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Condensing magnons in a degenerate ferromagnetic spinor Bose gas

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We observe the quasi-condensation of magnon excitations within an $F = 1$ ^{87}Rb spinor Bose-Einstein condensed gas. Magnons are pumped into a ferromagnetically ordered gas, allowed to equilibrate to a non-degenerate distribution, and then cooled evaporatively at near-constant net longitudinal magnetization whereupon they condense. The critical magnon number, spatial distribution, and momentum distribution indicate that magnons condense in a potential that is uniform within the volume of the ferromagnetic condensate. The macroscopic transverse magnetization produced by the degenerate magnon gas remains inhomogeneous within the ~ 10 s equilibration time accessed in our experiment, and includes signatures of Mermin-Ho spin textures that appear as phase singularities in the magnon quasi-condensate wavefunction.

Magnons are collective spin excitations of a magnetically ordered medium. At thermal equilibrium, these bosonic quasiparticles are not conserved in number, and, thus, are not expected to undergo Bose-Einstein condensation. Nevertheless, magnon condensation has been observed in non-equilibrium systems that are pumped with an excess population of magnons that decays slowly [1–3]. Magnon condensation also describes the onset of transverse canted magnetic order in spin-dimer compounds at high magnetic fields [4, 5]. Signatures of magnon condensation include properties of the critical point [5, 6], the accumulation of magnons in low-energy states [2], and spontaneous symmetry breaking indicated by the emergence of large regions of precessing transverse magnetization [1, 3, 7].

Here, we report on magnon quasi-condensation in a spinor Bose-Einstein condensate (BEC). We use $F = 1$ gases of ^{87}Rb , in which spin-dependent s-wave interactions favor Bose-Einstein condensation in a ferromagnetic state [8]. Because the longitudinal magnetization of such a gas is a nearly conserved quantity in the absence of external perturbations, we are able to produce a long-lived excess of magnon excitations above the critical number for magnon condensation. Like in other systems, we detect magnon condensation by determining the critical point, observing the accumulation of magnons in low-energy states, and probing for the spontaneous transverse magnetization of the atomic gas.

A key finding of our work is that, at the limit of low temperature, magnon condensation occurs as the Bose-Einstein condensation of free particles in a uniform potential. Within an inhomogeneous trapping potential, a ferromagnetic spinor BEC equilibrates at a non-uniform density and a uniform chemical potential. In the absence of spin-dependent potentials and dipolar interactions [9], rotational symmetry implies that magnon excitations of the ferromagnetic spinor BEC are gapless, so that their effective potential is uniform within the condensate volume. Like free particles, the magnons disperse quadrat-

ically with wavevector, with an effective mass that is, within a weakly interacting gas, nearly equal to that of a bare atom [9, 10]. A box-like trap for atoms in a single spin state was recently constructed using a finely tuned optical and magnetic trap [11]. In contrast, the box-like potential for magnon excitations is produced naturally by the gas itself, reducing sensitivity to experimental imperfections. Such potentials yield quantum gases at nearly uniform density, allowing for precise measurements of system properties that would otherwise be smeared out by inhomogeneous broadening. The condensation of minority spins in a spinor Bose gas was previously observed at high temperature where interactions play little role in modifying the inhomogeneous trapping potential [12].

Our second finding is that following passage through the magnon condensation transition, the magnon gas remains only quasi-condensed throughout the ~ 10 s equilibration times accessible to our experiment. In this state, the transverse magnetization becomes macroscopic and homogeneous over length scales longer than the thermal de Broglie wavelength, but does not become homogeneous over the system size.

