

This is the accepted manuscript made available via CHORUS. The article has been published as:

Tuning Bulk and Surface Conduction in the Proposed Topological Kondo Insulator SmB_{6}

Paul Syers, Dohun Kim, Michael S. Fuhrer, and Johnpierre Paglione

Phys. Rev. Lett. **114**, 096601 — Published 3 March 2015

DOI: [10.1103/PhysRevLett.114.096601](https://doi.org/10.1103/PhysRevLett.114.096601)

Tuning bulk and surface conduction in topological Kondo insulator SmB_6

Paul Syers*,¹ Dohun Kim*,^{1,2} Michael S. Fuhrer,^{1,3} and Johnpierre Paglione^{1,4}

¹*Center for Nanophysics and Advanced Materials, Department of Physics,
University of Maryland, College Park, MD 20742, USA**

²*Department of Physics, University of Wisconsin, Madison, WI 53706, USA**

³*School of Physics, Monash University, Victoria 3800, Australia*

⁴*Canadian Institute for Advanced Research, Toronto, Canada M5G 1Z8†*

Bulk and surface state contributions to the electrical resistance of single-crystal samples of the topological Kondo insulator compound SmB_6 are investigated as a function of crystal thickness and surface charge density, the latter tuned by ionic liquid gating with electrodes patterned in a Corbino disk geometry on a single (100) surface. By separately tuning bulk and surface conduction channels, we show conclusive evidence for a model with an insulating bulk and metallic surface states, with a crossover temperature that depends solely on the relative contributions of each conduction channel. The surface conductance, on the order of $100 e^2/h$, exhibits a field-effect mobility of $133 \text{ cm}^2/\text{Vs}$ and a large carrier density of $\sim 2 \times 10^{14} \text{ cm}^{-2}$, in good agreement with recent photoemission results. With the ability to gate-modulate surface conduction by more than 25%, this approach provides promise for both fundamental and applied studies of gate-tuned devices structured on bulk crystal samples.

Recent theoretical work has proposed that the intermediate valence compound SmB_6 may be a member of a newly classified family of strong topological insulators [1–3]. Called topological Kondo insulators, these systems differ from the conventional family of topological insulators [4, 5] such as Bi_2Se_3 because the bulk insulating band gap arises due to electronic correlations and opens at the Fermi energy. These materials are extremely interesting because of the potential for interplay between the topological states and other correlated electronic states, as well as the possibility to alleviate issues with chemical potential shifts due to intrinsic bulk doping [6, 7].

SmB_6 , one of the first known Kondo insulator materials, has been of interest for many decades due to a long debate about the nature of its insulating state [8, 9]. It is now well known to harbor a d - f hybridization gap that opens at low temperatures and has been well characterized by several experimental techniques to lie in the range of ~ 10 – 20 meV [10–20]. The electrical resistance $R(T)$ of SmB_6 exhibits a thermally activated behavior at intermediate temperatures below room temperature, before saturating at an approximately temperature-independent value below a few degrees Kelvin [8, 12, 14, 21–27]. This robust property has recently been considered a key signature of topologically protected surface states [1, 26], prompting many experimental efforts designed to probe the nature of the conducting states in this material [14, 16, 18–20, 28–32]. Most recently, strong evidence confirming the topologically protected nature of these states has been mounting [28, 33, 34].

Here we present resistance measurements probing the nature of surface conduction in bulk SmB_6 samples, using variations of bulk crystal geometry and surface ionic liquid gating techniques to, respectively, tune the bulk and surface conductance contributions. In both cases, $R(T)$ is well described by a low-temperature, temperature-

independent surface contribution in parallel with a thermally activated bulk contribution, with a crossover temperature that depends on the relative values of each conductance component. Gate-tuned measurements using a Corbino contact geometry placed on the (100) surface indicate a very large surface carrier density that can be dramatically changed by application of bias voltage. Our results strongly support the model of an insulating bulk with metallic surface states, as previously probed by other techniques [14, 25, 26], and characterize the tunability, mobility and carrier density of surface charge carriers, in good agreement with other spectroscopic techniques. Our study not only confirms the ability to tune the relative surface and bulk conductance contributions, but also paves the way for unique gate-controlled device construction on single-crystal samples of SmB_6 .

