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Phys. Rev. Lett. **117**, 017201 — Published 28 June 2016

DOI: [10.1103/PhysRevLett.117.017201](https://doi.org/10.1103/PhysRevLett.117.017201)

# Solitary magnons in the $S = \frac{5}{2}$ antiferromagnet- $\text{CaFe}_2\text{O}_4$

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(Dated: May 2, 2016)*

$\text{CaFe}_2\text{O}_4$  is a  $S = \frac{5}{2}$  anisotropic antiferromagnet based upon zig-zag chains having two competing magnetic structures, denoted as the A ( $\uparrow\uparrow\downarrow\downarrow$ ) and B ( $\uparrow\downarrow\uparrow\downarrow$ ) phases, which differ by the  $c$ -axis stacking of ferromagnetic stripes. We apply neutron scattering to demonstrate that the competing A and B phase order parameters results in magnetic antiphase boundaries along  $c$  which freeze on the timescale of  $\sim 1$  ns at the onset of magnetic order at 200 K. Using high resolution neutron spectroscopy, we find quantized spin wave levels and measure 9 such excitations localized in regions  $\sim 1$ -2  $c$ -axis lattice constants in size. We discuss these in the context of solitary magnons predicted to exist in anisotropic systems. The magnetic anisotropy affords both competing A+B orders as well as localization of spin excitations in a classical magnet.

Materials that display localized behavior have been studied extensively for the search of new properties including localized electronic [1] and electromagnetic [2] states. In the context of magnetism, single molecular magnets have been investigated owing to the ability to tune quantum properties as well as possible device applications [3–5] and mesoscopic magnetic structures [6] have been created to confine magnetic excitations. While localization in many of these systems is introduced through breaking up regular structures, spatial localization of energy has been known to exist in periodic structures that also contain strong nonlinear interactions. [7] Examples of such include solitary waves (solitons) or also localized breather modes. [8–10] Here we demonstrate the presence of such localized modes in a classical  $S = \frac{5}{2}$  periodic antiferromagnet where nonlinearity is introduced through magnetic anisotropy.

$\text{CaFe}_2\text{O}_4$  is a  $S = \frac{5}{2}$  antiferromagnetic based upon an orthorhombic (space group 62  $Pnma$ ) unit cell with dimensions  $a=9.922$  Å,  $b=3.017$  Å, and  $c=10.689$  Å. [11, 12] The magnetic structure [13–16] consists of two competing spin arrangements termed the A and B phases (illustrated in Fig. 1  $a - b$ ) which are distinguished by their  $c$ -axis stacking of ferromagnetic  $b$ -axis stripes. The high temperature B (Fig. 1  $a$ ) phase consists of stripes with antiferromagnetic alignment within the zig-zag chain network (denoted as  $\uparrow\downarrow\uparrow\downarrow$ ). At lower temperatures, this is replaced by the A phase (Fig. 1  $b$ ) with the spins aligned

parallel within the chain framework (denoted as  $\uparrow\uparrow\downarrow\downarrow$ ). The A ( $\uparrow\uparrow\downarrow\downarrow$ ) phase can be interpreted as the B ( $\uparrow\downarrow\uparrow\downarrow$ ) phase with an antiphase boundary along  $c$ . [14]

The competition between A ( $\uparrow\uparrow\downarrow\downarrow$ ) and B ( $\uparrow\downarrow\uparrow\downarrow$ ) phases can be motivated based on bond angles mediating superexchange interactions between the  $\text{Fe}^{3+}$  spins. The ferromagnetic alignment of the  $\text{Fe}^{3+}$  spins along the  $b$  axis originates from a superexchange interaction through a Fe-O-Fe bond angle of  $\sim 86^\circ$ . However, the interaction between chains within a zig-zag chain network is mediated by a bond angle of  $\sim 100^\circ$ . As outlined in Ref. 17 and 18 for the superexchange in the cuprates, the former bond angle is expected to be ferromagnetic while the second is marginal on the border of ferromagnetic and antiferromagnetic. Therefore, it is not clear based on bond angles alone whether the interaction between chains, within a given zig-zag chain network, is either ferromagnetic or antiferromagnetic and hence whether A ( $\uparrow\uparrow\downarrow\downarrow$ ) or B ( $\uparrow\downarrow\uparrow\downarrow$ ) is preferred. The preferred stability of the A ( $\uparrow\uparrow\downarrow\downarrow$ ) phase has been confirmed theoretically by electronic structure calculations. [19] The magnetic structure in both magnetic phases is unfrustrated unlike the case in anisotropic triangular magnets. [20]

