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Measurement of the Neutron Radius of ^{208}Pb Through Parity-Violation in Electron Scattering

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We report the first measurement of the parity-violating asymmetry A_{PV} in the elastic scattering of polarized electrons from ^{208}Pb . A_{PV} is sensitive to the radius of the neutron distribution (R_n). The result $A_{PV} = 0.656 \pm 0.060$ (stat) ± 0.014 (syst) ppm corresponds to a difference between the radii of the neutron and proton distributions $R_n - R_p = 0.33^{+0.16}_{-0.18}$ fm and provides the first electroweak observation of the neutron skin which is expected in a heavy, neutron-rich nucleus.

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Nuclear charge densities have been accurately measured with electron scattering and have become our picture of the atomic nucleus, see for example [1]. In contrast, our knowledge of neutron densities comes primarily from hadron scattering experiments involving, for example, pions [2], protons [3–5], or antiprotons [6, 7], the interpretation of which requires a model-dependent description of the non-perturbative strong interaction. Due to the fact that the weak charge of the neutron is much larger than that of the proton, the measurement of parity violation in electron scattering provides a model-independent probe of neutron densities that is free from most strong-interaction uncertainties [8].

In the Born approximation, the parity violating cross-section asymmetry for longitudinally polarized electrons elastically scattered from an unpolarized nucleus, A_{PV} , is proportional to the weak form factor $F_W(Q^2)$. This is the Fourier transform of the weak charge density, which is closely related to the neutron density, and therefore the neutron density can be extracted from an electro-weak measurement [8].

$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \approx \frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \frac{F_W(Q^2)}{F_{ch}(Q^2)} \quad (1)$$

where $\sigma_{R(L)}$ is the differential cross section for elastic scattering of right (R) and left (L) handed longitudinally polarized electrons, G_F is the Fermi constant, α the fine structure constant, and $F_{ch}(Q^2)$ is the Fourier transform of the known charge density. However, the Born approximation is not valid for a heavy nucleus and Coulomb-distortion effects must be included. These have been accurately calculated [9] because the charge density is well known, and many other details relevant for a practical parity-violation experiment to measure neutron densities have been discussed in a previous publication [10].

One system of particular interest is the doubly-magic nucleus ^{208}Pb , which has 44 more neutrons than protons; some of these extra neutrons are expected to be found in the surface, where they form a neutron-rich skin. The thickness of this skin is sensitive to nuclear dynamics and provides fundamental nuclear structure information. A number of mean-field-theory models have been developed that agree with the world's body of data on nuclear

charge distributions and other nuclear properties [11–15]. For ^{208}Pb , these are consistent with a radius of the point-neutron distribution R_n between 0.0 – 0.4 fm larger than that of the point-proton distribution R_p . In this paper we report a first measurement of A_{PV} from ^{208}Pb , which is sensitive to the existence of the neutron skin.

The value of the neutron radius of ^{208}Pb has important implications for models of nuclear structure and their application in atomic physics and astrophysics. There is a strong correlation between R_n of ^{208}Pb and the pressure of neutron matter P at densities near 0.1 fm^{-3} (about 2/3 of nuclear density) [16]. A larger P will push neutrons out against surface tension and increase R_n . Therefore measuring R_n constrains the equation of state (EOS), the pressure as a function of density, of neutron matter.

The correlation between R_n and the radius of a neutron star r_{NS} is also very interesting [17]. In general, a larger R_n implies a stiffer EOS, with a larger pressure, that will also suggest r_{NS} is larger. Recently there has been great progress in deducing r_{NS} from X-ray observations. From observations of X-ray bursts, Özel *et al.* [18] find r_{NS} is very small, near 10 km, implying that the EOS softens at high density which is suggestive of a transition to an exotic phase of QCD. In contrast, Steiner *et al.* [19] conclude that r_{NS} is near 12 km, leading to a prediction that $R_n - R_p = 0.15 \pm 0.02$ fm for ^{208}Pb . This implies a stiffer EOS which leaves little room for softening due to a phase transition at high density.