Our experiments begin with a Bose-Einstein condensed ^{87}Rb gas, prepared in the ferromagnetic states with all atoms in the $|F = 1, m_F = -1\rangle$ hyperfine state, and held in a state-independent optical trap. The gas is exposed to a 177 mG magnetic field that is uniform to about $10\ \mu\text{G}$ within the condensate. The gas temperature is controlled through evaporative cooling by the depth of the optical trap. With the gas at an initial temperature ranging from $T_i \simeq 80$ nK to 140 nK, we “pump” magnons into the gas by applying a spatially uniform, pulsed rf magnetic field at the 124 kHz Larmor precession frequency. The pulse tips the atomic spin by a small angle (up to 0.7 rad) at which the population of spin-flipped atoms – the magnon excitations – is dominantly in the $|F = 1, m_F = 0\rangle$ state. The gas is then allowed to thermalize at constant magnetization and at constant trap depth for 2.5 s. By measuring momentum distribution of the $m_F = 0$ gas, we

confirm that the magnon gas at this initial stage is not condensed.

We then cool this magnon-imbued gas by lowering the optical trap depth over a time $t_{\text{ramp}} = 5$ s and then maintain a constant depth for $t_{\text{hold}} = 2.5$ s so that the gas reaches a steady state temperature T_f before the gas is probed. During these times, not only does the total number of trapped atoms drop, but also the fractional population of magnon excitations drops owing to the preferential evaporation of magnons from the trap [13]. Nevertheless, provided that the initial magnon number is high enough, the magnon population at the time of probing can be sufficient to cross the magnon condensation transition. The coldest samples studied in this work are prepared at $T_f = 30$ nK in a trap with depth $k_B \times 122$ nK and trap frequencies $\omega_{x,y,z} = 2\pi\{23, 8.8, 160\} \text{ s}^{-1}$, with final magnon populations varying between 4% and 10% of the 6.5×10^5 atoms in the $m_F = -1$ condensate.

The properties of the magnon gas at the condensation transition are strongly influenced by the effective potential in which the magnons propagate. This effective potential $V_{\text{eff}}(\mathbf{r})$ is the sum of the external trapping potential $V(\mathbf{r})$ (with $V = 0$ at the trap center) and the interaction energy $V_{\text{int}}(\mathbf{r})$ of an $m_F = 0$ atom within the $m_F = -1$ condensate (Fig. 1a). According to mean-field theory, $V_{\text{int}}(\mathbf{r}) = \mu_{-1} - V(\mathbf{r})$ within the condensate volume and zero otherwise, with μ_{-1} the chemical potential of the $m_F = -1$ gas. We derive this expression using the Thomas-Fermi approximation and noting the equality $a_{-1,-1} = a_{-1,0}$ implied by rotational symmetry of the contact interaction; here, $a_{i,j}$ is the s-wave scattering length for a collision between atoms in the $|m_F = \{i, j\}\rangle$ states. The effective potential is then $V_{\text{eff}}(\mathbf{r}) = \max(V(\mathbf{r}), \mu_{-1})$ [14].

The first signature of the box-like potential for magnon excitations is the position-space distribution of the non-degenerate magnon gas. The normal magnon density is expected to be constant within the volume of the ferromagnetic condensate, and then, at low temperature (defined as $k_B T \ll \mu_{-1}$), to diminish rapidly outside that volume. We image this distribution within our coldest samples, for which $k_B T_f / \mu_{-1} \sim 0.7$, by applying a microwave pulse that drives atoms selectively from the $|F = 1, m_F = 0\rangle$ state to the $|F = 2\rangle$ hyperfine state, and then imaging selectively the $F = 2$ atoms with resonant probe light. The observed column density of the magnon gas \tilde{n}_0 , shown in Fig. 1b, is indeed large within the area defined by the Thomas Fermi radii of the condensate, $R_{x,y} = (2\mu_{-1}/m\omega_{x,y}^2)^{1/2} = \{21, 54\} \mu\text{m}$ in the imaged directions, and is described quantitatively by the function $\tilde{n}_0 \propto \max(1 - x^2/R_x^2 - y^2/R_y^2, 0)^\beta$ with $\beta = 1/2$. In contrast, the non-uniform density of the majority-spin condensate leads to a column density with a different exponent, $\beta = 3/2$. Both distributions deviate from that of a harmonically trapped Bose gas, $\propto g_2(\exp[-x^2/\sigma_x^2 - y^2/\sigma_y^2])$ with g_α being the polyloga-

rithm function of order α and $\sigma_{x,y} = (k_B T_f / m\omega_{x,y}^2)^{1/2} = \{12, 31\} \mu\text{m}$.