To investigate the phase transitions and also the dynamics of these competing magnetic order parameters, we have used neutron scattering. The experiments were performed both on powders and single crystals of  $\text{CaFe}_2\text{O}_4$  grown using a mirror furnace. [21]

We first discuss the magnetic critical dynamics in  $\text{CaFe}_2\text{O}_4$ . Diffraction results are displayed in Fig. 1 *c*) illustrating the onset of  $B$  (Fig. 1 *a*) phase ordering at 200 K, which persists to low temperatures where the  $A$  (Fig. 1 *b*) phase smoothly onsets over a similar temperature range and then eventually dominates at low temperatures. We have measured the temperature dependence of both  $A$  and  $B$  ordering using cold and thermal neutrons with both giving consistent results despite differing energy resolutions. This indicates true temporal long range ordering in contrast to reports in systems where the magnetic ordering consists of slow relaxations. [22–26] We have further confirmed the magnetic structure drawn in panels (*a*) and (*b*) on powder used to prepare the single crystal and also with polarized neutrons.

Concomitant with the magnetic ordering, a gap opens in the magnetic excitation spectrum shown in Fig. 1 (*d*) which plots a series of constant  $\vec{Q}=(1,0,2)$  scans taken on RITA-2 as a function of temperature. The presence of an excitation gap is characteristic of an anisotropy term in the magnetic Hamiltonian which describes the energy needed to overcome the local alignment of the spin along the  $b$  axis. We characterize the relative size of this anisotropy in relation to the exchange energy below.

While the magnetic ordering is long range with momentum resolution limited Bragg peaks, the total ordered moment in Fig. 1 (*c*) does not reach  $gS=5\mu_B$  expected based on  $\text{Fe}^{3+}$  ( $S=\frac{5}{2}$ ) spins. The missing spectral weight can be accounted for by a momentum broadened component illustrated in Fig. 2 (*a*) which shows momentum broadened rods of magnetic scattering extending along  $L$  from energy integrating diffraction measurements (taken using DNS at FRM2). The polarization analysis allows us to conclude the diffuse scattering is magnetic in origin and associated with magnetic moments predominately aligned along  $b$ . Given that the magnetic structure also (Fig. 1) consists of spins aligned along  $b$  we conclude the origin comes from incomplete stacking of the  $A$  and  $B$  phases. The data is compared against a calculation in Fig. 2 (*b*) where we consider antiphase boundaries for both the  $A$  and  $B$  phases along  $c$  and long-range order within the  $a-b$  plane. The cross section for this model takes the following form,

$$I(\mathbf{Q}) = \Lambda \frac{(\gamma r_0)^2}{4} m^2 g^2 f^2(Q) \times \quad (1)$$

$$[\theta_A |F_A(\uparrow\downarrow\uparrow\downarrow)|^2 + \theta_B |F_B(\uparrow\uparrow\downarrow\downarrow)|^2] \times$$

$$\delta(Q_x - Q_{x,0}) \times \left( \frac{\sinh(c/\xi)}{\cosh(c/\xi) - \cos(2\pi L)} \right)$$

where  $(\gamma r_0)^2$  is 0.292 b,  $f(Q)$  is the  $\text{Fe}^{3+}$  form factor [27],  $g$  the Landé factor,  $F_A(\uparrow\downarrow\uparrow\downarrow)$  and  $F_B(\uparrow\uparrow\downarrow\downarrow)$  are the magnetic structure factors for the  $A$  and  $B$  phases and  $\theta_A$  and  $\theta_B$  are the ratios of the  $A$  and  $B$  phases measured in Fig. 1 *c*).  $\delta(Q_x - Q_{x,0})$  is the  $\delta$  function along (100) indicat-

