



CHORUS

This is the accepted manuscript made available via CHORUS. The article has been published as:

Coherent Bichromatic Force Deflection of Molecules

Ivan Kozyryev, Louis Baum, Leland Aldridge, Phelan Yu, Edward E. Eyler, and John M. Doyle

Phys. Rev. Lett. **120**, 063205 — Published 8 February 2018

DOI: [10.1103/PhysRevLett.120.063205](https://doi.org/10.1103/PhysRevLett.120.063205)

Coherent Bichromatic Force Deflection of Molecules

Ivan Kozyryev,^{1,2,*} Louis Baum,^{1,2,*} Leland Aldridge,^{3,†} Phelan Yu,^{1,2} Edward E. Eyler,³ and John M. Doyle^{1,2}

¹Harvard-MIT Center for Ultracold Atoms, Cambridge, MA 02138, USA

²Department of Physics, Harvard University, Cambridge, MA 02138, USA

³Department of Physics, University of Connecticut, Storrs, CT 06269, USA

(Dated: January 19, 2018)

We demonstrate the effect of the coherent optical bichromatic force on a molecule, the polar free radical strontium monohydroxide (SrOH). A dual-frequency retroreflected laser beam addressing the $\tilde{X}^2\Sigma^+ \leftrightarrow \tilde{A}^2\Pi_{1/2}$ electronic transition coherently imparts momentum onto a cryogenic beam of SrOH. This directional photon exchange creates a bichromatic force that transversely deflects the molecules. By adjusting the relative phase between the forward and counter propagating laser beams we reverse the direction of the applied force. A momentum transfer of $70\hbar k$ is achieved with minimal loss of molecules to dark states. Modeling of the bichromatic force is performed via direct numerical solution of the time-dependent density matrix and is compared with experimental observations. Our results open the door to further coherent manipulation of molecular motion, including the efficient optical deceleration of diatomic and polyatomic molecules with complex level structures.

Laser manipulation of atomic motion has revolutionized atomic, molecular and optical (AMO) physics [1, 2]. The widely-used techniques of laser cooling and trapping made possible the creation of ultracold degenerate quantum gases [3], simulation of important condensed matter models [4] and development of new quantum sensors [5, 6] and clocks [7]. Laser deceleration and cooling of atomic beams – a necessary part of the trap loading process – typically requires scattering tens of thousands of photons in order to bring room (or oven) temperature atoms to velocities where they can be confined by electromagnetic traps for further studies [8]. While beam deceleration employing the spontaneous radiation pressure force has been a standard for atomic experiments, its application to slowing molecular beams has been limited by the small change in kinetic energy per scattered photon and the myriad of internal molecular states which inhibits photon cycling. At the same time, there is extreme interest in creating ultracold molecules for new physics applications [9].

Neutral diatomic molecules are predicted to play an important role in diverse research areas of modern physics such as quantum simulation [10] and computation [11], as well as searches for new particles and fields beyond the Standard Model [12, 13]. Larger polyatomic molecules will provide additional opportunities in physics and chemistry [14–17]. For example, exploring the origin of biomolecular homochirality [18, 19] and understanding primordial chemistry leading to the development of organic life requires the use of large molecules [20]. However, these molecules’ complexity presents significant challenges for direct laser slowing and cooling. Yet these are the key ingredients that allow for optical trapping, which, in turn, leads to long molecule-laser coherence times and high levels of quantum state control. Previously, the external motion of gas-phase polyatomic molecules has been manipulated with off-resonant laser fields [21, 22] as well as electric [23, 24], magnetic [25], and mechanical techniques [26]. Inspired by the success of laser control of atomic motion as well as the latest developments in high-power CW and pulsed laser technology, many experimental proposals and theoretical calculations have been presented on

