Search for B_{s}^{0}→μ^{+}μ^{-} and B^{0}→μ^{+}μ^{-} Decays with CDF II

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A search has been performed for $B_s^0 \to \mu^+\mu^-$ and $B^0 \to \mu^+\mu^-$ decays using $7\,\mathrm{fb}^{-1}$ of integrated luminosity collected by the CDF II detector at the Fermilab Tevatron collider. The observed number of B^0 candidates is consistent with background-only expectations and yields an upper limit on the branching fraction of $\mathcal{B}(B^0 \to \mu^+\mu^-) < 6.0 \times 10^{-9}$ at 95% confidence level. We observe an excess of B_s^0 candidates. The probability that the background processes alone could produce such an excess or larger is 0.27%. The probability that the combination of background and the expected standard model rate of $B_s^0 \to \mu^+\mu^-$ could produce such an excess or larger is 1.9%. These data are used to determine $\mathcal{B}(B_s^0 \to \mu^+\mu^-) = (1.8^{+1.1}_{-0.9}) \times 10^{-8}$ and provide an upper limit of $\mathcal{B}(B_s^0 \to \mu^+\mu^-) < 4.0 \times 10^{-8}$ at 95% confidence level.

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Studies of flavor-changing neutral current (FCNC) decays have played an important role in formulating the theoretical description of particle physics known as the standard model (SM). In the SM all neutral currents conserve flavor so that FCNC decays do not occur at lowest order. The decays of B_s^0 mesons (with a quark content of $\overline{b}s$) and B^0 mesons ($\overline{b}d$) into a dimuon pair ($\mu^+\mu^-$) [1] are examples of FCNC processes that can occur in the SM through higher order loop diagrams. Their branching fractions are predicted in the SM to be $(3.2 \pm 0.2) \times 10^{-9}$ and $(1.0 \pm 0.1) \times 10^{-10}$, respectively [2]. A wide variety of beyond-SM theories predict significant increases over the SM branching fraction [3], making the study of

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$B_s^0 \to \mu^+\mu^-$ and $B^0 \to \mu^+\mu^-$ decays one of the most sensitive indirect searches for new physics. Published upper limits [4–6] contribute significantly to our knowledge of the available new physics parameter space [7–11].

We report a search for $B_s^0 \to \mu^+\mu^-$ and $B^0 \to \mu^+\mu^-$ decays using $p\bar{p}$ data corresponding to an integrated luminosity of 7 fb^{-1} collected with the Collider Detector at Fermilab (CDF II). The sensitivity of this analysis is significantly improved with respect to the previous analysis [4] due to the higher integrated luminosity of the event sample, a 20% increase in the signal acceptance, and the use of an improved neural-network (NN) discriminant that provides approximately twice the background rejection for the same signal efficiency.

A detailed description of the CDF II detector can be found in Ref. [12]. A charged particle tracking system provides precise vertex determination and momentum measurements in a pseudorapidity range $|\eta| < 1.0$. Additionally, the system measures the ionization per unit path length dE/dx for particle identification. Beyond the tracking detectors are electromagnetic and hadronic calorimeters, which are surrounded by drift chambers used to detect muons in the central region (C) $|\eta| < 0.6$ and the forward region (F) $0.6 < |\eta| < 1.0$.

The online (trigger) requirements used to collect the data sample and the initial set of baseline requirements used in the analysis are the same as those described in Ref. [13]. The events are collected using a set of dimuon triggers [12] and must satisfy either of two sets of requirements corresponding to different topologies: CC events have both muon candidates detected in the central region, while CF events have one central muon and another muon detected in the forward region. Since the expected signal-to-background ratios are different, the two topologies are treated separately. The acceptance of the analysis is improved by 20% by using additional foward muon candidates and by using muon candidates that traverse detector regions previously excluded due to their rapidly changing trigger efficiency. The larger data sample has allowed us to obtain a detailed understanding of the trigger performance in these regions so that we can confidently include these muon candidates in the current analysis. The baseline selection requires high quality muon candidates with transverse momentum relative to the beam direction of $p_T > 2.0\,(2.2)\,\mathrm{GeV}/c$ in the central (forward) region. The muon pairs are required to have an invariant mass in the range $4.669 < m_{\mu\mu} < 5.969\,\mathrm{GeV}/c^2$ and are constrained to originate from a common well measured three-dimensional (3D) vertex. A likelihood method [14] together with a dE/dx based selection [15] are used to further suppress contributions from hadrons misidentified as muons. The baseline requirements also demand that the measured proper decay length of the B candidate λ with its uncertainty σ_λ satisfy $\lambda/\sigma_\lambda > 2$; the 3D opening angle between the momentum of the dimuon pair and the displacement vector between the primary $p\bar{p}$ collision vertex and the dimuon vertex $\Delta\Omega < 0.7$ rad; and the B-candidate track isolation [16] $I > 0.50$. There are 48 279 CC and 52 179 CF muon pairs that fulfill the trigger and baseline selection requirements.