Single crystals were grown using polycrystalline SmB_6 as the reactant and Al as the flux in a ratio of 1:200. Starting materials were placed in an alumina crucible and sealed in a quartz ampoule under partial Ar pressure. Ampoules were heated to 1250°C and maintained at that temperature for 120 hours, then cooled at -2°C/hr to 900°C , followed by faster cooling. Crystals were etched out of the flux using HCl, yielding mostly cubic-shaped crystals ranging in size from $\sim (0.2)^3 \text{ mm}^3$ to $\sim (1.2)^3 \text{ mm}^3$. Single-crystal x-ray diffraction at 200 K using a Bruker APEX-II CCD system with a graphite monochromator and a $\text{MoK}\alpha$ source yields excellent refinement of crystallographic parameters, with lattice constants $4.13308(8) \text{ \AA}$ for the Pm-3m cubic structure and goodness of fit convergence at $R_1 = 0.62\%$. All samples were sanded and polished prior to contact placement. Sample thickness was controlled by means of sanding and measured using an optical microscope, with uncertainties dominated by magnification resolution. Thickness-dependent electrical resistivity measurements

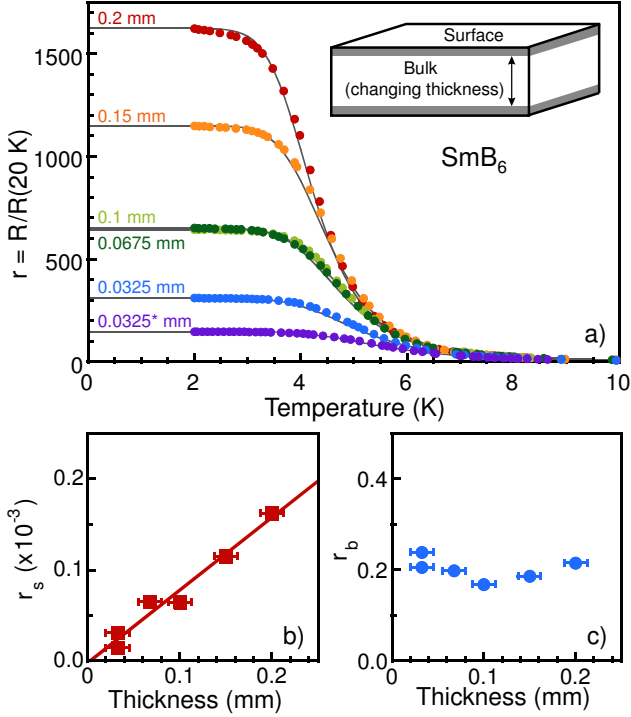


FIG. 1. a) Electrical resistance of a single crystal of SmB_6 as a function of sample thickness, normalized to its value at 20 K. The solid lines represent fits to the data using a two-channel conductance model [Eq. (1)], with fit parameters including the surface (panel b) and bulk (panel c) resistance ratio components $r_s \equiv R_s/R(20 \text{ K})$ and $r_b \equiv R_b/R(20 \text{ K})$, respectively, of the total conductance. [(*) Note geometry of sample for each 0.0325 mm thickness data set is slightly different due to loss of sample fraction.]

were performed using the standard AC technique, with four-wire geometry gold contact wires attached with silver conducting paint to the (110) face of the crystal. Gating experiments were also performed on the (100) face of the crystal, using a four-wire Corbino contact geometry pattern metalized with Au (200nm)/Ti (10nm) using thermal evaporation. After mounting on an insulating substrate, samples were covered with ionic liquid N,N-diethyl-N-methyl-N-(2-methoxyethyl)ammonium bis(trifluoromethylsulfonyl)imide (DEME-TFSI, Kanto Corporation) and an adjacent Au pad was used as a gate electrode. Ionic liquid application was performed inside a glove box and the sample was then transferred to the measurement cryostat within 5 minutes to minimize electrochemical reaction of DEME-TFSI with the ambient atmosphere. After each temperature sweep, gate voltage modulation was done by warming the sample to 230 K in-situ and changing gate voltage.