Recently Hebeler *et al.* [20] used chiral perturbation theory to calculate the EOS of neutron matter including important contributions from three-neutron forces. From their EOS, they predict $R_n - R_p = 0.17 \pm 0.03$ fm for ^{208}Pb . Monte Carlo calculations by Carlson *et al.* [21] also find sensitivity to three-neutron forces. The measurement of R_n provides an important check of fundamental neutron matter calculations, and constrains three-neutron forces.

The EOS of neutron-rich matter is closely related to the symmetry energy S . There is a strong correlation between R_n and the density dependence of the symmetry energy $dS/d\rho$, with ρ as the baryon density. The symmetry energy can be probed in heavy-ion collisions [22]. For example, $dS/d\rho$ has been extracted from isospin diffusion

data [23] using a transport model.

The symmetry energy S helps determine the composition of a neutron star. A large S at high density would imply a large proton fraction, which would allow the direct Urca process [24] of rapid neutrino cooling. If $R_n - R_p$ in ^{208}Pb were large, it is likely that massive neutron stars would cool quickly by direct Urca. In addition, the transition density from a solid neutron star crust to the liquid interior is strongly correlated with $R_n - R_p$ [25].

Reinhard and Nazarewicz claim that $R_n - R_p$ is tightly correlated with the dipole polarizability α_D [26] and Tamii et al. use this correlation to infer $R_n - R_p$ from a new measurement of α_D [27].

Atomic parity violation (APV) is also sensitive to R_n [10, 28, 29]. A future low-energy test of the standard model may involve the combination of a precise APV experiment along with PV electron scattering to constrain R_n [28]. Alternatively, measuring APV for a range of isotopes could provide information on neutron densities [30].

The measurement was carried out in Hall A at the Thomas Jefferson National Accelerator Facility. The experimental configuration is similar to that used previously for studies of the weak form factor of the proton and ^4He [31–33]. A 50 to 70 μA continuous-wave beam of longitudinally polarized 1.06 GeV electrons was incident on a 0.55 mm thick isotopically pure ^{208}Pb target foil. A 4 mm \times 4 mm square beam raster prevented the target from melting. Two 150 μm diamond foils sandwiched the lead foil to improve thermal conductance to a copper frame cooled to 20K with cryogenic helium. Elastically scattered electrons were focused onto thin quartz detectors in the twin High Resolution Spectrometers (HRS) [34]. The addition of a pair of dipole septum magnets between the target and the HRSs allowed us to achieve a forward scattering angle of $\theta_{lab} \sim 5^\circ$. The HRS momentum resolution ensured that only elastic events (and a negligible fraction of inelastic events from the 2.6 MeV first excited state) were accepted by the quartz detectors. Cherenkov light from each quartz bar traversed air light guides and were detected by 2-inch quartz-window photo-multipliers (PMT).

The polarized electron beam originated from a strained GaAsP photocathode illuminated by circularly polarized light [35]. The accelerated beam was directed into Hall A, where its intensity, energy, polarization, and trajectory on target were inferred from the response of several monitoring devices. The sign of the laser circular polarization determined the electron helicity; this was held constant for periods of 8.33 ms, referred to as “windows”. The integrated responses of detector PMTs and beam monitors were digitized by an 18-bit ADC and recorded for each window. Two “window quadruplet” patterns of helicity states (+ − + or − + −) ensured complementary measurements at the same phase relative to the 60 Hz line

power, thus canceling power-line noise from the asymmetry measurement. The right-left helicity asymmetry in the integrated detector response, normalized to the beam intensity, was computed for sets of complementary helicity windows in each quadruplet to form the raw asymmetry A_{raw} . The sequence of these patterns was chosen with a pseudo-random number generator.

Loose requirements were imposed on beam quality, removing periods of position, energy, or beam-intensity instability. No helicity-dependent cuts were applied, leaving a final data sample of 2×10^7 helicity-window quadruplets. The design of the apparatus ensured that, after all corrections, the fluctuations in the fractional difference of the PMT response between a pair of successive windows was dominated by scattered-electron counting statistics for rates up to 1 GHz. This facilitated the ability to achieve an A_{PV} precision significantly better than 100 parts per billion (ppb) in a reasonable length of time. Careful attention to the design and configuration of the photocathode laser optics [36] ensured that spurious beam-induced asymmetries were under control at this level.

