Measurement of b Hadron Lifetimes in Exclusive Decays Containing a J/ψ in pp[over ¯] Collisions at sqrt[s]=1.96 TeV

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We report on a measurement of b-hadron lifetimes in the fully reconstructed decay modes $B^+ \to J/\psi\, K^+$, $B^0 \to J/\psi\, K^*(892)^0$, $B^0 \to J/\psi\, K^0_s$, and $\Lambda^0_b \to J/\psi\, \Lambda^0$ using data corresponding to an integrated luminosity of 4.3 fb^{-1}, collected by the CDF II detector at the Fermilab Tevatron. The measured lifetimes are $\tau(B^+) = [1.639 \pm 0.009$ (stat) $\pm\ 0.009$ (syst)] ps, $\tau(B^0)$ = [1.507 ± 0.010 (stat) $\pm\ 0.008$ (syst)] ps and $\tau(\Lambda^0_b)$ = [1.537 ± 0.045 (stat) $\pm\ 0.014$ (syst)] ps. The lifetime ratios are $\tau(B^+)/\tau(B^0)$ = [1.088 ± 0.009 (stat) $\pm\ 0.004$ (syst)] and $\tau(\Lambda^0_b)/\tau(B^0)$ = [1.020 ± 0.030 (stat) $\pm\ 0.008$ (syst)]. These are the most precise determinations of these quantities from a single experiment.

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The lifetime of ground-state hadrons containing a b quark and lighter quarks is largely determined by the

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charged weak decay of the b quark. Interactions involving the lighter quarks, referred to as spectator processes, alter b-hadron lifetimes at approximately the 10% level. Lifetimes are important to probe our understanding of the low-energy strong interaction. While precise predictions for b-hadron lifetimes are difficult to calculate, ratios are predicted with fairly high accuracy by the Heavy Quark Expansion (HQE) [1]. This framework of theoretical calculation is used to predict low energy QCD effects in many flavor observables. For example, HQE predicts the decay-width of B_s mesons to final states common to B_s^0 and $\bar{B_s^0}$, Γ_{12}^s, which enters the decay-width difference in the B_s^0 system and several CP violation effects. The measurement of lifetime ratios provides a simple and accurate way to test the HQE framework as non standard model effects are expected to be highly suppressed in lifetimes. The ratio $\tau(B^+)/\tau(B^0)$, R_+, (charge conjugates are implied throughout) is predicted to be in the range 1.04-1.08 [1–4]. Predictions for the ratio $\tau(\Lambda_b^0)/\tau(B^0)$, R_Λ, in HQE, which do not presently incorporate next-to-leading order QCD corrections, lie in the range 0.88 ± 0.05 [2, 4, 5]. The first measurements of the Λ_b^0 lifetime have been at the lower end of that range. However, recent high precision measurements by the CDF experiment [6, 7] based on 1.0 fb^{-1} of data, are significantly higher than previous results. It's therefore useful to keep pursuing lifetime measurements with increased precision to settle the issue. In this letter we report precise measurements of b-quark meson lifetimes using the channels $B^+ \to J/\psi\, K^+$, $B^0 \to J/\psi\, K^{*0}$, and $B^0 \to J/\psi\, K_s^0$, in addition to the lifetime of the Λ_b^0 baryon using the $\Lambda_b^0 \to J/\psi\, \Lambda^0$ decay channel. Our data sample corresponds to an integrated luminosity of 4.3 fb^{-1} and consists of p$\bar{\text{p}}$ collisions at a center of mass energy $\sqrt{s} = 1.96$ TeV collected by the CDF II detector at the Fermilab Tevatron. The measurement reported here improves the previous CDF measurement [6] of the Λ_b^0 lifetime by updating it with significantly more data. In all decay modes, the decay position of the b hadron is estimated using only J/ψ decay products so that differences in decay time resolution between channels is reduced and certain systematic uncertainties cancel in ratios of lifetimes.

The components of the CDF II detector relevant to this analysis are described briefly here. Charged particles are reconstructed using six layers of silicon microstrip detectors with radii between 2.4 cm and 23 cm [8] and an open-cell drift chamber called the central outer tracker (COT) [9]. These are immersed in a 1.4 T solenoidal magnetic field and cover the range $|\eta| \leq 1$, where η is the pseudorapidity defined as $\eta \equiv -\ln\tan(\theta/2)$, and θ is the polar angle [10]. Four layers of planar drift chambers (CMU) [11] detect muons with $p_T > 1.4$ GeV/c within $|\eta| < 0.6$. Additional chambers and scintillators (CMX) [12] cover $0.6 < |\eta| < 1.0$ for muons with $p_T > 2.0$ GeV/c.

