  $^{60}\text{Co}$  [33]. The normalized residuals of this model and the binned data are shown in the top panel. The dashed (blue) curve shows the best-fit for a model with no  $0\nu\beta\beta$  decay component. The vertical band is centered at  $Q_{\beta\beta}$ ; the width of the band reflects the systematic uncertainty on the reconstructed energy.

Before unblinding the physics data, we fix the model and fitting strategy to search for the  $0\nu\beta\beta$  decay of  $^{130}\text{Te}$ . The ROI is taken from 2465 keV to 2575 keV. The model for each bolometer–dataset is composed of a  $0\nu\beta\beta$  decay peak, a peak for  $^{60}\text{Co}$  coincident  $\gamma$  rays, and a flat background. Each peak is modeled using the line shape discussed above, with the line width scaled by the resolution scaling extrapolated to the peak energy. All detectors are constrained to have the same  $0\nu\beta\beta$  decay rate  $\Gamma_{0\nu}$ , which we allow to vary freely in the fit; the position of the  $0\nu\beta\beta$  decay peak is fixed to  $Q_{\beta\beta}$  for each bolometer–dataset. The  $^{60}\text{Co}$  peak position is a dataset-dependent free parameter; the  $^{60}\text{Co}$  rate is a single free parameter but the known isotope half-life is used to account for its decay. The background rate is a dataset-dependent free parameter and is not scaled by the event selection efficiency.

Figure 3 shows the 155 candidate events in the ROI that pass all selection criteria together with the result of the UEML fit described above. The total  $\text{TeO}_2$  exposure is 86.3 kg·yr, corresponding 24.0 kg·yr for  $^{130}\text{Te}$ . The best-fit  $\Gamma_{0\nu}$  is  $(-1.0^{+0.4}_{-0.3} \text{ (stat.)} \pm 0.1 \text{ (syst.)}) \times 10^{-25} \text{ yr}^{-1}$ . With zero signal, the best-fit background in the ROI averaged over both datasets is  $(0.014 \pm 0.002)$  counts/(keV·kg·yr).

To evaluate the goodness of fit, we prepare a large set of pseudo-experiments, each with a number of events determined by a Poisson distribution with a mean of 155 and energy distributed according to the best-fit zero-signal model. Repeating our  $0\nu\beta\beta$  decay search fit on each of these, we find that 68% yield a negative log likelihood (NLL) larger than that obtained with our data.

TABLE II. Systematic uncertainties on  $\Gamma_{0\nu}$  for zero signal (additive) and as a percentage of nonzero signal (scaling).

	Additive ( $10^{-25} \text{ yr}^{-1}$ )	Scaling (%)
Line shape	0.02	2.4
Energy resolution	–	1.5
Fit bias	–	0.3
Energy scale	–	0.2
Background shape	0.05	0.8
Selection efficiency	2.4%	

We conclude there is no evidence for  $0\nu\beta\beta$  decay and set a 90% confidence Bayesian upper limit on the rate, finding  $\Gamma_{0\nu} < 0.50 \times 10^{-25} \text{ yr}^{-1}$  (stat. only) or  $T_{1/2}^{0\nu} > 1.4 \times 10^{25} \text{ yr}$ . In constructing the posterior pdf for  $\Gamma_{0\nu}$ , we approximate the marginalized likelihood with the profile likelihood and use a flat prior for  $\Gamma_{0\nu} \geq 0$ . This approximation speeds up the computation and is valid when the marginalization is dominated by the most probable values of the nuisance parameters. We expect this for our likelihood as the number of events is large and the background dominates. To confirm this we perform an independent analysis using the BAT toolkit [54] with the same prior but marginalize over the nuisance parameters. The results agree with those above to the percent level.

We repeat our analysis on a large set of pseudo-experiments generated in the same way as for the goodness of fit study. We find the median 90% confidence lower limit (sensitivity) for  $T_{1/2}^{0\nu}$  is  $7.0 \times 10^{24} \text{ yr}$ , and there is a 2% probability of obtaining a more stringent limit than the one obtained with our data.