Random fluctuations in beam position and energy contributed the largest source of noise beyond counting statistics in A_{raw} . Typical beam jitter in window-quadruplets was less than 2 parts per million (ppm) in energy, and 20 μm in position. This noise contribution was reduced by measuring window differences Δx_i using beam position monitors and applying a correction $A_{beam} = \sum c_i \Delta x_i$. The c_i ’s were measured several times each hour from calibration data in which the beam was modulated by using steering coils and an accelerating cavity. The largest of the c_i ’s was ~ 50 ppm/ μm . The noise in the resulting $A_{corr} = A_{raw} - A_{beam}$ was 210 (180) ppm per quadruplet, for a beam current of 50 (70) μA , dominated by counting statistics (~ 1 GHz at 70 μA). Non-uniformities in target thickness due to thermal damage caused window-to-window luminosity fluctuations from variations in the target area sampled by the rastered beam, leading to the degradation of A_{corr} by $\sim 40\%$. This source of noise was eliminated by locking the raster pattern frequency to a multiple of the helicity frequency. Low-current calibration data, triggered on individual scattered electrons, were regularly collected to evaluate the thickness of lead relative to diamond.

Sensitivity of A_{corr} to a transverse component of the beam polarization, coupled to the vector analyzing powers (A_T) for ^{208}Pb and ^{12}C , was studied using special runs with fully transverse beam polarization. The symmetry of the detector configuration as well as the measured A_T values (to be published separately) resulted in an upper bound for a possible correction to A_{corr} of 0.2%. The A_{raw} and A_{corr} window-pair distributions for the two complete data samples had negligible non-Gaussian tails over more than four orders of magnitude. To test the accuracy of error calculations and general statistical

behavior of the data, A_{corr} averages and statistical errors were studied for typical one-hour runs, consisting of $\sim 50k$ quadruplets each. This set of 316 average A_{corr} values, normalized by the corresponding statistical error, populated a Gaussian distribution of unit variance, as expected.

A half-wave ($\lambda/2$) plate was periodically inserted into the injector laser optical path, reversing the sign of the electron beam polarization relative to both the electronic helicity control signals and the voltage applied to the polarized source laser electro-optics. Roughly equal statistics were collected with this waveplate inserted and retracted, suppressing many possible sources of systematic error. An independent method of helicity reversal was feasible with a pair of Wien spin-rotators separated by a solenoid, providing an additional powerful check of systematic control. Reversing the direction of the solenoidal field reversed the electron beam helicity while the beam optics, which depend on the square of the solenoidal magnetic field, were unchanged. The $\lambda/2$ reversal was done about every 12 hours and the magnetic spin reversal was performed every few days. The dataset consisting of a period between two successive $\lambda/2$ or magnetic spin-reversals is referred to as a “slug”.

The spin reversals resulted in cumulative differences in beam position and energy of only 4 nm and 0.6 ppb respectively, leading to a run-averaged $A_{beam} = -39.0 \pm 5.9$ ppb. The asymmetry in beam charge, corrected by the intensity normalization of A_{raw} , was 84.0 ± 1.3 ppb, with the error determined using the correlation of measured beam intensity to PMT response which demonstrated the beam intensity monitors were linear to better than 1.5%. Nonlinearity in the PMT response was limited to 1% in bench-tests that mimicked running conditions. As shown in Table I, the values of A_{corr} are consistent within statistical errors for each of the reversal states. The reduced χ^2 for A_{corr} “slug” averages is close to one in every case, indicating that any residual beam-related systematic effects were small and randomized over the time period of $\lambda/2$ reversals. The final result is $A_{corr} = 594 \pm 50(\text{stat}) \pm 9(\text{syst})$ ppb where the systematic uncertainty includes possible effects from A_{beam} , non-linearity in the detectors or beam charge monitors, and transverse asymmetry. The physics asymmetry A_{PV} is formed from A_{corr} by correcting for the beam polarization P_b and background fractions f_i with asymmetries A_i

$$A_{PV} = \frac{1}{P_b} \frac{A_{corr} - P_b \sum_i A_i f_i}{1 - \sum_i f_i}. \quad (2)$$

These corrections are summarized in Table II.