The reconstruction of b-hadron candidates begins with the collection of $J/\psi \to \mu^+\mu^-$ candidates using a dimuon trigger. The extremely fast tracker (XFT) [13] uses COT hit information to measure the transverse momentum and azimuthal direction of charged tracks. Events with $J/\psi \to \mu^+\mu^-$ candidates are recorded for further analysis if two or more extrapolated tracks are matched to CMU or CMX track segments, opposite-charge and opening-angle requirements are met, and the J/ψ candidate has mass in the range 2.7 to 4.0 GeV/c^2. After offline reconstruction, tracks corresponding to two triggered muon candidates are constrained to originate from a common vertex to make a $J/\psi \to \mu^+\mu^-$ candidate. To ensure a high-quality vertex for the lifetime measurement, each muon track is required to have at least three hits in the silicon system. The reconstructed $\mu^+\mu^-$ invariant mass is required to be in the range $3.014 < m(\mu\mu) < 3.174$ GeV/c^2. The b hadron is assumed to originate from the average beamspot determined as a function of time using inclusive jet data. The primary vertex for a given event is the $x - y$ position of this beamspot at the average z coordinate of the muon tracks at their closest approach to the beamline. The typical beamline size is $\approx 30\ \mu m$ in $x - y$, and this dominates the uncertainty on the decay length. The projection of the transverse decay vector onto the b–hadron p_T direction, L_{xy}, and its uncertainty, σ_{xy}, are also obtained and are used to estimate the proper decay time, $ct = \frac{ML_{xy}}{p_T}$, and its uncertainty σ^{ct}, where M and p_T are the mass and transverse momentum of the b hadron. The uncertainties in the primary and J/ψ vertices, and the transverse momentum are all included in σ^{ct} which has a typical value around 0.1 ps. Uncertainties in transverse momentum have a negligible effect on ct measurement, in comparison to the uncertainty on the vertex positions.

We reconstruct $K^{*0} \to K^+\pi^-$, $K_s^0 \to \pi^+\pi^-$, and $\Lambda^0 \to p\pi^-$ candidates from pairs of oppositely-charged tracks fit to a common vertex. As K_s^0 and Λ^0 decays can occur outside some layers of the silicon system due to their long lifetime, their tracks are not required to have silicon hits. The fitted mass is required to be in a mass window; for the K^{*0} this window is $0.84 < m(K\pi) < 0.96$ GeV/c^2, (the lower range is selected in order to avoid reflections from the $\phi \to K^+K^-$, where one kaon is misreconstructed as a pion), for the K_s^0 it is $0.473 < m(\pi\pi) < 0.523$ GeV/c^2, and for the Λ^0 it is $1.107 < m(p\pi) < 1.125$ GeV/c^2. This corresponds to approximately $\pm 3\sigma$, where σ is the mass resolution of the reconstructed signal. We suppress K_s^0 and Λ^0 cross contamination by rejecting K_s^0 (Λ^0) candidates with proton-pion (pion-pion) invariant mass consistent with Λ^0 (K_s^0). We reconstruct the b-hadrons by performing a kinematic fit of all b-hadron final state tracks to the appropriate topology: two spatially separated vertices in the case of $\Lambda_b^0 \to J/\psi\Lambda^0$ and $B^0 \to J/\psi K_s^0$, one vertex in all other cases. A mass constraint is applied in the J/ψ fit, and the reconstructed momenta of the K_s^0 and Λ^0 are required to point back to the J/ψ vertex. We exclude candidates with $\sigma^{ct} > 100\ \mu m$ to ensure well measured vertices. Additional selection requirements implying consistency with the fit assumptions (common vertex or vertices, mass and pointing constraints) are also applied. Further selection requirements on the transverse momenta of the b-hadrons and daughter particles, invariant mass of the K_s^0, K^{*0}, and Λ^0, the vertex probability of the b-hadrons, and the L_{xy} significance of the K_s^0 and Λ^0 were obtained via an optimization procedure, which maximizes the quantity $\mathcal{S}/\sqrt{\mathcal{S}+\mathcal{B}}$ over all of the selection requirements. The number of signal events ($\mathcal{S}$) is estimated from simulation and the number of background events ($\mathcal{B}$) from the mass sidebands in data. Sidebands are events away from the mass peak and form a sample of pure background.