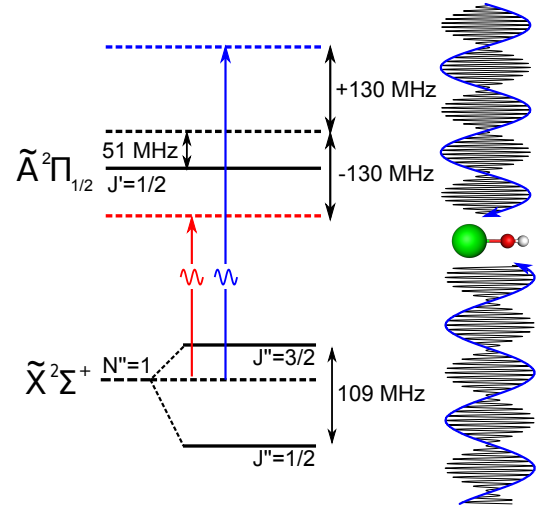


Figure 1. The energy diagram on the left-hand side depicts optimal bichromatic frequency detunings for SrOH addressed on the $\tilde{X}^2\Sigma^+ \rightarrow \tilde{A}^2\Pi_{1/2} P(N''=1)$ transition for 10 W/cm^2 irradiance per each frequency component with 688 nm laser light. Solid horizontal lines indicate positions of the energy levels. The parameters were estimated using numerical solutions of optical Bloch equations. Amplitude modulated optical pulses seen by SrOH molecules are shown on the right-hand side.

using *stimulated* light forces for molecular beam slowing [27–33], yet there has been no successful experimental implementation.

In this Letter, we demonstrate and characterize the optical bichromatic force (BCF) for molecules by deflecting a cryogenic buffer-gas beam (CBGB) of SrOH. Using dual-frequency high-power standing light waves we achieve significant force enhancement compared to radiative deflection. The coherent nature of the directional momentum transfer allows multiple $\hbar k$ of momentum change per spontaneous emission cycle. We also perform theoretical calculations of BCF in complex multilevel systems and compare to our data. Our results enable accurate predictions for other molecular species as well as different experimental parameters. Thirty years af-

ter the initial theoretical proposals [34, 35], the bichromatic force for molecules has been conclusively demonstrated.

Radiation pressure beam slowing using white-light [36–38] and chirped techniques [39, 40] has been achieved for a few simple diatomic species with highly diagonal Franck-Condon factors (FCFs). However, the reduced scattering rate (due to multiple ground state sublevels) leads to a long slowing distance and lowers the capture efficiency of molecules from a beam into a magneto-optical trap (MOT) [41]. Significant reduction of the slowing distance in order to mitigate transverse pluming of the molecular beam would result in enhanced MOT densities, desired for diverse applications. Furthermore, extending beam slowing to species with non-diagonal FCFs (the typical situation for molecules) is extremely challenging due to the rapid loss to dark states after only a few scattered photons.

The idea of using *stimulated* slowing and cooling (first for atoms) dates back to the end of 1980s when the rectified dipole force was theoretically proposed [34, 42] and experimentally observed for sodium atoms [35, 43]. Using two standing light waves of different frequencies and detunings, the dipole force can be rectified to apply force on atoms on the macroscopic spatial scale. The discovery of the bichromatic force, which derives from a rapid coherent sequence of absorption-stimulated emission cycles, led to a technique with a wide velocity capture range and robustness under experimental imperfections. This allowed significant advances in atom manipulation [44]. For example, short-distance deceleration of Cs atomic beam with BCF was demonstrated, achieving a force ten times greater than radiative force [45]. Additionally, BCF has been used for He [46] and Ar [47] atomic beam collimation. However, in atoms stimulated forces are often cast aside in favor of technically simpler technologies (e.g. radiative or Zeeman slowing) that depend on the rigorous selection rules present only in specific atoms.

BCF allows rapid directional momentum exchange of multiple $\hbar k$ between the light field and atoms in a single spontaneous emission cycle [45]. Intuitively, the BCF mechanism can be thought of as the result of a series of optical pulses formed by the beating of the dual-frequency light. If applied from two counter propagating directions and timed carefully, a pulse from one direction will drive the molecule into an excited state and a pulse from the opposite direction will stimulate coherent emission back to the original state. One pair of pulses will leave the molecule in the original internal state but will have removed $2\hbar k$ of momentum. Optical pulses created by the bichromatic fields detuned $\pm\delta$ relative to the transition resonance ω are shown in Fig. 1. Under the properly chosen laser intensity, each beat envelope of duration π/δ drives a state inversion, transferring $\hbar k$ momenta from the laser field [48]. In a two-level system, the magnitude of the achieved force becomes $F_{\text{BCF}} = \hbar k \delta / \pi$, which can be much larger than $F_{\text{rad},2\text{-level}} = \hbar k \gamma / 2$ for $\delta \gg \gamma$. Rotational motion in molecules leads to a significant reduction in the magnitude of the maximum radiative force compared to $F_{\text{rad},2\text{-level}}$, further motivating the need for coherent stimulated slowing. Since many $\hbar k$

of momentum can be transferred for every spontaneous emission, BCF can be effectively applied to molecules that have optically accessible electronic transitions but suffer from the loss to dark states (e.g. excited electronic or vibrational levels). Particularly, application of BCF to polyatomic radicals that have multiple vibrational modes or diatomic molecules with complex electronic structures or non-diagonal Franck-Condon factors could allow production of slow, velocity controlled molecular beams for physics and chemistry applications.