A sample of $B^+ \to J/\psi\ K^+$ events serves as a normalization mode. The $B^+ \to J/\psi\ K^+$ sample is collected using the same dimuon triggers and selection requirements so that common systematic uncertainties are suppressed. An additional requirement on the kaon candidate $p_T > 1$ GeV/c is made to limit the p_T range to a region where the tracking efficiency is well understood.

For the final selection, we define search regions around the known B_s^0 and B^0 masses [17]. These regions correspond to approximately $\pm 2.5\sigma_m$, where $\sigma_m \approx 24\,\mathrm{MeV}/c^2$ is the estimated two-track mass resolution. The sideband regions $5.0 < m_{\mu\mu} < 5.169\,\mathrm{GeV}/c^2$ and $5.469 < m_{\mu\mu} < 5.969\,\mathrm{GeV}/c^2$ are used to estimate combinatorial backgrounds. Backgrounds from $B \to h^+h'^-$ decays (where $h, h' = \pi^\pm$ or $K^\pm$), which peak in the signal mass region, are estimated separately.

Fourteen variables are used to construct a NN discriminant ν_N that ranges from 0 to 1 and enhances the signal-to-background ratio [18]. The variables include dimuon vertex related information (e.g. λ/σ_λ), the impact parameters with respect to the primary vertex and transverse momenta of the muons, the isolation of the B candidate, and the opening angle $\Delta\Omega$. The NN is trained with background events sampled from the sideband regions and signal events generated with a simulation described below. Only a fraction of the total number of background and simulated signal events are used to train the NN. The remainder are used to test for NN overtraining and to determine the signal and background efficiencies. Several tests are done to ensure ν_N is independent of $m_{\mu\mu}$.

All selection criteria were finalized before revealing the content of the signal regions. The optimization of the criteria used the expected upper limit on the $B_s^0 \to \mu^+\mu^-$ branching fraction as a figure of merit. To exploit the difference in the $m_{\mu\mu}$ distributions between signal and background and the improved suppression of combinatorial background at large ν_N, the data is divided into sub-samples in the $(\nu_N, m_{\mu\mu})$ plane. The CC and CF samples are each divided into 40 sub-samples. There are eight bins in ν_N with bin boundaries $0.70, 0.76, 0.85, 0.90, 0.94, 0.97, 0.987, 0.995$ and 1. Within each ν_N bin we employ five $m_{\mu\mu}$ bins, each 24 MeV/c^2 wide, centered on the world average $B_s^0(B^0)$ mass. The expected backgrounds and efficiencies are calculated in each bin separately.

For measuring efficiencies, estimating backgrounds, and optimizing the analysis, samples of $B_s^0(B^0) \to \mu^+\mu^-$, $B^+ \to J/\psi\ K^+$, and $B \to h^+h'^-$ are generated with the PYTHIA program [19] and a CDF II detector simulation. The p_T spectrum and the I distribution of the B-mesons are weighted to match distributions measured in samples of $B^+ \to J/\psi\ K^+$ and $B_s^0 \to J/\psi\ \phi$ events [12].