The second manner in which the box-like potential for magnon condensation is evident is the variation of the critical magnon number with temperature. At each temperature T_f , ranging from 30 to 114 nK, we examine the momentum distribution of the $m_F = 0$ atoms as a function of total magnon number. This momentum distribution is measured by releasing all atoms from their trap, and then using state-selective magnetic-field focusing and absorption imaging [13]. Excluding data from the central region of the image, we fit the momentum-space column density to several parametrizations of the distribution of a non-condensed Bose gas at variable magnon chemical potential and temperature T_f . These parametrizations include $g_\alpha(z(\mathbf{p}))$ with α chosen among several values [11], which describes the expected momentum distribution of bosons in various power-law potentials, and also the distribution expected for bosons trapped in the effective potential V_{eff} ; here, $z(\mathbf{p}) = \exp[(\mu - p^2/2m)/k_B T_f]$ with $\mu = \mu_0 - \mu_{-1}$ being the magnon chemical potential referenced to μ_{-1} and ignoring the Zeeman energy ($\mu = 0$ at the magnon condensation transition), and \mathbf{p} being the momentum in the imaged plane. The size of the excluded region is chosen such that these functions give about the same temperature. The magnon condensate number is then determined by subtracting the fitted function from the observed distribution and summing over the image, including the central region. This magnon condensate number rises linearly from zero with increasing total magnon number (Fig. 2c inset) above a number $N_{\text{mag,c}}$ that we identify as the critical magnon number for condensation. Different parametrizations of the non-condensed magnon distribution give critical magnon numbers that vary by around 10% at the same T_f for our coldest samples.

In Fig. 2, we compare $N_{\text{mag,c}}$ to predicted values for ideal-gas Bose-Einstein condensation in one of three potentials: a harmonic trapping potential, a hard-walled box potential with a volume matching that of the $m_F = -1$ condensate, and the effective trapping potential V_{eff} , using the experimentally determined $\omega_{x,y,z}$, μ_{-1} , and T_f at each setting [15]. The measurements agree well with predictions based on the condensation of magnons in the effective potential V_{eff} . At low temperature, the measured critical magnon number tends toward that predicted for a hard-walled box potential while at high temperatures ($k_B T_f \gg \mu_{-1}$), $N_{\text{mag,c}}$ tends toward the prediction for a harmonically trapped gas, as expected.

The details of the critical magnon momentum distribution serve as a third signature of condensation in a box-like potential. We assume that the calculations based on magnons propagating in the effective potential give the correct critical magnon number, and use this number to identify the momentum-space image of the magnon gas just below the magnon condensation transition (Fig.