We estimate the systematic uncertainties following the same procedure used for CUORE-0 [33]. We perform a large number of pseudo-experiments with zero and nonzero signals assuming different detector line shape models and background shapes (flat and first-order polynomial), varying the energy resolution scaling parameters within their uncertainty, and shifting the position of  $Q_{\beta\beta}$  by  $\pm 0.5 \text{ keV}$  to account for the energy reconstruction uncertainty. The results are summarized in Table II. We find the fit bias on  $\Gamma_{0\nu}$  to be negligible. Including these systematic uncertainties, the 90% confidence limits are  $\Gamma_{0\nu} < 0.52 \times 10^{-25} \text{ yr}^{-1}$  and  $T_{1/2}^{0\nu} > 1.3 \times 10^{25} \text{ yr}$ . A frequentist analysis [55] yields  $T_{1/2}^{0\nu} > 2.1 \times 10^{25} \text{ yr}$  at 90% C.L. with a median 90% C.L. lower limit sensitivity for  $T_{1/2}^{0\nu}$  of  $7.6 \times 10^{24} \text{ yr}$ .

We combine our profile likelihood curve with those from 9.8 kg·yr of  $^{130}\text{Te}$  exposure from CUORE-0 [32] and 19.8 kg·yr from Cuoricino [56] (see Figure 4). The combined 90% C.L. limits are  $\Gamma_{0\nu} < 0.47 \times 10^{-25} \text{ yr}^{-1}$  and  $T_{1/2}^{0\nu} > 1.5 \times 10^{25} \text{ yr}$ . The frequentist technique yields  $\Gamma_{0\nu} < 0.31 \times 10^{-25} \text{ yr}^{-1}$  and  $T_{1/2}^{0\nu} > 2.2 \times 10^{25} \text{ yr}$ .

We interpret the combined half-life limit,

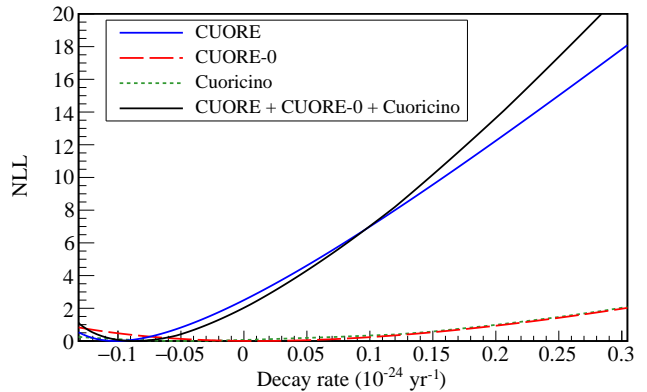


FIG. 4. Profile negative-log-likelihood curves for CUORE, CUORE-0, Cuoricino, and their combination.

$T_{1/2}^{0\nu} > 1.5 \times 10^{25} \text{ yr}$ , as a limit on the effective Majorana neutrino mass ( $m_{\beta\beta}$ ) in the framework of models of  $0\nu\beta\beta$  decay mediated by light Majorana neutrino exchange. We use phase-space factors from [57], nuclear matrix elements from a broad range of models [58–68], and assume the axial coupling constant  $g_A \simeq 1.27$ ; this yields  $m_{\beta\beta} < (110 - 520) \text{ meV}$  at 90% C.L., depending on the nuclear matrix element estimate employed.

In summary, we find no evidence for  $0\nu\beta\beta$  decay of  $^{130}\text{Te}$  and place the most stringent limit to date on this decay half-life. The observed background,  $(0.014 \pm 0.002) \text{ counts}/(\text{keV}\cdot\text{kg}\cdot\text{yr})$ , is in line with our expectations [51]. The characteristic energy resolution at  $Q_{\beta\beta}$  is  $(7.7 \pm 0.5) \text{ keV}$ , which we foresee improving to  $\sim 5 \text{ keV}$  by optimizing operating conditions and through analysis improvements. A study of our future sensitivity for a number of scenarios is presented in Ref. [69]. The experimental progress in  $0\nu\beta\beta$  decay searches has been dramatic in recent years; half-lives greater than  $10^{25} \text{ yr}$  are now probed by several experiments [70–73]. CUORE is the first ton-scale cryogenic detector array in operation, more than an order of magnitude larger than its predecessors. The successful commissioning and operation of this large-mass, low-background, cryogenic bolometer array represents a major advancement in the application of this technique to  $0\nu\beta\beta$  decay searches.