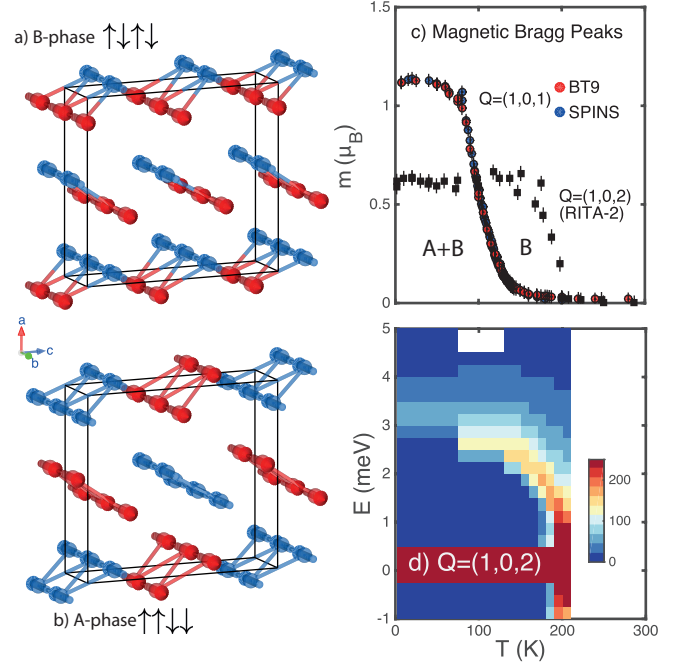


FIG. 1. (*a – b*) Schematic illustration of the refined magnetic structures for the  $A$  ( $\uparrow\downarrow\uparrow\downarrow$ ) and  $B$  ( $\uparrow\uparrow\downarrow\downarrow$ ) phases (from powder diffraction - BT1 and polarized single crystal - DNS). The lattice constants are  $a=9.21$ ,  $b=3.01$ , and  $c=10.68$  Å. Only the Fe atoms and their moments are shown for clarity. The magnetic moments are aligned along  $b$ . (*b*) plots the magnetic moments of the  $A$  and  $B$  phases as a function of temperature illustrating that both coexist at low temperatures. (*c*) Neutron spectroscopy data showing the opening of an anisotropy gap at 200 K which increases as temperature decreases.

ing long-range magnetic order along  $a$ . The parameter  $\xi=10 \pm 2$  Å indicates nearest neighbor correlations along  $c$ . The calibration constant  $\Lambda$  was calculated based on 12 nuclear Bragg peaks from which the absolute moment in this momentum broadened part of the cross section was derived to be  $m=0.40 \pm 0.10 \mu_B$ . The cross section is similar to that used to analyze short-range stripe order in the cuprates and nickelates. [23, 28] Our results indicate a significant amount of the  $\text{Fe}^{3+}$  spins are associated with regions which are disordered by antiphase boundaries.

We now investigate the dynamics of these antiphase boundaries using neutron spin-echo [29] (Fig. 2 *c, d*) which probes fluctuations on the  $\sim$  GHz timescale. The data was taken on IN11 with the multidetector setup integrating in  $\vec{Q}=(1.0 \pm 0.1, 0, 1.5 \pm 0.3)$ . Spin-echo finds that the diffuse scattering consists of both a static and dynamic component between 100-200 K, with only a static component observable at temperatures below 100 K. Example scans are shown in panel (*c*) with solid curves fits to the relaxational form  $I(Q, t)/I(Q, 0) = \alpha + (1 - \alpha)t^{1-\beta} \exp[-(\frac{t}{\tau})^\beta]$  with a temperature independent relaxation time  $\tau=1.6 \pm 0.3$  ps. [30, 31] The temperature independent characteristic timescale is consistent