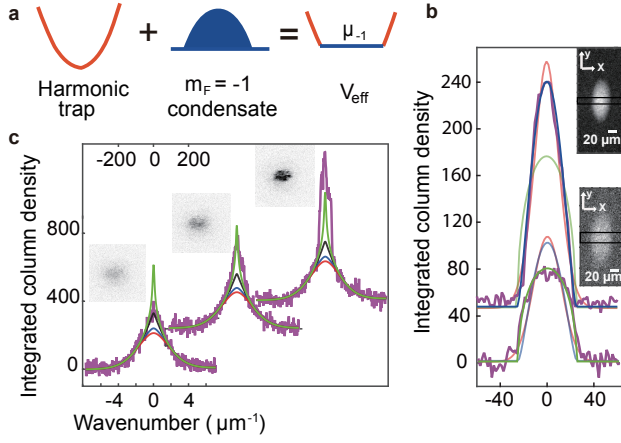


FIG. 1. Magnons in a flat potential. (a) The effective potential V_{eff} for magnon excitations is the sum of the harmonic trap potential (red) and the repulsive potential from the inhomogeneous ferromagnetic condensate (blue, maximum potential of μ_{-1}), shown schematically in one dimension. (b) The column density of a low temperature non-degenerate magnon gas (bottom column density image) is large within the volume of the $m_F = -1$ condensate (top column density image). Line densities (integrated within black rectangles) are fitted to predicted critical density in a harmonic trap (red), constant density within condensate volume (green, best fit for magnon distribution), or Thomas-Fermi condensate distribution (blue, best fit for the $m_F = -1$ condensate). (c) Images and integrated line profiles (purple) of magnon momentum distribution below (left), at (center) and above (right) critical number for magnon condensation. Profiles are fitted to Bose distributions $\propto g_\alpha(z(p))$ with $z(p)$ defined in the text and $\alpha = \{1, 1.5, 2\}$ (black, blue, red lines, respectively), and the distribution expected in V_{eff} (green).

1b). This distribution is more sharply peaked at low momentum than predicted for atoms expanding from a harmonic potential. The data are consistent with the expected momentum-space distribution of magnons in the effective potential, although we do not observe the expected sharp cusp at zero momentum ($\alpha = 1/2$ for a uniform potential), likely owing to limited momentum resolution.

Magnon condensation is a symmetry breaking phase transition. In our spinor gas system, prior to magnon condensation, the longitudinally magnetized spinor Bose-Einstein condensate retains the $O(2)$ spin-gauge symmetry that describes the combination of rotations about the longitudinal (magnetic field) axis and the multiplication of the condensate order parameter by a phase. Upon magnon condensation, this symmetry is broken, the ferromagnet now acquiring a non-zero transverse magnetization that serves as the order parameter for the magnon condensate.

We detect this transverse magnetization using absorptive state-selective *in-situ* imaging [9]. We apply an rf pulse to rotate the atomic spins by $\pi/2$ and then mea-

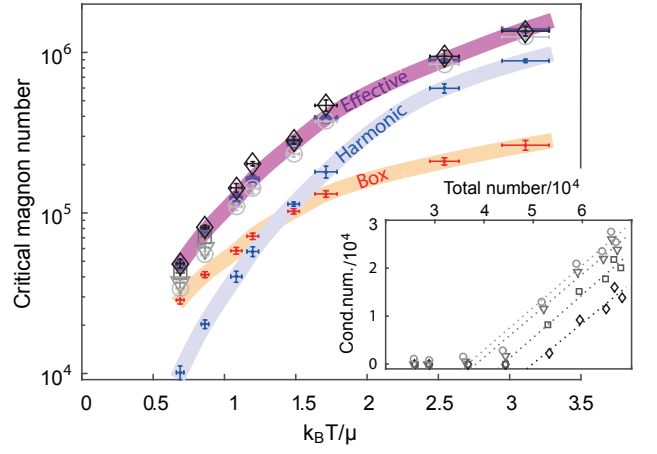


FIG. 2. The measured critical magnon number $N_{\text{mag},c}$ (black markers) is compared to the expected value (blue, red, purple points, respectively) for magnons in a harmonic, box (size defined by condensate volume), and V_{eff} . Shaded lines are guides to the eye. Error bars on model predictions account for statistical uncertainties in experimental parameters. Inset: The magnon condensate number, shown for $k_B T_i / \mu_{-1} = 0.7$, is the integrated residual above the fits to the non-condensed magnon momentum distributions using one of four fit functions: $g_\alpha(z(p))$ with $\alpha = 1$ (squares), 1.5 (triangles), or 2 (circles), or the predicted distribution in the effective potential (diamonds). A linear fit shows the magnon condensate number increasing linearly above $N_{\text{mag},c}$; error bars are statistical. Extracted measured critical magnon numbers from the four fitting functions are shown in the main plot (diamonds, squares, triangles, circles with a decrease in grayscale).

sure the column density in each of the Zeeman sublevels using a sequence of brief microwave and optical probe pulses. The difference in the $m_F = \pm 1$ column densities gives one component of the dimensionless transverse column magnetization \tilde{M}_x at the start of the several-ms-long measurement procedure. We then apply a second, carefully timed $\pi/2$ spin rotation, again measure the Zeeman state distribution, and subtract the $m_F = \pm 1$ images to obtain \tilde{M}_y . We intersperse spin-echo pulses within the measurement sequence to control for magnetic field variations, similar to Ref. [16].