SrOH provides a convenient platform for exploring the effects of BCF on complex molecules. As a triatomic molecule, it contains a significant increase in complexity compared to diatomics, including a large number of low-lying bending mode states [49]. A single non-bonding valence electron allows laser addressing of internal quantum states but the loss to excited stretching and bending vibrational levels inhibits the photon cycling process [50, 51]. Upon electronic excitation to the \tilde{A} state, molecules predominantly decay to the ground vibrational state (95.5%) with a small fraction decaying to the Sr-O stretching mode with 1 (4.3%) and 2 (0.17%) quanta as well as bending vibrations without vibrational angular momentum (0.04%) [51].

The exact technical details for the cryogenic production, optical state manipulation, and photon cycling detection of the SrOH beam have been provided in our previous publications [50, 52] as well as in the Supplemental Material (SM). SrOH in a CBGB is collimated using a 0.5×2 mm aperture before encountering two collinearly superimposed standing waves of $1/e^2$ Gaussian beam diameter $\varnothing \approx 0.9$ mm. We use a bichromatic detuning $\delta = \pm 130$ MHz per each frequency component and a peak on-axis Gaussian beam intensity of 10-11 W/cm². The 688 nm laser light is generated by a tapered amplifier seeded with an injection locked diode laser. An acousto-optic modulator operating at 260 MHz generates the required frequency offset as described in the SM. Figure 1 presents the relevant molecular structure of SrOH and detunings used in the experiment. A magnetic field of ~ 19 Gauss is used to rapidly remix the magnetic dark states created by addressing $J_g > J_e$ molecular structure with a single laser polarization.

The beat notes interacting with the molecular beam from both directions are generated by the same bichromatic traveling wave upon the reflection from a stationary mirror. The relative phase ϕ of the irradiance of the counter propagating traveling waves is controlled by the distance d_ϕ between the retroreflecting mirror and molecular beam. Since $\phi = 4d_\phi\delta/c$ [53], the direction of the molecular deflection can be controlled by choosing the appropriate mirror position d_ϕ . Thus, for $\pi/2$ BCF phase shift we use $d_\phi = 14.4$ cm and for $3\pi/2$ phase $d_\phi = 43.2$ cm.

The effect of the BCF on the SrOH molecular beam is illustrated in Fig. 2. The comparison of the radiative force shift (left) with the bichromatic force shift (right) clearly demonstrates larger momentum transfer. This data was taken in the following way. The bichromatic force configuration has two applied beams. To take the radiative force data, the retrore-