The fraction of the accepted flux from ^{12}C in the detectors varied with time due to changes in the target; averaged over the run, the fraction $f = (6.3 \pm 0.6)\%$. The asymmetry of this background was determined to

| $\lambda/2$ plate | Spin-rotator | A_{corr} (ppb) | δA_{corr} (ppb) | $\chi^2/\text{d.o.f.}$ |
|-------------------|--------------|------------------|-------------------------|------------------------|
| OUT | RIGHT | 606 | 113 | 1.03 |
| IN | RIGHT | 492 | 107 | 0.74 |
| OUT | LEFT | 565 | 95 | 1.12 |
| IN | LEFT | 687 | 92 | 1.03 |
| Average | | 594 | 50 | 0.99 |

TABLE I: Values of A_{corr} and the statistical error, for each helicity reversal state and for the grand average. The χ^2 per degree of freedom for each average is also shown.

| Correction | Absolute (ppb) | Relative(%) |
|-----------------------------|-----------------|-----------------|
| Beam Charge Normalization | -84.0 ± 1.5 | -12.8 ± 0.2 |
| Beam Asymmetries A_{beam} | 39.0 ± 7.2 | 5.9 ± 1.1 |
| Target Backing | -8.8 ± 2.6 | -1.3 ± 0.4 |
| Detector Nonlinearity | 0 ± 7.6 | 0 ± 1.2 |
| Transverse Asymmetry | 0 ± 1.2 | 0 ± 0.2 |
| Polarization P_b | 70.9 ± 8.3 | 10.8 ± 1.3 |
| Total | 17.1 ± 13.7 | $2.6 \pm 2.1\%$ |

TABLE II: Corrections to A_{PV} and systematic errors.

be $A_{PV}^C = 817 \pm 41$ ppb using the Standard Model value for the e -N weak neutral isoscalar coupling and the measured kinematics, with the uncertainty bounded by the precision measurement of A_{PV} from ^4He [31]. This was the only non-negligible background. An additional possible systematic error in $\langle Q^2 \rangle$ lay in the determination of the absolute value of θ_{lab} . A nuclear recoil technique using a water cell target [32] limited the scale error on $\langle Q^2 \rangle$ to 1%.

The spectrometer acceptance function $\epsilon(\theta)$ characterizes the probability, as a function of scattering angle θ ,

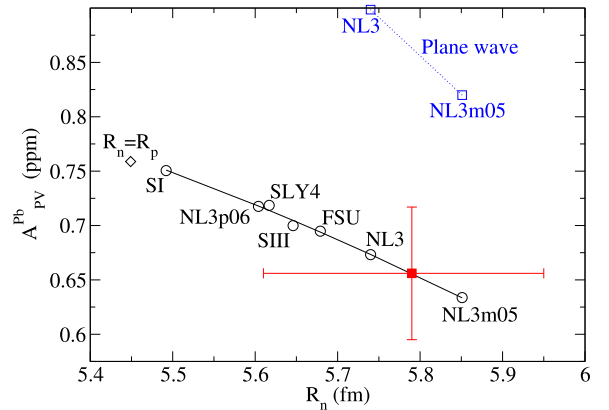


FIG. 1: Result of this experiment (red square) vs neutron point radius R_n in ^{208}Pb . Distorted-wave calculations for seven mean-field neutron densities are circles [39], while the diamond marks the expectation for $R_n = R_p$. The blue squares show plane wave impulse approximation results.