Fig. 1a) presents the temperature dependence of the longitudinal resistance of a single crystal as a function of thickness variation, with measurements taken subsequent to each thickness adjustment. To eliminate uncertainty in the geometric factor arising from varying contact ge-

ometry, we plot the resistance ratio $r \equiv R/R(20 \text{ K})$, or $R(T)$ normalized to the resistance value at 20 K (approximately the temperature at which the Kondo gap is fully open). The resistance curves exhibit similar qualitative behavior to each other and to those reported in the literature [11, 21, 25, 35] over the entire temperature range up to 300 K, with 2 K resistivity values ranging between 0.5–2.9 Ωcm .

The crossover from high-temperature, thermally activated behavior to a low-temperature plateau in $R(T)$ has been interpreted as a transition from bulk state-dominated conduction to surface state-dominated conduction [25, 26, 35]. This picture is consistent with the thickness dependence of resistance presented in Fig. 1a), which exhibits a clear separation of $r(T)$ curves from a single trace at higher temperatures to distinct plateau values for each thickness at low temperatures. In other words, the relative bulk-to-surface ratio of conductance shrinks with decreasing thickness, as expected due to the reduction of overall bulk conductance.

A simple parallel conductance model is used to extract the relative contributions, with total conductance described by $G = G_s + G_b$, where $G_s = 1/R_s$ is the surface contribution (assumed temperature-independent) and $G_b = 1/R_b$ is the bulk contribution, assumed to be activated in temperature due to a bulk energy gap Δ . Therefore $G_b = W \cdot t / (\rho_b L) e^{-\Delta/k_B T}$ with sample length L , width W , and thickness t ; bulk resistivity ρ_b in the high-temperature limit; and Boltzmann constant k_B . Thus, for the dimensionless and geometry-independent [36] resistance ratio,

$$r(T)^{-1} = r_s^{-1} + [r_b e^{-\Delta/k_B T}]^{-1}, \quad (1)$$

where $r_s \equiv R_s/R(20 \text{ K})$ and $r_b \equiv R_b/R(20 \text{ K})$ are the dimensionless, normalized surface and bulk resistance ratios, respectively.

Fits to this model using r_s , r_b and Δ as free parameters are shown as solid lines in Fig. 1a). For all thicknesses, we obtain a thickness-independent energy gap of $\Delta = 3.3 \pm 0.2 \text{ meV}$, consistent with other transport measurements [8, 12, 21, 24, 25]. The values of r_s and r_b are presented in Figs. 1b) and c), respectively, showing a clear contrast in their relationship with thickness; r_s exhibits a clear linear trend with thickness, while r_b is independent of thickness. Understood in the context of their normalized nature, the linear relation of $r_s(t)$ translates to a linearly decreasing relative contribution of surface conductance compared to bulk conductance with increasing sample thickness. Conversely, the extrapolated value of $r_s(0) = 0 \pm 0.00001$ at the zero-thickness limit translates to zero electrical conductance through the bulk, as expected in the bulk-surface model at low temperatures.

Previous experiments with electrochemical gating of Bi_2Se_3 thin films have shown great success in shifting the Fermi energy from well within the bulk conducting band into the bulk gap, allowing the isolated Dirac surface