For B^+ and B^0 modes, only candidates with a reconstructed B mass between 5.17 and 5.39 GeV/c^2 are used for the lifetime measurements. For the Λ_b^0 mode, the mass range is set to $5.48 - 5.76$ GeV/c^2. These ranges provide a sufficient number of events in the sideband regions to constrain the background shape while avoiding regions where the mass distribution has complex structure. The invariant mass distributions for are shown in Fig. 1, where the sideband regions are indicated. The hadron masses are consistent with world average values. We observe the following yields of signal events: 45000 ± 230 (B^+), 16860 ± 140 ($B^0 \to J/\psi\, K^{*0}$), 12070 ± 120 ($B^0 \to J/\psi\, K_s^0$), and $1710 \pm 50(\Lambda_b^0)$.

The lifetimes are extracted using an unbinned maximum likelihood method. The likelihood function $\mathcal{L}$ is multivariate, and is based on the probability of observing a candidate i with reconstructed mass, m_i, decay time, ct_i, decay time uncertainty, σ_i^{ct}, and mass uncertainty, σ_i^m. The PDF is factorized in the following form:

$$\begin{aligned} \mathcal{L} &= \prod_i [f_s \cdot P_m^s(m_i|\sigma_i^m) \cdot T_t^s(ct_i|\sigma_i^{ct}) \cdot S_{\sigma^{ct}}^s(\sigma_i^{ct}) \quad (1) \\ &+ (1-f_s) \cdot P_m^b(m_i) \cdot T_t^b(ct_i|\sigma_i^{ct}) \cdot S_{\sigma^{ct}}^b(\sigma_i^{ct})], \end{aligned}$$

where P_m, T_{ct}, and $S_{\sigma^{ct}}$ and are the normalized probability density functions (PDF) for observables m_i, ct_i and σ_i^{ct}, the superscripts s or b refer to the PDF for signal or background candidates, respectively, and f_s is the fraction of signal events. The PDF $S_{\sigma^{ct}}$ is substantially different for signal and background events and therefore needs to be taken into account as discussed in Ref. [14]. A PDF term for σ_i^m can be ignored as the distribution of σ^m is observed to be similar for both signal and background and hence represents a constant in the log-likelihood.

The signal mass distribution, P_m^s, is modeled as

$$P_m^s = \sum_i \sum_{k=1}^{2} \frac{F_K exp\left(-\frac{(m_0-m_i)^2}{2(u_k\sigma_i^m)^2}\right)}{\sqrt{2\pi}u_k\sigma_i^m}, \qquad (2)$$

where $F_1 + F_2 = 1$, m_0 is the mass of the hadron, and u_k are scale factors to account for the misestimation of the mass resolutions. We find that two Gaussians are sufficient to model the data. The background mass distribution, P_m^b, is modeled as a linear function.

The signal ct distribution, T_{ct}^s, is modeled by an exponential ($e^{-ct_i/c\tau}/c\tau$) convolved with a detailed detector ct-resolution model, $\mathcal{R}$. The background ct distribution, T_{ct}^b, has four components: a $\delta-$function convolved with $\mathcal{R}$ to account for backgrounds from prompt J/ψ's originating from the primary vertex, and one increasing and two decreasing exponentials that account for mismeasured decay vertices and background from other heavy-flavor decays. These exponential components are convolved with a single Gaussian of width σ_i^{ct} multiplied by a scale factor. The relative contribution of each background component is determined by the data. The parameters of the background model are mainly determined from the candidates in the mass sidebands. Studies of inclusive b–hadron decays have shown that after the selection requirements the contamination from other b decays is very low and, furthermore, that the mass distribution of the long lived background components is flat in the fitted mass range, and hence the mass sidebands can provide a realistic background model for candidates in the signal mass range.

The same resolution model, $\mathcal{R}$, is used for signal and prompt background events. The detector resolution is based upon a Gaussian with width $s_j \cdot \sigma_i^{ct}$, where s_j is the scale factor that accounts for the misestimation the σ_i^{ct}. Motivated by a study of resolution in an inclusive sample of J/ψ events, where prompt J/ψ events dominate, $\mathcal{R}$ is modeled as $\mathcal{R} = \sum_{j=1}^{3} f_j/\left(\sqrt{2\pi}s_j\sigma_i^{ct}\right) \cdot \exp\left(-t^2/2(s_j\sigma_i^{ct})^2\right)$, where $f_1 + f_2 + f_3 = 1$. The Gaussians are centered at zero as no evidence of an offset in the data samples is observed. Small differences in $\mathcal{R}$ arise between decay channels due to different χ^2 distributions for the vertex fits of decays with different number of tracks. Therefore the parameters f_j and s_j are obtained separately for each channel from a fit to data in the mass sidebands. This yields an accurate determination of $\mathcal{R}$ since the background events are primarily expected to originate from the interaction vertex.