The CUORE Collaboration thanks the directors and staff of the Laboratori Nazionali del Gran Sasso and our technical staff for their valuable contribution to building and operating the detector. We thank Danielle Speller for contributions to detector calibration, data analysis, and discussion of the manuscript. We thank Giorgio Frossati for his crucial support in the commissioning of the dilution unit. This work was supported by the Istituto Nazionale di Fisica Nucleare (INFN); the National Science Foundation under Grant Nos. NSF-PHY-0605119, NSF-PHY-0500337, NSF-PHY-0855314, NSF-PHY-0902171, NSF-PHY-0969852, NSF-PHY-1307204,

and NSF-PHY-1404205; the Alfred P. Sloan Foundation; and Yale University. This material is also based upon work supported by the US Department of Energy (DOE) Office of Science under Contract Nos. DE-AC02-05CH11231 and DE-AC52-07NA27344; and by the DOE Office of Science, Office of Nuclear Physics under Contract Nos. DE-FG02-08ER41551, DE-FG03-00ER41138, DE-SC0011091, and DE-SC0012654. This research used resources of the National Energy Research Scientific Computing Center (NERSC).

- 
- <sup>a</sup> Present address: University College Leuven-Limburg, 3590 Diepenbeek, Belgium
- <sup>b</sup> Present address: Istituto Nazionale di Ottica, Consiglio Nazionale delle Ricerche (INO-CNR), Firenze I-50125, Italy
- <sup>c</sup> Applied Physics Laboratory, The Johns Hopkins University, Baltimore, MD 20723, USA
- <sup>d</sup> Corresponding author: [cuore-spokesperson@lngs.infn.it](mailto:cuore-spokesperson@lngs.infn.it)
- <sup>e</sup> Present address: INFN – Laboratori Nazionali del Gran Sasso, Assergi (L’Aquila) I-67100, Italy
- <sup>f</sup> Present address: Research Square, Durham, NC 27701, USA
- <sup>g</sup> Deceased
- <sup>h</sup> Present address: Physics Division, P-23, Los Alamos National Laboratory, Los Alamos, NM 87544, USA
- <sup>i</sup> Present address: Department of Physics, Drexel University, Philadelphia, PA 19104, USA
- <sup>j</sup> Present address: European Spallation Source (ESS), 225 92 Lund, Sweden
- <sup>k</sup> Present address: ITER Organization, 13067 St. Paul-lez-Durance, France
- <sup>l</sup> Present address: CSNSM, Univ. Paris-Sud, CNRS/IN2P3, Universit Paris-Saclay, 91405 Orsay, France
- <sup>m</sup> Present address: INFN – Laboratori Nazionali di Frascati, Frascati (Roma) I-00044, Italy
- <sup>n</sup> Present address: Laboratoire Commun de Métrologie, LNE-CNAM, 93210 La Plaine Saint-Denis, France
- <sup>o</sup> Present address: Max-Planck-Institut für Physik, 80805 München, Germany
- <sup>p</sup> Present address: ABB Corporate Research Center, 722 26 Västerås, Sweden
- <sup>q</sup> Present address: CryoVac GmbH & Co KG, 53842 Troisdorf, Germany
- <sup>r</sup> Present address: Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA
- [1] C. Patrignani *et al.* (Particle Data Group), *Chin. Phys. C* **40**, 100001 (2016).
- [2] E. Majorana, *Nuovo Cimento* **14**, 171 (1937).
- [3] M. Fukugita and T. Yanagida, *Phys. Lett. B* **174**, 45 (1986).
- [4] G. Racah, *Nuovo Cimento* **14**, 322 (1937).
- [5] W. H. Furry, *Phys. Rev.* **56**, 1184 (1939).
- [6] B. Pontecorvo, *Sov. Phys. JETP* **26**, 984 (1968).
- [7] J. Schechter and J. Valle, *Phys. Rev. D* **25**, 2951 (1982).
- [8] S. Dell’Oro, S. Marocci, M. Viel, and F. Vissani, *Adv. High Energy Phys.* **2016**, 2162659 (2016).
- [9] D. R. Artusa *et al.*, *Adv. High Energy Phys.* **2015**,