Assuming the majority spin condensate order parameter $\psi_{-1}(\mathbf{r})$ is known, the magnon condensate wavefunction $\psi_0(\mathbf{r})$ is characterized interferometrically by the transverse magnetization through the relation $M_T(\mathbf{r}) = M_x(\mathbf{r}) + iM_y(\mathbf{r}) = \sqrt{2}\psi_{-1}^*(\mathbf{r})\psi_0(\mathbf{r})$. This relation is valid when the longitudinal magnetization of the gas is sufficiently large. We measure instead the transverse column magnetization, \tilde{M}_T , which can be taken as $\tilde{M}_T(\boldsymbol{\rho}) = e^{i\varphi(\boldsymbol{\rho})} \sqrt{2\tilde{n}_{-1,c}(\boldsymbol{\rho})\tilde{n}_{0,c}(\boldsymbol{\rho})}$ if we assume the atomic spin state of the condensed atoms is constant along the imaging axis. Here, $\tilde{n}_{m_F,c}$ is the column density of the m_F component of the condensate, φ is the magnon condensate phase (up to a uniform offset), and $\boldsymbol{\rho}$ is the position

in the imaged plane. We assume that the majority-spin condensate has uniform phase given that it is prepared by very gradual evaporative cooling, and that it contains no vortex excitations as confirmed by high-resolution imaging after a short time of flight from the optical trap.

The images taken on samples with sufficiently large magnon populations (Fig. 3a) show significant transverse magnetization within the boundary of the ferromagnetic condensate. The magnon condensate number implied by these images, taken as the integral of $|\tilde{M}_T|^2/2\tilde{n}_{-1,c}$ over the condensate area, is consistent with the magnon condensate numbers determined from momentum-space measurements [17].

This transverse magnetization is inhomogeneous, both in magnitude and in phase, so that the degenerate magnon gas is more aptly described as a quasi-condensate [18]. The spatial Fourier power spectra of such images (Fig. 3b) are concentrated in a narrow and nearly isotropic ring with spatial wavenumber k_r [19]. The indicated order-parameter domain size $\xi = k_r/(2\pi) \sim 17 \mu\text{m}$ is similar to that observed in previous studies of non-equilibrium ^{87}Rb spinor Bose-Einstein condensates [16, 20, 21]. These domains, and the associated ring in Fourier space, are observed only for magnon numbers above the critical number; images for samples with fewer magnons show no significant transverse magnetization.

The images of the magnon quasi-condensate include singularities, regions encircled by a path along which the amplitude of the transverse magnetization is non-zero while the magnon condensate phase φ winds around by $\pm 2\pi$. We observe several such singularities in each repetition of the experiment, with the sign of the phase winding varying randomly. We identify these features as Mermin-Ho spin textures [22–24], in which the orientation of the magnetization spans a small cap about the longitudinal axis.

The inhomogeneity of the magnon quasi-condensate may be the result of the gas having been quenched rapidly across the magnon condensation transition. This transition occurs simultaneously throughout the uniform effective potential experienced by the magnons. The characteristic domain size ξ in our measurements does not agree directly with the Kibble-Zurek theory [25, 26] in two regards. First, the Kibble-Zurek theory predicts a domain size, just after the quench, on the order of the thermal deBroglie wavelength $\lambda_{\text{dB}} = (2\pi\hbar^2/mk_B T_f)^{1/2} = 1.1 \mu\text{m}$, which is far smaller than observed. Second, we imaged magnon condensates produced with variable ramp times t_{ramp} , tuning the quench time $t_Q = T_c/(dT/dt)$ between 3 and 18 s, and observed no significant change in the characteristic domain size. In contrast, Kibble-Zurek theory predicts a variation by a factor of 1.8 [27–29].