The functional form of the PDF $S_{\sigma^{ct}}$ is determined empirically using data in the B hadron mass sidebands as indicated on Fig. 1. The parameters of the function, which are different for signal and background are determined from the final fit to data. After the resolution model parameters are determined from the mass-sideband only fit, the likelihood is calculated for each candidate and the product is maximized in each of the four channels to extract the lifetime, signal yield and the other parameters required to describe the mass, background decay time, and σ^{ct} distributions. Decay time projections of the likelihood function are compared with the data in Fig. 2.

We considered correlated and uncorrelated systematic uncertainties. Correlated uncertainties affect all mea-

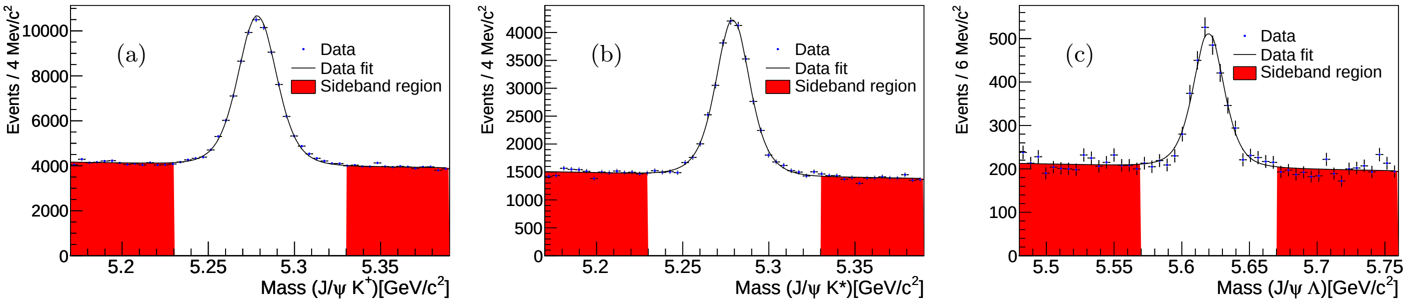

FIG. 1: Invariant mass together with mass fit projection for (a) $B^+ \to J/\psi K^+$, (b) $B^0 \to J/\psi K^*$, and (c) $\Lambda_b^0 \to J/\psi \Lambda^0$.

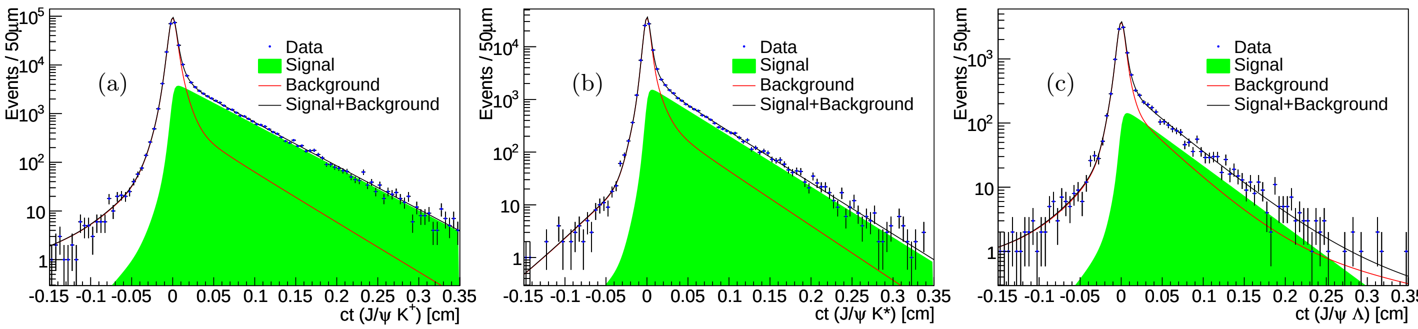

FIG. 2: Decay time distributions for (a) $B^+ \to J/\psi K^+$, (b) $B^0 \to J/\psi K^*$, and (c) $\Lambda_b^0 \to J/\psi \Lambda^0$ candidates.