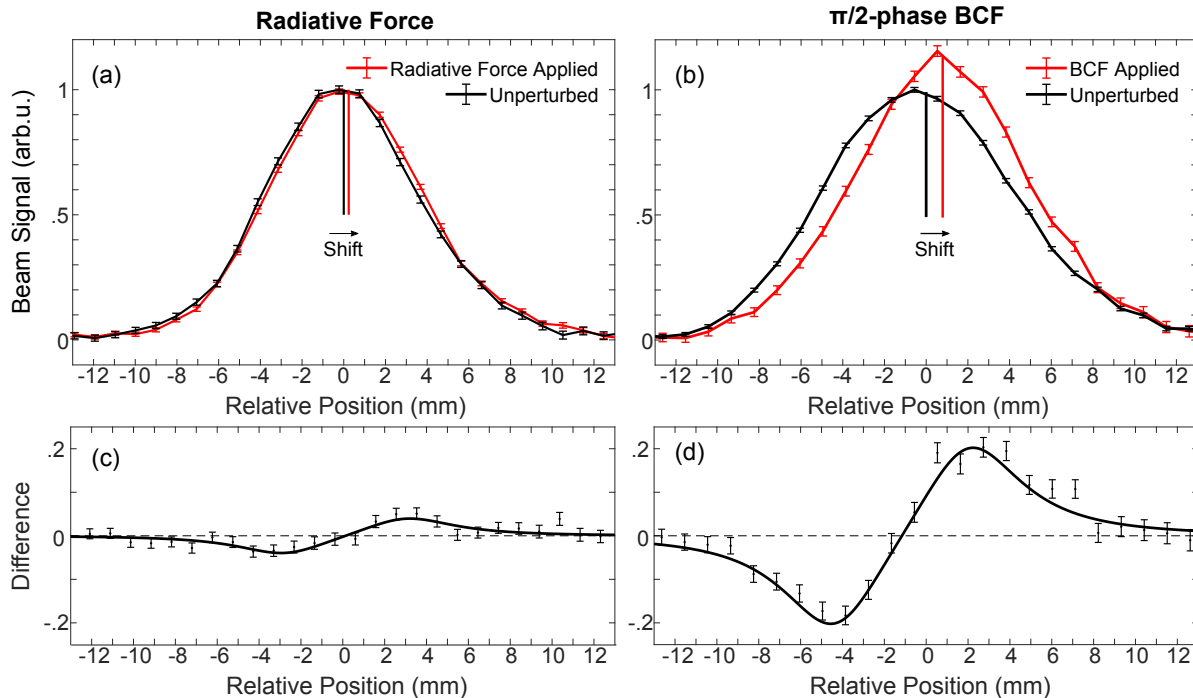


Figure 2. Data demonstrating bichromatic force deflection of SrOH beam with the $\pi/2$ relative phase shift between the counter-propagating laser beams. Plots (a) and (c) illustrate the effect of purely radiative force on the molecular beam. Plots (b) and (d) compare the beam profiles for unperturbed molecular beam and with BCF applied. The same vertical axis is used for (a) and (b) as well as (c) and (d) plots.

flected beam is simply blocked. Because laser intensities are deliberately balanced, in the absence of the directional momentum transfer, there should be no net shift in the center of the molecular beam. However, we see a deflection of the molecular beam with a larger magnitude than with the radiative force alone.

In order to extract the magnitude of the BCF on SrOH and characterize several experimental conditions we plot the absolute value of the beam deflection in Fig. 3. While the direction of the radiative deflection is defined to be positive, by choosing a proper relative phase $\phi = 3\pi/2$ between the dual-frequency laser beams transversely addressing the molecules from opposite sides the direction of the BCF can be reversed. The reversal of the deflection direction with $3\pi/2$ -phase BCF relative to the radiative force with the addition of the counter-propagating beam conclusively indicates the presence of the BCF effect in our experiment. By comparing the ratio of the absolute deflection magnitude for $\pi/2$ and $3\pi/2$ phases we obtain that BCF is 3.7 ± 0.7 and 2.7 ± 0.5 times greater than radiative force, correspondingly. A slightly weaker force for $\phi = 3\pi/2$ configuration arises from the larger relative beam diameter imbalance because of the longer d_ϕ distance. This challenge could be resolved by using separate dual-frequency laser beams shaped identically to the original but with a phase delay introduced. This approach has already been shown to work well for atomic BCF slowing [45]. For the π -phase configuration there is no directional momentum transfer between the light beams and molecules and therefore no molec-

ular beam deflection as seen in Fig. 3, as expected for BCF.

From the BCF deflection data we can estimate that $(68 \pm 5) \hbar k$ units of momentum have been coherently transferred from the laser fields to molecules leading to a deflection of 0.3° . With the 631 nm clean-up beam for repumping molecules that decayed to the first excited Sr-O stretching mode (see SM for details), we observe no loss of molecules to other dark states within the experimental uncertainties. Thus, for the operating configuration with the repumper present, larger deflection magnitudes can be achieved by extending the molecule-laser interaction time while preserving the proper irradiance condition to achieve efficient stimulated π -pulse transfer.

The BCF effect on SrOH involves a coherent momentum exchange between molecules and optical fields consisting of two phase-shifted dual-frequency laser beams, which couple 12 ground state magnetic sublevels to 4 excited states. Unlike for atoms [54, 55], no analytical solutions exist to describe this process and statistical scaling of the two-level expressions [27] has proven to be highly inaccurate [30]. Therefore, detailed numerical modeling of the bichromatic force in molecules is necessary. In order to estimate the bichromatic force effect on SrOH we have employed direct numerical solution for the time-dependent density matrix in the rotating-wave approximation [56].