We use a relative normalization to determine the $B_s^0 \to$

$\mu^+\mu^-$ branching fraction:

$$\mathcal{B}(B_s^0 \to \mu^+\mu^-) = \frac{N_s}{N_+}\,\frac{\alpha_+}{\alpha_s}\,\frac{\epsilon_+}{\epsilon_s}\,\frac{1}{\epsilon_N}\,\frac{f_+}{f_s}\,\mathcal{B}(B^+),$$

where N_s is the number of $B_s^0 \to \mu^+\mu^-$ candidate events. The observed number of $B^+ \to J/\psi\ K^+$ candidates is $N_+ = 22\,388 \pm 196$ and $9\,943 \pm 138$ in the CC and CF channels, respectively. The contribution of $B^+ \to J/\psi\ \pi^+$ events is negligible. We use $\mathcal{B}(B^+) = \mathcal{B}(B^+ \to J/\psi\ K^+ \to \mu^+\mu^- K^+) = (6.01 \pm 0.21) \times 10^{-5}$ [17] and the ratio of B-meson production fractions $f_+/f_s = 3.55 \pm 0.47$ [17]. The parameter α_s (α_+) is the acceptance of the trigger and ϵ_s (ϵ_+) is the efficiency of the reconstruction requirements for the signal (normalization) mode. The reconstruction efficiency includes trigger, track, muon, and baseline requirement efficiencies. The NN efficiency ϵ_N only applies to the signal mode since it is not used to select the $B^+ \to J/\psi\ K^+$ sample. The expression for $\mathcal{B}(B^0 \to \mu^+\mu^-)$ is derived by replacing B_s^0 with B^0 and f_+/f_s with $f_+/f_d = 1$. The ratios of acceptances α_+/α_s are 0.307 ± 0.018 and 0.197 ± 0.014 for the CC and CF topologies, respectively. These ratios are measured using simulated events. The uncertainties include contributions from systematic variations of the modeling of the B-meson p_T distributions and the longitudinal beam profile. The ratio of reconstruction efficiencies is $\epsilon_+/\epsilon_s = 0.81 \pm 0.03$ as determined from studies using samples of $J/\psi \to \mu^+\mu^-$ and $B^+ \to J/\psi\ K^+$ events collected with the same triggers. The uncertainty in ϵ_+/ϵ_s is dominated by kinematic differences between $J/\psi \to \mu^+\mu^-$ and $B_s^0(B^0) \to \mu^+\mu^-$ decays. The ϵ_N is estimated from the simulation. We assign a relative systematic uncertainty on ϵ_N of 4-7% , depending on ν_N bin, using comparisons of the NN performance in simulated and observed $B^+ \to J/\psi\ K^+$ event samples, and the statistical uncertainty on studies of the p_T and I distributions from observed $B_s^0 \to J/\psi\ \phi$ event samples. The $B^0 \to \mu^+\mu^-$ decay is determined to have the same acceptances and efficiencies. Treating CC and CF together, about 90% of simulated $B_s^0 \to \mu^+\mu^-$ events surviving the initial requirements have $\nu_N > 0.70$, with about 45% having $\nu_N > 0.995$. The expected SM yield of $B_s^0 \to \mu^+\mu^-$ events ranges from 0.05 in the lowest ν_N bin to 1.0 events in the highest ν_N bin summing the CC and CF contributions. The expected SM yield of $B^0 \to \mu^+\mu^-$ events is about thirty times smaller.