- 879871 (2015).
- [10] E. Fiorini and T. Niinikoski, *Nucl. Instrum. Meth.* **224**, 83 (1984).
- [11] C. Enss and D. McCammon, *J. Low Temp. Phys.* **151**, 5 (2008).
- [12] M. A. Fehr, M. Rehkämper, and A. N. Halliday, *Int. J. Mass Spectrom.* **232**, 83 (2004).
- [13] M. Redshaw, B. J. Mount, E. G. Myers, and F. T. Avignone, III, *Phys. Rev. Lett.* **102**, 212502 (2009).
- [14] N. D. Scielzo *et al.*, *Phys. Rev. C* **80**, 025501 (2009).
- [15] S. Rahaman *et al.*, *Phys. Lett. B* **703**, 412 (2011).
- [16] C. Arnaboldi *et al.*, *J. Cryst. Growth* **312**, 2999 (2010).
- [17] E. E. Haller, N. P. Palaio, M. Rodder, W. L. Hansen, and E. Kreysa, in *Neutron Transmutation Doping of Semiconductor Materials* (Springer US, Boston, MA, 1984) pp. 21–36.
- [18] A. Alessandrello *et al.*, *Nucl. Instrum. Meth. A* **412**, 454 (1998).
- [19] E. Andreotti, C. Brofferio, L. Foggetta, A. Giuliani, B. Margesin, C. Nones, M. Pedretti, C. Rusconi, C. Salvioni, and M. Tenconi, *Nucl. Instrum. Meth. A* **664**, 161 (2012).
- [20] Leiden Cryogenics DRS-CF3000 continuous-cycle.
- [21] Cryomech PT415-RM.
- [22] A. Alessandrello *et al.*, *Nucl. Instrum. Meth.* **142**, 163 (1998).
- [23] F. Alessandria *et al.*, “The CUORE cryostat: A 10 mK infrastructure for large bolometric arrays,” (In preparation).
- [24] A. Bersani *et al.*, “The detector suspension system of the CUORE experiment,” (In preparation).
- [25] G. Bucci, P. Carniti, L. Cassina, C. Gotti, A. Pelosi, G. Pessina, M. Turqueti, and S. Zimmermann, *JINST* **12**, P12013 (2017).
- [26] C. Arnaboldi, P. Carniti, L. Cassina, C. Gotti, X. Liu, and G. Pessina, “Front-end electronic systems for large arrays of bolometric experiments,” (2017), accepted to *JINST*, [arXiv:1710.06365](https://arxiv.org/abs/1710.06365).
- [27] G. Benato, D. Biare, C. Bucci, L. D. Paolo, A. Drobizhev, R. Kadel, Y. Kolomensky, J. Schreiner, V. Singh, T. Sipla, J. Wallig, and S. Zimmermann, *Journal of Instrumentation* **13**, P01010 (2018).
- [28] F. Alessandria *et al.*, *Astropart. Phys.* **45**, 13 (2013).
- [29] E. Buccheri, M. Capodiferro, S. Morganti, F. Orio, A. Pelosi, and V. Pettinacci, *Nucl. Instrum. Meth. A* **768**, 130 (2014).
- [30] C. Alduino *et al.* (CUORE Collaboration), *J. Instrum.* **11**, P07009 (2016).
- [31] E. Andreotti *et al.*, *JINST* **4**, P09003 (2009).
- [32] K. Alfonso *et al.* (CUORE Collaboration), *Phys. Rev. Lett.* **115**, 102502 (2015).
- [33] C. Alduino *et al.* (CUORE Collaboration), *Phys. Rev. C* **93**, 045503 (2016).
- [34] C. Alduino *et al.* (CUORE Collaboration), *Eur. Phys. J. C* **77**, 13 (2017).
- [35] J. L. Ouellet, *The Search for  $0\nu\beta\beta$  Decay of  $^{130}\text{Te}$  with CUORE-0*, Ph.D. thesis, University of California, Berkeley (2015).
- [36] A. D’Addabbo, C. Bucci, L. Canonica, S. Di Domizio, P. Gorla, L. Marini, A. Nucciotti, I. Nutini, C. Rusconi, and B. Welliver, “An active noise cancellation technique for the CUORE Pulse Tube Cryocoolers,” Submitted to *Cryogenics*, [arXiv:1712.02753](https://arxiv.org/abs/1712.02753).
- [37] J. S. Cushman *et al.*, *Nucl. Instrum. Meth. A* **844**, 32