Figure 4 presents the results of the numerical calculations for the BCF effect on SrOH as a function of molecular velocity along the laser propagation direction. The average theo-

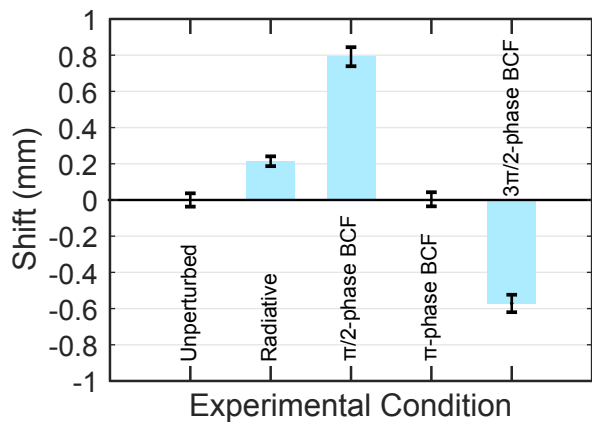


Figure 3. Summary of the molecular beam centroid shift for various experimental conditions: unperturbed molecular beam, radiative deflection with retroreflected beam blocked and BCF deflection with $\pi/2$, π and $3\pi/2$ relative phase shift. Reversal of the BCF directionality with the relative phase change is observed.

retically calculated BCF force exerted on molecules with velocities $v \leq 25$ m/s is $(1.89 \pm 0.04) \hbar k \gamma / 2$, which exceeds the maximum radiative force in the experiment by 5.7 times. Using the measured optimal BCF deflections for SrOH we obtain F_{BCF} which is 3.7 ± 0.7 times the measured radiative force, which is 35% lower compared to the predicted value. Experimental observations of BCF for atoms have seen significant reduction in the force magnitude under realistic experimental conditions [47]. Increasing the remixing magnetic field above 25 Gauss should lead to 18% gain in the bichromatic force [57]. In our experiment, some of additional explanations for the discrepancy are power imbalance between the two beat note trains, less than unity overlap efficiency between the laser and molecular beams, and a dependence of the two forces on irradiance which favors the radiative force in the wings of the Gaussian laser profile.

In earlier experiments, both longitudinal [45] and transverse [46] cooling of atomic beams have been observed, including in the ultimate limit of $\lesssim 1$ scattered photon [58]. While there is contention in the community over the presence of BCF induced cooling without spontaneous emission [59], our work does not shed any light on this debate because of the presence of spontaneously emitted photons in the experiment. Our experimental results and theoretical calculations indicate compression of the transverse velocity distribution. From Fig. 2 we conclude that the full width at half maximum of the beam distribution is reduced from 10.6 ± 0.2 mm to 9.3 ± 0.2 mm, providing evidence for cooling. Interpreted in combination with previous atomic experiments [46], our results indicate the possibility of creating optical molasses-like bichromatic field configurations characterized by rapid damping forces and wide velocity capture range using two sequential spatially separated regions with opposite ϕ .

In summary, we demonstrate a coherent optical force on polyatomic molecules using bichromatic laser fields. SrOH

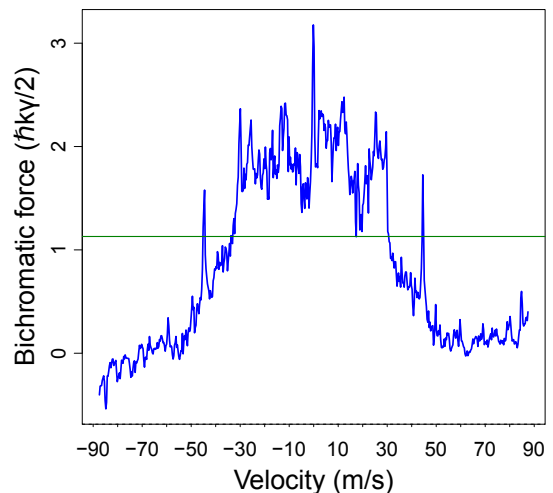


Figure 4. Calculated profile of the bichromatic force exerted on SrOH molecules as a function of velocity under optimal experimental conditions. The force magnitude is normalized to the size of the maximum radiative force in an ideal two-level system ($F_{\text{rad}, 2\text{-level}} = \hbar k \gamma / 2$). Sharp resonances around 0 and ± 45 m/s represent Doppler resonances and have been previously seen in BCF profiles for atoms [47]. Solid horizontal line shows the average bichromatic force achieved in the experiment. The calculated force profile assumes an irradiance of 10 W/cm^2 , a magnetic field of 27.2 G skewed 33° with respect to the polarization axis and the frequency parameters depicted in Fig. 1.