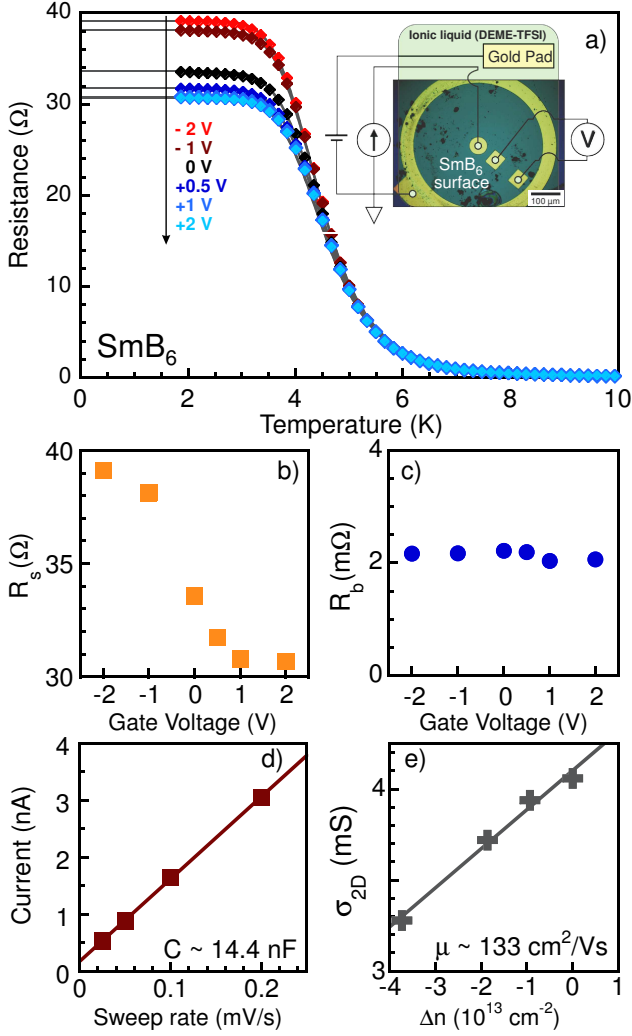


FIG. 2. a) Electrical resistance of SmB₆ sample as a function of ionic liquid gate voltage, measured with a Corbino disk lead geometry placed on the surface of a single crystal, as shown in the inset diagram (superimposed on sample photo, with green shaded area representing the area covered by the ionic liquid gate structure). Solid lines are fits to the two-channel conductance model [Eq. (2)]. Panels b) and c) present the surface and bulk resistance contributions, respectively, extracted from the two-channel fits as a function of gate voltage. Panel d) plots the transient current between the gate pad and ground as a function of gate voltage sweep rate, measured at 230 K. The solid line is a fit to a charging capacitor model with a capacitance of 14.4 nF. Panel e) shows the two-dimensional sheet conductivity as a function of the change in carrier concentration induced by gating. The solid line is a fit to a constant-mobility model with a mobility $\mu = 133 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$.

states to be probed directly [37–39]. Gating is, however, usually only effective at shifting the chemical potential of thin films or two dimensional systems, not bulk materials, as the gate electric field is confined to a thin region near the surface and heavily screened by bulk charge car-

riers. In the case of SmB₆, applying a gate voltage to the surface of a bulk crystal is a simple, yet clear test of the surface versus bulk contribution of charge carriers. If the transport is dominated by surface conduction [40], the Corbino geometry ensures that the electronic transport occurs only on the surface of a single side of the sample.

For the IL gating measurements, a four-probe Corbino geometry (see inset of Fig. 2) was patterned using e-beam lithography on a polished (100) surface of SmB₆. Fig. 2a) presents the $R(T)$ data for a single corbino device with various values of applied gate voltage V_g . Similar to the case of thickness variation (c.f. Fig. 1), the variation of V_g has no effect on $R(T)$ at higher temperatures, as exhibited by the collapse of all data onto a single trace above ~ 5 K. However, at lower temperatures a clear voltage-dependent splitting of $R(T)$ occurs, suggesting an identical tuning of bulk-to-surface contributions to the measured conductance, now controlled by a gate-controlled shift of the surface state chemical potential.

The same two-conductivity model can be applied, with the exception that resistance ratios are no longer needed since no geometries are varying. We therefore fit $R(T)$ to the following form:

$$R(T)^{-1} = R_s^{-1} + [R_b e^{-\Delta/k_B T}]^{-1} \quad (2)$$

where R_s is the (constant) surface resistance and R_b the bulk resistance in the high temperature limit, and Δ is the gap energy as before.