The expected background is obtained by summing contributions from the combinatorial background and from $B \to h^+h'^-$ decays. To estimate the combinatorial background, we fit the $m_{\mu\mu}$ distribution of sideband events with $\nu_N > 0.70$ to a linear function. We only use events with $m_{\mu\mu}$> 5 GeV/c^2 in order to suppress contributions from $b \to \mu^+\mu^- X$ decays. The slopes are then fixed, and the normalization is determined for each ν_N bin separately using the relevant sideband events. In addition to the statistical uncertainties of the slope and normalization parameters, systematic uncertainties are assigned by comparing results derived using alternative fit functions and ranges. The systematic uncertainties vary from about 7% for the lower ν_N bins to about 45% for the highest ν_N bins. The $B \to h^+h'^-$ contributions are estimated using efficiencies determined from the simulation, probabilities of misidentifying hadrons as muons measured in data, and normalizations derived from their branching fractions [15, 17]. The hadron misidentification probabilities are parameterized as a function of hadron p_T and instantaneous luminosity using a $D^0 \to K^-\pi^+$ data sample obtained from $D^{*+} \to D^0\pi^+$ decays. In addition to the statistical uncertainties from the D^0 sample, systematic uncertainties are assigned to account for residual variations of the misidentification probability due to variations in detector performance (primarily arising from occupancy and calibration effects) and for branching fraction uncertainties. For the B_s^0 modes there is an additional uncertainty from f_+/f_s. The estimated $B \to h^+h'^-$ background is approximately one quarter of the total background in the $B^0 \to \mu^+\mu^-$ search while in the $B_s^0 \to \mu^+\mu^-$ search it is a factor of ten smaller than both the combinatorial background and the SM signal. The expected background is shown in Fig. 1 for the $B_s^0 \to \mu^+\mu^-$ and $B^0 \to \mu^+\mu^-$ searches. The background estimates are cross-checked using three sets of independent control samples: $\mu^+\mu^-$ events with $\lambda < 0$ and $\mu^\pm\mu^\pm$ events, both of which are dominated by combinatorial backgrounds, and a misidentified-muon enhanced $\mu^+\mu^-$ sample with at least one muon candidate failing the muon quality requirements. The latter sample has a significant contribution from $B \to h^+h'^-$ backgrounds. We compare the predicted and observed number of events in each of these control samples for all 80 sub-samples and observe no significant discrepancies.

Two fits are performed on the data, a background-only fit (b) and a signal-plus-background fit ($s+b$) for which the branching fraction of the signal is left floating. A log-likelihood ratio is formed, $-2\ln Q$, where $Q = \mathcal{L}(s+b|\text{data})/\mathcal{L}(b|\text{data})$ and $\mathcal{L}(h|x)$ is the likelihood of hypothesis h given observation x; this likelihood is obtained by multiplying Poisson probabilities over all 80 sub-samples and is minimized with respect to the nuisance parameters that model our systematic uncertainties. To evaluate the consistency of the data in the signal region with our background model, we compare the observed value of $-2\ln Q$ with the distribution of $-2\ln Q$ obtained from an ensemble of background-only simulated experiments. The effects of systematic uncertainties are included in the simulated experiments by randomly choosing the nuisance parameters from Gaussian distributions. The fraction of simulated experiments with a value of $-2\ln Q$ less than that observed in the data is used to determine the p-value for the background-only hypothesis.

The data in the signal regions are shown in Fig. 1 using the $(\nu_N, m_{\mu\mu})$ binning from the optimization. In the

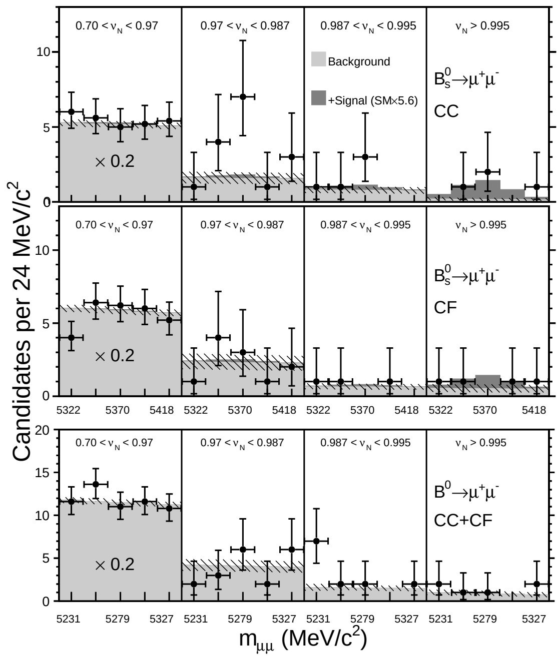

FIG. 1: For the B_s^0 and B^0 signal regions, the observed number of events (points) is compared to the total expected background (light grey) and its uncertainty (hatched) using the $(\nu_N, m_{\mu\mu})$ bins from the optimization. The background uncertainty is the quadrature sum of the relevant systematic uncertainties. The top and middle rows show the results in the B_s^0 mass signal region for the CC and CF channels, respectively. The bottom row shows the results in the B^0 mass signal region for the CC and CF channels combined. The results for the first 5 ν_N bins are combined (and scaled by 0.2) while the results for the last three bins are each shown separately. Also shown is the expected contribution from $B_s^0 \to \mu^+\mu^-$ events (dark gray) using a branching fraction that corresponds to the central value from the fit to the data, which is 5.6 times the expected SM value.