has 12 ground state sublevels coupled to 4 excited states, yet we achieve $F_{\text{BCF}} \approx 1.1 \hbar k \gamma / 2$ which is effectively the size of the radiative force on an ideal two-level system. This can be compared with the radiative force experiments on diatomic [60] and polyatomic molecules [52], which have previously shown $F_{\text{rad}} \approx \hbar k \gamma / 7$ and $\hbar k \gamma / 4$, respectively. Thus, large resonant optical forces on molecules are possible with the use of coherent bichromatic optical fields to induce stimulated forces.

The magnitude of the bichromatic force achieved in this experiment is limited by the available optical power from our laser system. Our calculations indicate that with 1 W per frequency component and using a 2 mm $1/e^2$ Gaussian beam diameter, the average value of the BCF will be $7 \hbar k \gamma / 2$ for molecules spanning a 60 m/s range [61]. Slowing of a cryogenic beam of SrOH from 60 m/s to near zero in a distance of $\lesssim 5$ mm could be achieved with only 500 scattered photons, compared to 10,000 photons (and a much longer distance) required for radiative slowing. The use of polychromatic forces with higher harmonics could further enhance the magnitude of stimulated forces by reducing the overlap of counter-propagating beat pulses, in addition to extending the velocity capture range [62, 63]. The powerful combination of cryogenic buffer-gas beam technology and stimulated optical forces has the potential to deliver slow molecules with complex internal structures for precision beam measurements or for trapping. For example, BCF can be used for slowing heavy polyatomic molecules like YbOH, YbCH₃ and

YbOCH₃ to search for new fundamental particles at the PeV energy scale [64]. Additionally, we expect that bichromatic force slowing can be applied to a large number of monovalent calcium and strontium derivatives with optically accessible electronic transitions [65] including monomethyls (MCH₃), alkoxides (MOR) and monocyclopentadienides (MC₅H₅) to deliver slow beams for trapping and exploration of cold controlled collisions and chemistry [66].

We would like to thank Kyle Matsuda and Yicheng Bao for building the TA laser system. We also thank Scott Galica, Zack Lasner, Loic Anderegg, and Ben Augenbraun for many insightful discussions. This work has been supported by the NSF, grant # PHY-1505961.

* These authors contributed equally. ivan@cua.harvard.edu, louis@cua.harvard.edu

† Present address: Department of Physics, Yale University, New Haven, CT 06520, USA