Similar to the thickness case, we obtain a voltage-independent value of $\Delta = 3.78 \pm 0.01$ meV. Presented in Figs. 2b) and c) are the results for R_s and R_b , respectively, as a function of V_g , showing that the variation of V_g has a dramatic effect on the surface resistance R_s , modulating it by over 25% through the accessible voltage range, while the bulk resistance R_b remains unaffected and relatively constant. The large change in saturation resistance with IL gating (Fig. 2b) but negligible change in bulk resistivity (Fig. 2c) confirms that the additional carriers injected by IL gating are confined to a region very near the surface. This tuning directly confirms the surface-born origin of low-temperature charge carriers in SmB₆ and demonstrates the unique ability of controlling surface state conduction via device construction on the surface of a bulk crystal. In contrast, bulk doping [41, 42] is known to suppress the hybridization gap, thus reducing the bulk resistivity, whereas we see no change in our analysis of the bulk resistivity with IL gate tuning. We can also rule out that the IL gating effect is due to changes in bulk doping in a layer near the surface, which would manifest as a change in activation energy and R_b , not seen in our experiment.

The Corbino gating experiment also provides information on the sign of charge carriers, their areal density and their mobility. As shown in Fig. 2a), the variation of R_s with V_g is a decreasing function, consistent

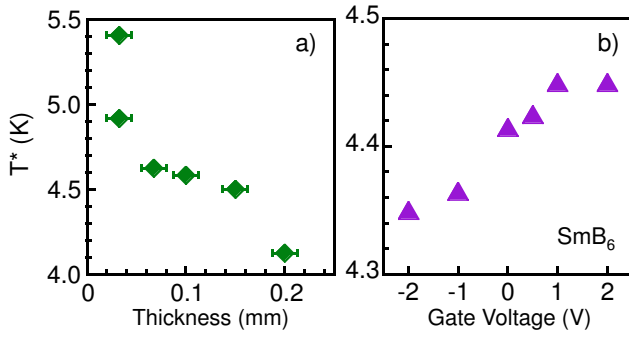


FIG. 3. a) Evolution of the crossover temperature T^* between bulk- and surface-dominated conductance, defined as the inflection point in resistance temperature dependence as a function of (a) crystal thickness variation and (b) ionic liquid gate voltage tuning. The non-constant evolution of T^* is consistent with the two-channel conduction model (see text), showing the crossover depends solely on the relative bulk and surface contributions to the total electrical conduction.

with the presence of dominant electron-like charge carriers at the surface. This may seem to be at odds with some measurements [26, 43] and consistent with others [8, 18, 19, 25, 29, 44], but it should be noted that traditional Hall effect experiments are no longer trivial in a situation with surface dominated transport, where conduction may be non-uniform. Recent angle-resolved photoemission measurements on SmB_6 have observed two large and one small electron-like surface state bands, centered at the Γ and X points, respectively, in addition to the expected hole-like bulk bands [18, 19, 25, 29], suggesting a large difference in mobilities of each carrier type may be consistent with our observations.

Using a simple capacitor model to fit the transient gate current dependence on V_g sweep rate (Fig. 2d)), allows the determination of the gate capacitance (14 nF), or specific capacitance $c_g = 3 \mu\text{F}/\text{cm}^2$, and hence the change in surface carrier density $\Delta n = c_g V_g / e$, where e is the elementary charge. The measured sheet conductivity σ_{2D} is approximately linear in the gate-induced change in Δn , as shown in Fig. 2e), indicating a constant field-effect and a surface carrier mobility of $133 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ (based on the measured resistance and the distance between voltage probe contacts – see Fig. 2 inset). Extrapolation of the linear relationship to $\sigma_{2D} = 0$ provides an estimate of $n \approx 2 \times 10^{14} \text{ cm}^{-2}$ for the total carrier concentration. Using the unit cell area of $(4.13 \text{ \AA})^2$ this amounts to roughly one third of an electron per unit cell of the surface, which indicates the unlikelihood that the surface state arises due to impurities on the surface.

While the uncertainty of $\sim 30\%$ for μ and n is large due to difficulties in estimating sample area and geometric factor for the Corbino geometry, the absolute carrier density is in excellent agreement with recent photoemission results [45]. This agreement should be considered in the context of two very different surface preparation tech-

niques (*i.e.*, UHV-cleaved versus oxidized), which may result in non-trivial energy band shifts that render such agreement fortuitous. The large carrier density suggests, however, that relatively large Fermi surface pockets make a significant contribution to the overall conduction. The relatively low surface state mobility then appears to be a natural consequence of the very low Fermi velocity for these surface states, which could possibly explain the difficulty in observing Shubnikov de Haas oscillations[28].