for an electron to reach the detector after elastically scattering from ^{208}Pb . For example, the asymmetry averaged over the acceptance would be

$$\langle A \rangle = \frac{\int d\theta \sin \theta A(\theta) \frac{d\sigma}{d\Omega} \epsilon(\theta)}{\int d\theta \sin \theta \frac{d\sigma}{d\Omega} \epsilon(\theta)} \quad (3)$$

where $\frac{d\sigma}{d\Omega}$ is the cross section. See Supplemental Material at [URL will be inserted by publisher] for the acceptance function $\epsilon(\theta)$. The observed distribution of events corrected for the cross section, the background from the carbon (diamond) backing, and the effects of multiple scattering is used to extract $\epsilon(\theta)$; corrections for energy loss in the target were negligible. To compare to predictions, one must integrate the theoretical asymmetry and the Q^2 over $\epsilon(\theta)$. The systematic error in $\epsilon(\theta)$ was evaluated from reasonable variations in the parameters of the simulation and resulted in an additional equivalent error in $\langle Q^2 \rangle$ of 0.8%. Added in quadrature to the error arising from knowledge of $\langle \theta \rangle$, we obtain an overall error in $\langle Q^2 \rangle$ of 1.3%. We do not include this uncertainty in the total systematic uncertainty of the asymmetry. Using a calculation by Horowitz [9], dA_{PV}/dQ^2 is approximately 30 ppm/GeV², which would correspond to an additional systematic uncertainty on A_{PV} of 3 ppb (0.5% of A_{PV}).

The beam polarization was continuously monitored by a Compton polarimeter. Helicity-dependent asymmetries in the integrated signal from backscattered Compton photons yielded $P_b = (88.2 \pm 0.1 \pm 1.0)\%$ averaged over the duration of the run. The beam polarization was stable within systematic errors. An independent Møller polarimeter making nine measurements at different times during the run gave $P_b = (90.3 \pm 0.1 \pm 1.1)\%$. We used an average of these two measurements, $P_b = (89.2 \pm 1.0)\%$ which conservatively accounts for the correlated systematic errors between the two measurements.

After all corrections,

$$A_{PV}^{Pb} = 656 \pm 60 \text{ (stat)} \pm 14 \text{ (syst) ppb}$$

at $\langle Q^2 \rangle = 0.00880 \pm 0.00011 \text{ GeV}^2$. This result is displayed in Figure 1, in which models predicting the point-neutron radius illustrate the correlation of A_{PV}^{Pb} and R_n [39].

Seven non-relativistic and relativistic mean field models were chosen that have charge densities and binding energies in good agreement with experiment, and that span a large range in R_n . The weak charge density ρ_w was calculated from model point proton ρ_p and neutron ρ_n densities, $\rho_w(r) = q_p \rho_{ch}(r) + q_n \int d^3r' [G_E^p \rho_n + G_E^n \rho_p]$, using proton $q_p = 0.0721$ and neutron $q_n = -0.9878$ weak charges that include radiative corrections. Here G_E^p (G_E^n) is the Fourier transform of the proton (neutron) electric form factor. The Dirac equation was solved [9] for an electron scattering from ρ_w and the experimental ρ_{ch} [1], and the resulting $A_{PV}(\theta)$ integrated over the acceptance, Eq. 3, to yield the open circles in Fig. 1. The

importance of Coulomb distortions is emphasized by indicating results from plane-wave calculations, which are not all contained within the vertical axis range of the figure. A least squares fit of the model results yields $R_n \approx 6.156 + 1.675 \langle A \rangle - 3.420 \langle A \rangle^2 \text{ fm}$ (with $\langle A \rangle$ in ppm), as illustrated. Comparing this to the measured A_{PV}^{Pb} implies a value for $R_n = 5.78^{+0.16}_{-0.18} \text{ fm}$. More details of this analysis, along with additional information such as the weak charge form factor and weak radius, will be presented in a future publication [37].

Assuming a point-proton radius of 5.45 fm [38], corresponding to the measured charge radius of 5.50 fm [1], implies that the neutron distribution is 1.8σ larger than that of the protons: $R_n - R_p = 0.33^{+0.16}_{-0.18} \text{ fm}$ [39] (see also [40]). A future run is planned which will reduce the quoted uncertainty by a factor of three [41], to discriminate between models and allow predictions relevant for the description of neutron stars and parity violation in atomic systems.

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