- [1] S. Chu, *Rev. Mod. Phys.* **70**, 685 (1998).
- [2] W. D. Phillips, *Rev. Mod. Phys.* **70**, 721 (1998).
- [3] W. Ketterle, *Rev. Mod. Phys.* **74**, 1131 (2002).
- [4] I. Bloch, J. Dalibard, and W. Zwerger, *Rev. Mod. Phys.* **80**, 885 (2008).
- [5] S. M. Dickerson, J. M. Hogan, A. Sugarbaker, D. M. Johnson, and M. A. Kasevich, *Phys. Rev. Lett.* **111**, 083001 (2013).
- [6] A. D. Cronin, J. Schmiedmayer, and D. E. Pritchard, *Rev. Mod. Phys.* **81**, 1051 (2009).
- [7] A. D. Ludlow, M. M. Boyd, J. Ye, E. Peik, and P. O. Schmidt, *Rev. Mod. Phys.* **87**, 637 (2015).
- [8] H. J. Metcalf and P. Van der Straten, *Laser cooling and trapping* (New York: Springer Science & Business Media, 2012).
- [9] L. D. Carr, D. DeMille, R. V. Krems, and J. Ye, *New J. Phys.* **11**, 055049 (2009).
- [10] J. L. Bohn, A. M. Rey, and J. Ye, *Science* **357**, 1002 (2017).
- [11] D. DeMille, *Phys. Rev. Lett.* **88**, 067901 (2002).
- [12] D. DeMille, *Phys. Today* **68**, 34 (2015).
- [13] D. DeMille, J. M. Doyle, and A. O. Sushkov, *Science* **357**, 990 (2017).
- [14] M. L. Wall, K. Maeda, and L. D. Carr, *Ann. Phys. (Berlin)* **525**, 845 (2013).
- [15] Q. Wei, S. Kais, B. Friedrich, and D. Herschbach, *J. Chem. Phys.* **135**, 154102 (2011).
- [16] C. M. Tesch and R. de Vivie-Riedle, *Phys. Rev. Lett.* **89**, 157901 (2002).
- [17] M. Wall, K. Maeda, and L. D. Carr, *New J. Phys.* **17**, 025001 (2015).
- [18] M. Quack, *Angew. Chem. Int. Ed.* **41**, 4618 (2002).
- [19] M. Quack, *Angew. Chem. Int. Ed.* **28**, 571 (1989).
- [20] A. Lazcano and S. L. Miller, *Cell* **85**, 793 (1996).
- [21] R. Fulton, A. I. Bishop, and P. F. Barker, *Phys. Rev. Lett.* **93**, 243004 (2004).
- [22] H. Stapelfeldt, H. Sakai, E. Constant, and P. B. Corkum, *Phys. Rev. Lett.* **79**, 2787 (1997).
- [23] H. L. Bethlem, G. Berden, F. M. Crompvoets, R. T. Jongma, A. J. Van Roij, and G. Meijer, *Nature* **406**, 491 (2000).
- [24] A. Prehn, M. Ibrügger, R. Glöckner, G. Rempe, and M. Zeppenfeld, *Phys. Rev. Lett.* **116**, 063005 (2016).
- [25] Y. Liu, M. Vashishta, P. Djuricanin, S. Zhou, W. Zhong, T. Mitterreiner, D. Carty, and T. Momose, *Phys. Rev. Lett.* **118**, 093201 (2017).
- [26] S. Chervenkov, X. Wu, J. Bayerl, A. Rohlfes, T. Gantner, M. Zeppenfeld, and G. Rempe, *Phys. Rev. Lett.* **112**, 013001 (2014).
- [27] M. Chieda and E. Eyler, *Phys. Rev. A* **84**, 063401 (2011).
- [28] D. Dai, Y. Xia, Y. Fang, L. Xu, Y. Yin, X. Li, X. Yang, and J. Yin, *J. Phys. B At. Mol. Opt. Phys.* **48**, 085302 (2015).
- [29] X. Yang, C. Li, Y. Yin, S. Xu, X. Li, Y. Xia, and J. Yin, *J. Phys. B At. Mol. Opt. Phys.* **50**, 015001 (2016).
- [30] L. Aldridge, S. Galica, and E. Eyler, *Phys. Rev. A* **93**, 013419 (2016).
- [31] E. Ilinova, J. Weinstein, and A. Derevianko, *New J. Phys.* **17**, 055003 (2015).
- [32] A. Jayich, A. Vutha, M. Hummon, J. Porto, and W. Campbell, *Phys. Rev. A* **89**, 023425 (2014).
- [33] V. Romanenko, Y. G. Udovitskaya, A. Romanenko, and L. Yatsenko, *Phys. Rev. A* **90**, 053421 (2014).
- [34] A. Kazantsev and I. Krasnov, *JETP Lett.* **46**, 264 (1987).
- [35] V. Voitsekhovich, M. Danileiko, A. Negriiko, V. Romanenko, and L. Yatsenko, *JETP Lett.* **49**, 161 (1989).
- [36] J. F. Barry, E. S. Shuman, E. B. Norrgard, and D. P. Demille, *Phys. Rev. Lett.* **108**, 103002 (2012).
- [37] B. Hemmerling, E. Chae, A. Ravi, L. Anderegg, G. K. Drayna, N. R. Hutzler, A. L. Collopy, J. Ye, W. Ketterle, and J. M. Doyle, *J. Phys. B: At. Mol. Opt. Phys.* **49**, 174001 (2016).
- [38] L. Anderegg, B. Augenbraun, E. Chae, B. Hemmerling, N. R. Hutzler, A. Ravi, A. Collopy, J. Ye, W. Ketterle, and J. M. Doyle, *Phys. Rev. Lett.* **119**, 103201 (2017).
- [39] S. Truppe, H. Williams, M. Hambach, N. Fitch, T. Wall, E. Hinds, B. Sauer, and M. Tarbutt, *New J. Phys.* **19**, 022001 (2017).
- [40] M. Yeo, M. T. Hummon, A. L. Collopy, B. Yan, B. Hemmerling, E. Chae, J. M. Doyle, and J. Ye, *Phys. Rev. Lett.* **114**, 223003 (2015).
- [41] D. DeMille, J. Barry, E. Edwards, E. Norrgard, and M. Steinecker, *Mol. Phys.* **111**, 1805 (2013).
- [42] A. Kazantsev and I. Krasnov, *JOSA B* **6**, 2140 (1989).
- [43] R. Grimm, Y. B. Ovchinnikov, A. Sidorov, and V. Letokhov, *Phys. Rev. Lett.* **65**, 1415 (1990).
- [44] M. Cashen and H. Metcalf, *JOSA B* **20**, 915 (2003).
- [45] J. Söding, R. Grimm, Y. B. Ovchinnikov, P. Bouyer, and C. Salomon, *Phys. Rev. Lett.* **78**, 1420 (1997).
- [46] M. Partlow, X. Miao, J. Bochmann, M. Cashen, and H. Metcalf, *Phys. Rev. Lett.* **93**, 213004 (2004).
- [47] Z. Feng, S. Ebser, L. Ringena, F. Ritterbusch, and M. K. Oberthaler, *Phys. Rev. A* **96**, 013424 (2017).
- [48] While the π -pulse model presented here is the most intuitive to understand, full description of the bichromatic force on a two-level system involves using a doubly-dressed atom model and can be found in Ref. [55].
- [49] P. I. Presunka and J. A. Coxon, *Chem. Phys.* **190**, 97 (1995).
- [50] I. Kozyryev, L. Baum, K. Matsuda, B. L. Augenbraun, L. Anderegg, A. Sedlack, and J. M. Doyle, *Phys. Rev. Lett.* **118**, 173201 (2017).
- [51] D.-T. Nguyen, T. C. Steimle, I. Kozyryev, H. Meng, and A. B. McCoy, *arXiv:1711.08033* (2017).
- [52] I. Kozyryev, L. Baum, K. Matsuda, B. Hemmerling, and J. M. Doyle, *J. Phys. B: At. Mol. Opt. Phys.* **49**, 134002 (2016).
- [53] M. Williams, F. Chi, M. Cashen, and H. Metcalf, *Phys. Rev. A* **61**, 023408 (2000).
- [54] L. Podlecki, R. Glover, J. Martin, and T. Bastin, preprint *arXiv:1702.05410* (2017).
- [55] L. Yatsenko and H. Metcalf, *Phys. Rev. A* **70**, 063402 (2004).
- [56] Ref. [30, 63] and SM provide the details of how the calculations