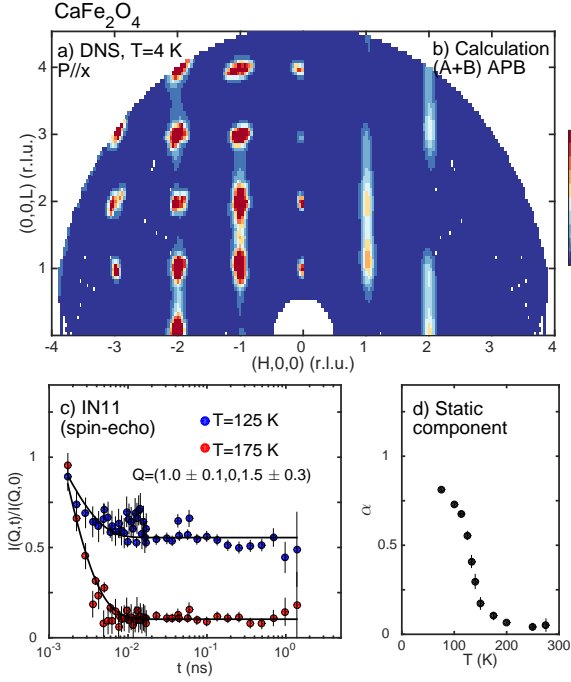


FIG. 2. (a) Energy integrating polarized diffuse scattering measurements confirming the magnetic nature of the diffuse rods extending along  $L$ . The data is compared against a calculation in panel (b) with antiphase boundaries (APB). (c) shows time dependence of the real part of the normalized intermediate scattering function at 125 and 175 K. (d) plots the static component as a function of temperature.

with slow relaxations investigated in random field Ising magnets with finite magnetic domains. [32–34] This form approaches the single relaxational lineshape when  $\beta \rightarrow 1$  and we found a temperature independent  $\beta = 0.90 \pm 0.05$  described the data well. The stretched exponential was required to fit the data and indicates a distribution of relaxation times. The static component  $\alpha$  (measured as the baseline in Fig. 2 c) as a function of temperature is shown in Fig. 2 (d) illustrating that the antiphase boundaries start to become static (on the timescale of  $\sim 1$  ns) at 200 K. At 100 K,  $\alpha = 0.75$  indicating a substantial amount of spectral weight remains dynamic even at low temperatures.

To understand the microscopic origin for the antiphase boundaries, we investigate the spatial spin coupling along the three crystallographic directions by measuring the magnetic dispersion curves (Fig. 3 at  $T = 5$  K). Figure 3 (a) shows a constant energy slice taken on MAPS illustrating well defined rings of scattering in the  $(H,K)$  plane. Dispersion curves along the  $a$  and  $b$  crystallographic axes are plotted in panels (b) and (c) illustrating strong and nearly isotropic dispersion in this plane indicative of strong exchange coupled spins

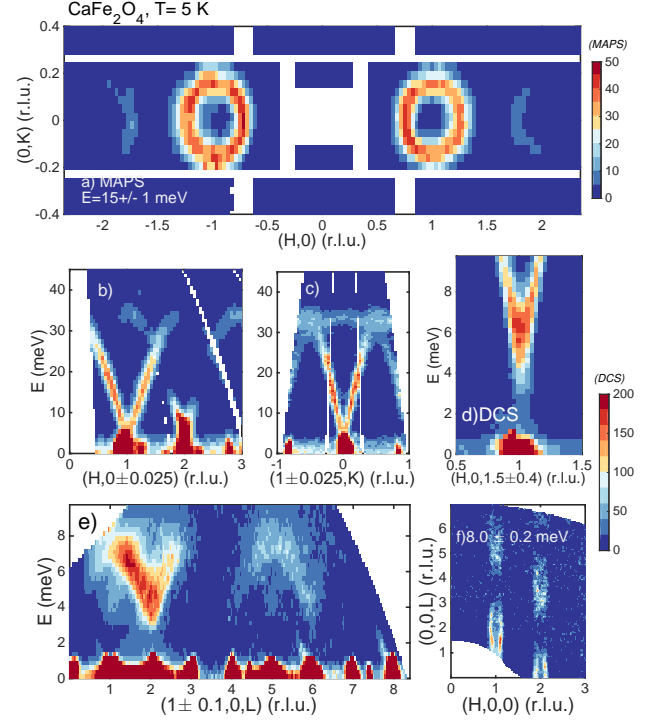


FIG. 3. (a) Constant energy slice illustrating rings of magnetic scattering in the  $(H,K)$  plane taken on MAPS. (b, c) display strong dispersion along the  $H$  and  $K$  directions from MAPS. (d) illustrates high resolution DCS data showing a low temperature anisotropy gap. (e) shows the dispersion of the excitations along  $c$  with a low energy constant energy slice illustrated in (f) in the  $(H,0,L)$  plane.