The discrepancy with Kibble-Zurek theory can be explained by the coarsening of the magnon condensate order parameter following a quench. By varying t_{hold} between 3.5 and 11 s with a fixed $t_{\text{ramp}} = 5$ s, we observe

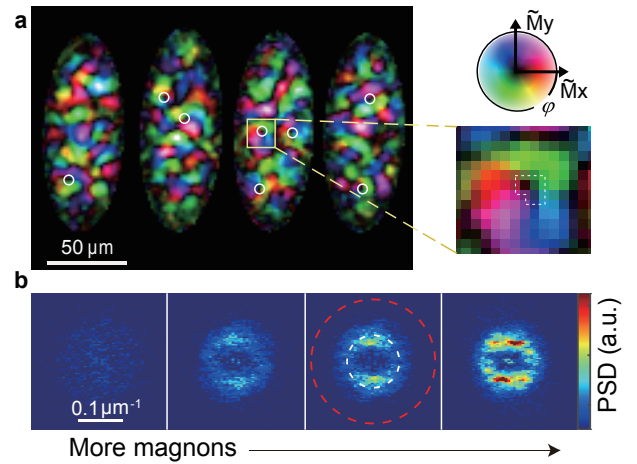


FIG. 3. Inhomogeneous transverse magnetization measured for four experimental repetitions, in coldest samples and with $\sim 10\%$ magnon fraction, with hue and brightness representing orientation and amplitude, as indicated by color wheel. Identified phase singularities are marked by white circles. The inset highlights one such singularity, with the dashed white line showing a path along which the magnon order parameter phase winds by 2π . (b) Spatial Fourier power spectral density (PSD) of transverse magnetization averaged over 20 experimental repetitions, for magnon fraction increasing from 7% (below critical value) to 10% (above critical value). White dashed circle marks $2\pi/k_r$. Red dashed circle represents the imaging resolution.

only slight coarsening that increases ξ from 16 to $22 \mu\text{m}$, similar to observations made in Ref. [16]. However, it may be that the magnon condensate is initially disordered in three dimensions, with characteristic domain size on the order of λ_{dB} , and then rapidly coarsens until the correlation length becomes a few times the length of the $m_F = -1$ condensate in its narrowest (vertical) dimension, at which point our images, which are column-integrated and limited to an imaging resolution of around $7 \mu\text{m}$, can finally reveal significant transverse magnetization; for this length, we may take the condensate Thomas-Fermi diameter of $6 \mu\text{m}$ at the low-temperature trap setting. Upon the dimensional crossover to a two-dimensional system, coarsening dynamics might slow significantly.

Magnon condensation in a spinor gas offers a new system in which to study the Bose-Einstein condensation of quasi-particle excitations and of quantum gases in a uniform potential. Our measurements of the critical magnon number, density, and momentum distribution agree with predictions based on equilibrium thermodynamics of particles held in a flat-bottomed effective potential. However, detailed measurements of the non-equilibrium dynamics of the magnon quasi-condensate present challenges for further study. Measurements of the real-time evolution of the gas magnetization, using re-

peated non-destructive spin-sensitive imaging [9, 30] with finer spatial resolution than achieved here, may address these challenges.

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- [18] For uniform magnon condensate, we expect a homogeneous transverse magnetization. Corresponding spatial Fourier power spectra should be peaked at wavenumber $k=0$. We experimentally exclude the existence of uniform condensate. See Supplemental Material at [URL will be inserted by publisher].
- [19] By imaging deliberately prepared helical spin textures of varying wavevector, we confirm that k_r lies within the resolution of our spin-sensitive imaging system; see Supplemental Material at [URL will be inserted by publisher] for details.
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