- (2017).
- [38] P. Carniti, L. Cassina, C. Gotti, M. Maino, and G. Pessina, *Rev. Sci. Instrum.* **87**, 054706 (2016).
- [39] C. Arnaboldi *et al.*, *Rev. Sci. Instrum.* **86**, 124703 (2015).
- [40] C. Arnaboldi, M. Cariello, S. Di Domizio, A. Giachero, and G. Pessina, *Nucl. Instrum. Meth. A* **617**, 327 (2010).
- [41] A. Giachero, *Characterization of cryogenic bolometers and data acquisition system for the CUORE experiment*, Ph.D. thesis, Università degli studi di Genova (2008).
- [42] S. Di Domizio, *Search for double beta decay to excited states with CUORICINO and data acquisition system for CUORE*, Ph.D. thesis, Università degli studi di Genova (2009).
- [43] S. Copello, *Dark Matter Induced Annual Modulation Analysis in CUORE-0*, Ph.D. thesis, Università degli studi di Genova (2017).
- [44] S. Di Domizio, F. Orio, and M. Vignati, *J. Instrum.* **6**, P02007 (2011).
- [45] C. Alduino *et al.* (CUORE Collaboration), *Eur. Phys. J. C* **77**, 857 (2017).
- [46] P. Carniti, L. Cassina, A. Giachero, C. Gotti, and G. Pessina, “A High Precision Pulse Generation and Stabilization System for Bolometric Experiments,” (2017), accepted to JINST, [arXiv:1710.05565](https://arxiv.org/abs/1710.05565).
- [47] E. Gatti and P. F. Manfredi, *Nuovo Cimento* **9**, 1 (1986).
- [48] P. C. Mahalanobis, *Proc. Natl. Inst. Sci. India*, 49 (1936).
- [49] G. Cowan, K. Cranmer, E. Gross, and O. Vitells, *Eur. Phys. J. C* **71**, 1554 (2011).
- [50] S. Agostinelli *et al.*, *Nucl. Instr. Meth. A*, 250 (2003).
- [51] C. Alduino *et al.* (CUORE Collaboration), *Eur. Phys. J. C* **77**, 543 (2017).
- [52] A. D. Bryant, *A Search for Neutrinoless Double Beta Decay of  $^{130}\text{Te}$* , Ph.D. thesis, University of California, Berkeley (2010).
- [53] M. Carrettoni, *Data analysis for Neutrinoless Double Beta Decay*, Ph.D. thesis, Università degli Studi di Milano-Bicocca (2011).
- [54] A. Caldwell, D. Kollar, and K. Kroninger, *Comput. Phys. Commun.* **180**, 2197 (2009).
- [55] W. A. Rolke, A. M. Lopez, and J. Conrad, *Nucl. Instrum. Meth. A* **551**, 493 (2005), (bounded method).
- [56] E. Andreotti *et al.*, *Astropart. Phys.* **34**, 822 (2011).
- [57] J. Kotila and F. Iachello, *Phys. Rev. C* **85**, 034316 (2012).
- [58] J. Engel and J. Menéndez, *Rept. Prog. Phys.* **80**, 046301 (2017), [arXiv:1610.06548](https://arxiv.org/abs/1610.06548) [nucl-th].
- [59] J. Barea, J. Kotila, and F. Iachello, *Phys. Rev. C* **91**, 034304 (2015).
- [60] F. Šimkovic, V. Rodin, A. Faessler, and P. Vogel, *Phys. Rev. C* **87**, 045501 (2013).
- [61] J. Hyvärinen and J. Suhonen, *Phys. Rev. C* **91**, 024613 (2015).
- [62] J. Menendez, A. Poves, E. Caurier, and F. Nowacki, *Nucl. Phys. A* **818**, 139 (2009).
- [63] T. R. Rodriguez and G. Martinez-Pinedo, *Phys. Rev. Lett.* **105**, 252503 (2010).
- [64] N. López Vaquero, T. R. Rodríguez, and J. L. Egido, *Phys. Rev. Lett.* **111**, 142501 (2013).
- [65] J. M. Yao, L. S. Song, K. Hagino, P. Ring, and J. Meng, *Phys. Rev. C* **91**, 024316 (2015).
- [66] M. T. Mustonen and J. Engel, *Phys. Rev. C* **87**, 064302 (2013), [arXiv:1301.6997](https://arxiv.org/abs/1301.6997) [nucl-th].
- [67] A. Neacsu and M. Horoi, *Phys. Rev. C* **91**, 024309 (2015).
- [68] A. Meroni, S. T. Petcov, and F. Šimkovic, *JHEP* **02**, 025 (2013), [arXiv:1212.1331](https://arxiv.org/abs/1212.1331) [hep-ph].
- [69] C. Alduino *et al.* (CUORE Collaboration), *Eur. Phys. J. C* **77**, 532 (2017).
- [70] A. Gando *et al.* (KamLAND-Zen Collaboration), *Phys. Rev. Lett.* **117**, 082503 (2016).
- [71] J. B. Albert *et al.* (EXO-200), *Nature* **510**, 229 (2014).
- [72] M. Agostini *et al.* (GERDA), *Nature* **544**, 47 (2017).
- [73] J. B. Albert *et al.* (EXO), (2017), [arXiv:1707.08707](https://arxiv.org/abs/1707.08707) [hep-ex].