Finally, in both thickness and gating variation experiments above, the crossover from bulk- to surface-dominated conductance with decreasing temperature is observed to change as a function of the control parameter, as expected due to the change in relative weighting of each contribution. Fig. 3a) and b) present variation in this crossover temperature T^* , defined as the point of inflection of $R(T)$. The variation of T^* with that of both bulk (thickness) and surface (gating) contributions to conductance confirms the two-channel model and disproves the prevailing idea that there is a static transition temperature; rather, it is merely defined by the relative contributions from the two conduction channels.

In conclusion, we have demonstrated methods to tune both the bulk and surface contributions to the electrical conduction, thereby providing conclusive evidence for surface-dominated transport in SmB_6 at low temperature. The evolution of transport with both sample thickness and surface gate tuning fits well to a two-channel conduction model involving a bulk, activated channel and a surface metallic channel. Furthermore, the measurement of gated surface conducting states using a Corbino lead geometry allows for the direct determination of the electron-like sign of the surface charge carriers as well as the charge carrier density ($\sim 2 \times 10^{14} \text{ cm}^{-2}$) and mobility ($\sim 133 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$). The charge carrier sign and density are in good agreement with previous photoemission results for SmB_6 . This study adds valuable information to our understanding of the proposed topological surface conduction in SmB_6 and provides promising use of gate-tuned devices structured on bulk crystal samples for both fundamental and applied studies of these unique states.

The authors would like to acknowledge R. L. Greene, X. Wang and Y. Nakajima for valuable discussion and experimental assistance. Research at the University of Maryland was supported by AFOSR-MURI (FA9550-09-1-0603) and NSF (DMR-0952716 and DMR-1105224). M.S.F. is supported by an ARC Laureate Fellowship.

* These authors contributed equally to this work

† paglione@umd.edu

- [1] M. Dzero, K. Sun, V. Galitski, and P. Coleman, Phys. Rev. Lett. **104**, 106408 (2010).
- [2] M. Dzero, K. Sun, P. Coleman, and V. Galitski, Phys. Rev. B **85**, 045130 (2012).