within the  $a - b$  plane. Higher resolution data (panel d) shows the presence of an anisotropy gap of  $\Delta = 5.0 \pm 0.3$  meV. Fits to dispersion relations along  $a$  and  $b$  find a spin-wave velocity of  $\hbar v = 78 \pm 5$  meV  $\text{\AA}$ . The results along the  $(0,0,L)$  direction are different and shown in panels (e) where a significantly smaller dispersion is found with a fitted  $J_c/J_{ab} = 0.14 \pm 0.03$ . This is confirmed by a constant energy slice ( $8.0 \pm 2.0$  meV) (f) which find rods of scattering following the structure factor used to model the  $A$  and  $B$  diffuse scattering above. Based on these dispersion curves, we find the spin coupling in  $\text{CaFe}_2\text{O}_4$  is strongly two dimensional with weak spin coupling along the  $(0,0,L)$  direction. The strong two dimensionality characterized by the large difference in spin exchange ( $J_c/J_{ab}$ ) facilitates the formation of antiphase boundaries along  $c$  discussed above.

We now investigate the spin excitations using the fine energy resolution afforded by neutron backscattering. High resolution temperature dependent spectroscopy of the magnetic excitations is illustrated in Fig. 4 with an elastic resolution of  $\delta E = 0.025$  meV (full-width at half maximum). At high temperatures of 235 K (panel a)

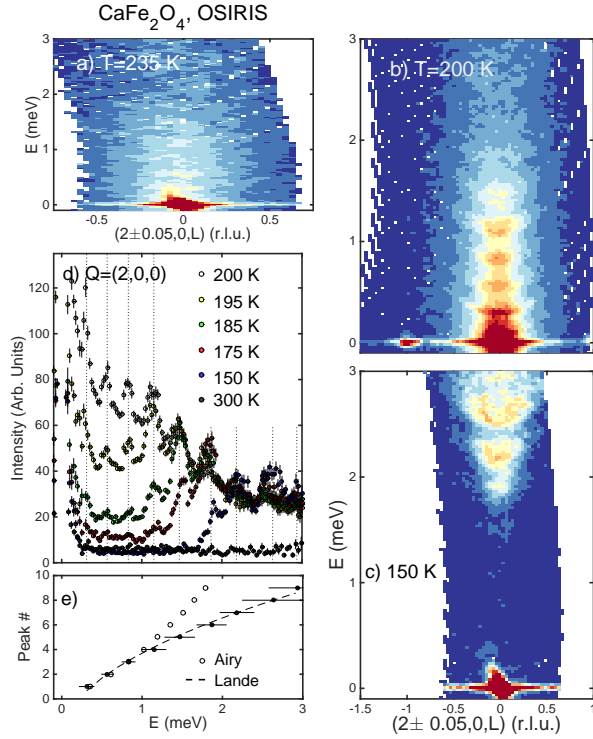


FIG. 4. High resolution spectroscopy illustrating the presence of countable magnetic excitations. (a) shows temporally and momentum broadened excitations at high temperatures. (b – c) illustrate quantized spin excitations at 200 and 150 K with a summary of a series of temperatures plotted in (d). The peak position is plotted against energy in (e) and compared against expectations based on the Airy analysis and a heuristic model based on the Landé interval theorem and also proposed by theory for anisotropic classical chains.

momentum and temporally broadened correlations are observed consistent with paramagnetic critical scattering for the formation of the  $B$  phase. At high temperatures, thermal fluctuations dominate the dynamics.

However, at 200 K (panel b), where static antiphase boundaries begin to form (Fig. 2 d), this is replaced by a series of discrete quantized excitations. The hierarchy of excitations persists for temperatures below the magnetic ordering temperature of 200 K and is shown in panel (c) at 150 K. A summary of a series of temperatures from 150 K-300 K is shown in panel d) where 9 discrete energies are observed. The spectral weight distribution as a function of temperature tracks the recovery of the anisotropy gap studied with coarser energy resolution in Fig. 1. The discrete nature is most clearly resolved in scans along  $L$  given the weak spin coupling and hence softer spin wave dispersion along this direction.