B^0 search region the data are consistent with the background prediction and have a p-value of 23%. In the B_s^0 search region the data exceed the background prediction and have a p-value of 0.27%. The excess is concentrated in bins with $\nu_N > 0.97$. If we restrict ourselves to only the two highest ν_N bins ($\nu_N > 0.987$), which together account for 85% of the signal acceptance, we find a p-value of 0.66%. For the $B_s^0 \to \mu^+\mu^-$ analysis we also produce an ensemble of simulated experiments that includes a $B_s^0 \to \mu^+\mu^-$ contribution at the expected SM branching fraction [2] and yields a p-value of 1.9%. The corresponding p-value for the two highest ν_N bins alone is 4.3%.

We use a modified frequentist approach [20, 21] that includes the effects of systematic uncertainties to calculate expected and observed limits. We calculate expected limits of $\mathcal{B}(B^0 \to \mu^+\mu^-) < 4.6 \times 10^{-9}$ and $\mathcal{B}(B_s^0 \to \mu^+\mu^-) < 1.5 \times 10^{-8}$ at the 95% confidence level (C.L.), a factor 3.3 improvement relative to our previous analysis [4]. We calculate observed limits of $\mathcal{B}(B^0 \to \mu^+\mu^-) < 6.0(5.0) \times 10^{-9}$ and $\mathcal{B}(B_s^0 \to \mu^+\mu^-) < 4.0(3.5) \times 10^{-8}$ at 95% (90%) C.L. If we assume the observed excess in the B_s^0 region is due to signal, we determine $\mathcal{B}(B_s^0 \to \mu^+\mu^-)=(1.8^{+1.1}_{-0.9}) \times 10^{-8}$ using the data $-2\ln Q$ distribution and taking the central value from the minimum and the associated uncertainty as the interval corresponding to a change of one unit. By examining the interval corresponding to a change of 2.71 units we set bounds of $4.6 \times 10^{-9} < \mathcal{B}(B_s^0 \to \mu^+\mu^-) < 3.9 \times 10^{-8}$ at the 90% C.L. As a cross check we use a Bayesian technique to make a point estimate and to derive bounds at 90% C.L. and obtain results very similar to those reported here. Using the central value for the fitted $B_s^0 \to \mu^+\mu^-$ branching fraction we produce an ensemble of simulated experiments and find a p-value of 50%.

The source of the data excess in the $0.970 < \nu_N < 0.987$ bin of the B_s^0 signal region is investigated. The same events, same fits, and same methodologies are used for both the B_s^0 and B^0 searches. Because the data in the B^0 search region shows no excess, problems with the background estimates are ruled out. In particular, the only peaking background in this mass region is from $B \to h^+h'^-$ decays, whose contribution to the B^0 search region is ten times larger than to the B_s^0 search region. Problems with the NN are ruled out by the many studies performed. These NN studies find no evidence of a ν_N- $m_{\mu\mu}$ correlation, no evidence of overtraining, and no evidence of a significant mis-modeling of the ν_N shape, even in the region $0.995 < \nu_N$. In short, there is no evidence that the excess in this bin is caused by a mistake or systematic error in our background estimates or our modeling of the ν_N performance and distribution. The most plausible remaining explanation is that this is a statistical fluctuation. For our central result we use the full set of bins that had been established a priori since this represents an unbiased choice. As discussed above, if we remove the $0.970 < \nu_N < 0.987$ bin the results are not significantly affected.

In summary, we have performed a search for $B^0 \to \mu^+\mu^-$ and $B_s^0 \to \mu^+\mu^-$ decays using 7 fb^{-1} of integrated luminosity collected by the CDF II detector at the Fermilab Tevatron. The data in the B^0 search region are consistent with background expectations and the world's most stringent upper limit on $\mathcal{B}(B^0 \to \mu^+\mu^-)$ is established. The data in the B_s^0 search region are in excess of the background predictions with a p-value of 0.27%. A fit to the data determines $\mathcal{B}(B_s^0 \to \mu^+\mu^-)= (1.8^{+1.1}_{-0.9})\times 10^{-8}$ including all uncertainties.

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