- [3] V. Alexandrov, M. Dzero, and P. Coleman, *Physical Review Letters* **111**, 226403 (2013).
- [4] C. L. Kane and E. J. Mele, *Phys. Rev. Lett.* **95**, 146802 (2005).
- [5] L. Fu, C. L. Kane, and E. J. Mele, *Phys. Rev. Lett.* **98**, 045302 (2007).
- [6] N. P. Butch, P. Syers, A. B. Sushkov, G. S. Jenkins, H. D. Drew, and J. Paglione, *Phys. Rev. B* **81**, 241301R (2010).
- [7] D.-X. Qu, Y. S. Hor, J. Xiong, R. J. Cava, and N. P. Ong, *Science* **329**, 821 (2010).
- [8] J. W. Allen, B. Batlogg, and P. Wachter, *Phys. Rev. B* **20**, 4807 (1979).
- [9] J. W. Allen, R. M. Martin, B. Batlogg, and P. Wachter, *J. Appl. Phys.* **49**, 2078 (1978).
- [10] B. Gorshunov, N. Sluchanko, A. Volkov, M. Dressel, G. Knebel, A. Loidl, and S. Kunii, *Phys. Rev. B* **59**, 1808 (1999).
- [11] S. Gabani, K. Flachbart, V. Pavlik, M. Orendac, E. Konovalova, Y. Paderno, and J. Sebek, *Czech. J. Phys.* **52**, 279 (2002).
- [12] N. E. Sluchanko, V. V. Glushkov, B. P. Gorshunov, S. V. Demishev, M. V. Kondrin, A. A. Pronin, A. A. Volkov, A. K. Savchenko, G. Grner, Y. Bruynseraede, V. V. Moshchalkov, and S. Kunii, *Phys. Rev. B* **61**, 9906 (2000).
- [13] K. Flachbart, K. Gloos, E. Konovalova, Y. Paderno, M. Reiffers, P. Samuely, and P. Švec, *Phys. Rev. B* **64**, 085104 (2001).
- [14] X. Zhang, N. P. Butch, P. Syers, S. Ziemak, R. L. Greene, and J. Paglione, *Phys. Rev. X* **3**, 011011 (2013).
- [15] M. M. Yee, Y. He, A. Soumyanarayanan, D.-J. Kim, Z. Fisk, and J. E. Hoffman, *ArXiv e-prints* (2013), arXiv:1308.1085 [cond-mat.str-el].
- [16] S. Rössler, T.-H. Jang, D.-J. Kim, L. H. Tjeng, Z. Fisk, F. Steglich, and S. Wirth, *Proc. Natl. Acad. Sci.* **111**, 4798 (2014).
- [17] W. Ruan, C. Ye, M. Guo, F. Chen, X. Chen, G.-M. Zhang, and Y. Wang, *Phys. Rev. Lett.* **112**, 136401 (2014).
- [18] N. Xu, X. Shi, P. K. Biswas, C. E. Matt, R. S. Dhaka, Y. Huang, N. C. Plumb, M. Radović, J. H. Dil, E. Pomjakushina, K. Conder, A. Amato, Z. Salman, D. M. Paul, J. Mesot, H. Ding, and M. Shi, *Phys. Rev. B* **88**, 121102 (2013).
- [19] M. Neupane, N. Alidoust, S. Y. Xu, T. Kondo, Y. Ishida, D. J. Kim, C. Liu, I. Belopolski, Y. J. Jo, T. R. Chang, H. T. Jeng, T. Durakiewicz, L. Balicas, H. Lin, A. Bansil, S. Shin, Z. Fisk, and M. Z. Hasan, *Nature Comm.* **4**, 2991 (2013).
- [20] E. Frantzeskakis, N. de Jong, B. Zwartsenberg, Y. K. Huang, Y. Pan, X. Zhang, J. X. Zhang, F. X. Zhang, L. H. Bao, O. Tegus, A. Varykhalov, A. de Visser, and M. S. Golden, *Phys. Rev. X* **3**, 041024 (2013).
- [21] A. Menth, E. Buehler, and T. H. Geballe, *Phys. Rev. Lett.* **22**, 295 (1969).
- [22] J. C. Nickerson, R. M. White, K. N. Lee, R. Bachmann, T. H. Geballe, and G. W. Hull, *Phys. Rev. B* **3**, 2030 (1971).
- [23] T. Caldwell, A. P. Reyes, W. G. Moulton, P. L. Kuhns, M. J. R. Hoch, P. Schlottmann, and Z. Fisk, *Phys. Rev. B* **75**, 075106 (2007).
- [24] J. Beille, M. B. Maple, J. Wittig, Z. Fisk, and L. E. DeLong, *Phys. Rev. B* **28**, 7397 (1983).
- [25] S. Wolgast, C. Kurdak, K. Sun, J. W. Allen, D. J. Kim, and Z. Fisk, *Phys. Rev. B* **88**, 180405 (2013).
- [26] D. J. Kim, S. Thomas, T. Grant, J. Botimer, Z. Fisk, and J. Xia, *Scientific Reports* **3**, 150 (2013).
- [27] W. A. Phelan, S. M. Koochpayeh, P. Cottingham, J. W. Freeland, J. C. Leiner, C. L. Broholm, and T. M. McQueen, *Phys. Rev. X* **4**, 031012 (2014).