for SrOH were performed.

- [57] See Supplemental Material at [URL will be inserted by publisher] for dependence of BCF on the magnitude of applied magnetic field, the angle between the field and laser polarization axis as well as additional Refs. [67–69].
- [58] C. Corder, B. Arnold, and H. Metcalf, Phys. Rev. Lett. **114**, 043002 (2015).
- [59] C. Corder, B. Arnold, X. Hua, and H. Metcalf, JOSA B **32**, B75 (2015).
- [60] E. S. Shuman, J. F. Barry, D. R. Glenn, and D. DeMille, Phys. Rev. Lett. **103**, 223001 (2009).
- [61] See Ref. [63] and SM.
- [62] S. E. Galica, L. Aldridge, and E. E. Eyler, Phys. Rev. A **88**, 043418 (2013).
- [63] L. Aldridge, *The Bichromatic Force in Multi-Level Systems*, Ph.D. thesis, University of Connecticut (2016).
- [64] I. Kozyryev and N. R. Hutzler, Phys. Rev. Lett. **119**, 133002 (2017).
- [65] P. F. Bernath, Adv. Photochem. **23**, 1 (1997).
- [66] N. Balakrishnan, J. Chem. Phys. **145**, 150901 (2016).
- [67] Y. Cunyun, *Tunable external cavity diode lasers* (Hackensack: World Scientific, 2004).
- [68] D. Berkeland and M. Boshier, Phys. Rev. A **65**, 033413 (2002).
- [69] J. Nakagawa, R. F. Wormsbecher, and D. O. Harris, J. Mol. Spectrosc. **97**, 37 (1983).