We now discuss the origin of the quantized spin waves observed below  $T_N$  shown in Fig. 4. A possible mechanism for quantized magnetic fluctuations is

through a binding interaction as postulated to exist in  $\text{CoNb}_2\text{O}_6$  [35, 36] where discrete excitations are stabilized by interactions between spinons. Our data is inconsistent with a theory describing this situation based upon a Schrödinger equation where the potential energy is proportional to the separation of the spinons (in the case of  $S=\frac{1}{2}$ ). Confirming this conclusion, Fig. 4 (e) shows a plot of the energy position of each mode compared against the negative roots of the Airy function which are the solutions of this Schrödinger equation. The data diverges markedly from the prediction at higher energy and larger peak number ( $N$ ) indicating the discrete nature does not originate from interactions.

Another possibility is that the discrete excitations originate from a localized spin object where the underlying Hamiltonian is  $H = J \sum_n \vec{S}_n \cdot \vec{S}_{n+1}$ , with the sum over a finite number of spins. Such a Hamiltonian, which is bilinear in angular momentum, is subject to the Landé interval rule which states that the interval between two neighboring levels is proportional to the higher  $j$  (quantum number associated with the total angular momentum operator) value of the pair. [37] Hence, for increasing energy and hence  $j$ , assuming antiferromagnetic interactions, the level spacing is predicted to increase in agreement with the data. In panel (e), we fit the peak position to the heuristic form  $N = \Pi\sqrt{E} + \Phi$ , where  $\Pi$  and  $\Phi$  are fitting constants,  $E$  is the energy position, and  $N$  is peak integer. This heuristic model motivated by the interval theorem provides a good description of the data over the 9 observed peaks implicating localized regions of confined classical spins as the origin of the discrete spin waves.

The discrete excitations are only observed below  $T_N=200$  K where static magnetic order and antiphase boundaries are present. An estimate of the size of these localized regions is given by the correlation length derived from the diffraction data presented above in Fig. 2 with a lengthscale of only 1-2 unit cells along  $c$ . The size of these localized regions in the  $a-b$  plane is determined by the resolution ( $\sim 100$  Å) indicating that they are highly anisotropic spatially. The observed modes are analogous to discrete breathers as they are spatially localized, time periodic, and stable excitations. However, unlike localized nonlinear solitons which can propagate without dampening, breathers are known to decay [8, 38, 39], consistent with the temperature independent relaxation on the 1 GHz timescale measured in Fig. 2.

Stabilizing localized excitations has been predicted in lattices which have nonlinear terms in the energy expansion. In the context of the antiferromagnetic Hamiltonian, this can occur through anisotropic energies originating from a distorted octahedra around the magnetic  $\text{Fe}^{3+}$  ion. An analogous case has been considered in a number of studies and most notably in a classic spin-chain [40–44]. Theoretical studies [45–48] found the semiclassical (high spin) anisotropic Hamiltonian to map onto the discrete nonlinear Schrödinger equation and to sup-

port soliton excitations where the energy scaled as  $\propto N^2$  - the result argued above based on the Landé interval rule. A similar increase in frequency spacing has also been modelled in nonlinear electronic circuits [49], lattices [50], and waveguides [51] which also are based upon similar nonlinear equations.

Localized antiferromagnetic excitations have been observed through driving large amplitude spin-waves in low-dimensional antiferromagnets. [38] Countable excitations were measured as a step in the time dependence of the emission signal. [52] The anisotropy in  $\text{CaFe}_2\text{O}_4$  allows such localized countable excitations to exist even under mild perturbations with neutrons.

In summary, we have reported the magnetic excitations and transitions in  $\text{CaFe}_2\text{O}_4$ . We find competing anisotropic order parameters which support countable spin excitations.  $\text{CaFe}_2\text{O}_4$  therefore displays solitary excitations confining energy locally.

This work was supported by the EPSRC, Carnegie Trust for the Universities of Scotland, Royal Society of London, Royal Society of Edinburgh, STFC, and EU-NMI3. Further details about open access data is given the supplementary information.

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