- [28] G. Li, Z. Xiang, F. Yu, T. Asaba, B. Lawson, P. Cai, C. Tinsman, A. Berkley, S. Wolgast, Y. S. Eo, D.-J. Kim, C. Kurdak, J. W. Allen, K. Sun, X. H. Chen, Y. Y. Wang, Z. Fisk, and L. Li, *Science* **346**, 1208 (2014).
- [29] J. Jiang, A. Li, T. Zhang, Z. Sun, F. Chen, Z. R. Ye, M. Xu, Q. Q. Ge, S. Y. Tan, X. H. Niu, M. Xia, B. P. Xie, Y. F. Li, X. H. Chen, H. H. Wen, and D. L. Feng, *Nature Comm.* **4**, 3010 (2013).
- [30] Z. H. Zhu, A. Nicolau, G. Levy, N. P. Butch, P. Syers, X. F. Wang, J. Paglione, G. A. Sawatzky, I. S. Elfimov, and A. Damascelli, *Phys. Rev. Lett.* **111**, 216402 (2013).
- [31] J. D. Denlinger, J. W. Allen, J. S. Kang, and K. Sun, arXiv:1312.6637: [cond-mat.str-el].
- [32] S. Suga, K. Sakamoto, T. Okuda, K. Miyamoto, K. Kuroda, A. Sekiyama, J. Yamaguchi, H. Fujiwara, A. Irizawa, T. Ito, S. Kimura, T. Balashov, W. Wulfskel, S. Yeo, F. Iga, and S. Imada, *J. Phys. Soc. Japan* **83**, 038001 (2014).
- [33] N. Xu, P. K. Biswas, J. H. Dil, R. S. Dhaka, G. Landolt, S. Muff, C. E. Matt, X. Shi, N. C. Plumb, M. Radović, E. Pomjakushina, K. Conder, A. Amato, S. V. Borisenko, R. Yu, H.-M. Weng, Z. Fang, X. Dai, J. Mesot, H. Ding, and M. Shi, *Nature Comm.* **5**, (2014).
- [34] Y. Nakajima, P. S. Syers, X. Wang, R. Wang, and J. Paglione, *ArXiv e-prints* (2013), arXiv:1312.6132 [cond-mat.str-el].
- [35] A. Kebede, M. C. Aronson, C. M. Buford, P. C. Canfield, J. H. Cho, B. R. Coles, J. C. Cooley, J. Y. Coulter, Z. Fisk, G. D. Goette, W. L. Hults, A. Lacerda, T. D. McLendon, P. Tiwari, and J. L. Smith, *Physica B* **223**, 256 (1996).
- [36] We neglect a small dependence of surface conductance on changing surface area due to reduction in cross section perimeter.
- [37] S. Cho, N. P. Butch, J. Paglione, and M. S. Fuhrer, *Nano Lett.* **11**, 1925 (2011).
- [38] D. Kim, S. Cho, N. P. Butch, P. Syers, K. Kirshenbaum, S. Adam, J. Paglione, and M. S. Fuhrer, *Nature Phys.* **8**, 460 (2012).
- [39] D. Kim, P. Syers, N. P. Butch, J. Paglione, and M. S. Fuhrer, *Nature Comm.* **4**, (2013).
- [40] The experiment is designed to show the difference between a quasi-two-dimensional (surface) state vs a fully bulk state, but cannot place an estimate on the exact thickness of the surface state.
- [41] D. J. Kim, J. Xia, and Z. Fisk, *Nature Materials* **13**, 466 (2014).
- [42] S. Yeo, K. Song, N. Hur, Z. Fisk, and P. Schlottmann, *Phys. Rev. B* **85**, 115125 (2012).
- [43] N. E. Sluchanko, A. A. Volkov, V. V. Glushkov, B. P. Gorshunov, S. V. Demishev, M. V. Kondrin, A. A. Pronin, N. A. Samarin, Y. Bruynseraede, V. V. Moshchalkov, and S. Kunii, *Journal of Experimental and Theoretical Physics* **88**, 533 (1999).
- [44] J. C. Cooley, M. C. Aronson, Z. Fisk, and P. C. Canfield, *Phys. Rev. Lett.* **74**, 1629 (1995).
- [45] Recent angle-resolved photoemission experiments gener-

ally observe a small “ α ” band with Fermi surface area $0.020(2\pi/a)^2$ and two large “ β ” bands each with area $0.134(2\pi/a)^2$ [18, 19, 29]. With a unit cell of $(4.13 \text{ \AA})^2$ and assuming each β band is electron-like and singly

degenerate, this yields a total electron density of $\sim 1.7 \times 10^{14} \text{ cm}^{-2}$, which is very close to our estimate.