sured lifetimes identically, and cancel in ratios. We estimate uncertainties due to any residual misalignments of the silicon detector using Monte Carlo samples generated with radial displacements of individual sensors (internal alignment) and relative translation and rotation of the silicon detector with respect to the COT (global alignment). The XFT triggers on tracks assuming they originate from the center of the beam, which may introduce a bias for triggering long-lived decays. No indication of any bias was found in a study of the XFT response in a large sample of simulated events but a small uncertainty is assigned due to the limited statistical precision of the evaluation method. The systematic uncertainty that results from ignoring the correlation between reconstructed mass and σ^{ct} in the likelihood is found to be negligible. The remainder of systematic uncertainties are treated as uncorrelated. They were determined using pseudo-experiments in which many statistical trials are generated according to alternate PDFs where the alternate parameters are derived from data. The shift in data due to the alternate PDFs were consistent with the shift observed with the pseudo-experiments. As the time-resolution is determined from the prompt events, and the shape of those events is sensitive to the modeling of long-lived (positive and negative) background, uncertainties in the background modeling can affect the lifetime through the resolution function. We account for that uncertainty by including an extra long-lived component in the background model. This alternate description produces a substantial change in the fraction of prompt events (approximately 7%), and has a small but non-negligible effect on the lifetime. A further small uncertainty arising from the functional form of $\mathcal{R}$ is also assessed and included in the total resolution uncertainty. To evaluate uncertainties in the mass model, alternate parametrizations, including a 2nd order polynomial for background, and a single Gaussian to describe signal events, were considered. Alternatives to the background PDF included extra long lived and Gaussian components. We determined the uncertainty due to the σ^{ct} parametrization by using a reasonable alternate model. We also considered the effect of ignoring any differences between signal and background mass uncertainties by using distributions determined from data to generate the values of the mass uncertainty in the pseudo-experiments. We also determined the systematic uncertainty due to the presence of the Cabibbo suppressed channel $B^+ \to J/\psi\pi^+$ in the charged B decays, and the effect of swapping the kaon and pion hypotheses in K^{*0} reconstruction. The possibility of a systematic biases caused by the σ^{ct} and p_T selection requirements were found to be negligible. The systematic uncertainties are summarized in Table I.

We measure $\tau(B^+) = [1.639 \pm 0.009(\text{stat}) \pm 0.009(\text{syst})]$ ps and $\tau(B^0) = [1.507 \pm 0.010(\text{stat}) \pm 0.008(\text{syst})]$ ps where the two B^0 measurements have been combined. These results are consistent and of similar precision to the leading measurements from Belle [15] which are $\tau(B^+) = [1.635 \pm 0.011(\text{stat}) \pm 0.011(\text{syst})]$ ps and $\tau(B^0) = [1.534 \pm 0.008(\text{stat}) \pm 0.010(\text{syst})]$ ps. The similarities between the decay channels al-

TABLE I: Summary of systematic uncertainties.

	$J/\psi\, K^+$ (fs)	$J/\psi\, K^{*0}$ (fs)	$J/\psi\, K^0_s$ (fs)	$J/\psi\, \Lambda^0$ (fs)	R_+	R_Λ
Resolution function	2.5	3.5	3.0	8.9	0.0024	0.0061
Background ct model	1.0	2.3	4.1	4.6	0.0017	0.0034
Mass model	2.8	2.8	2.8	2.8	0.0020	0.0017
Proper decay time uncertainty	1.7	1.7	1.7	4.3	0.0010	0.0029
Mass uncertainty	3.0	3.0	3.0	3.0	0.0020	0.0012
Total uncorrelated	± 5.2	± 6.2	± 6.8	± 11.7	0.0042	0.0079
Alignment	6.7	6.7	6.7	6.7	–	–
Cabibbo suppressed mode in B^+	0.7	–	–	–	0.0004	–
Swapped track assignment in B^0	–	0.7	–	–	–	–
Possible trigger bias	1.7	1.7	1.7	1.7	–	–
σ^{ct} -m correlation	0.7	0.7	0.7	0.7	–	–
Total	±8.7	±9.3	±9.7	± 13.7	0.0043	0.0079

low for the accurate determination of the ratio $R_+ = [1.088 \pm 0.009\text{ (stat)} \pm 0.004\text{ (syst)}]$ which favors a slightly higher value than the current average of 1.071±0.009 [2]. These results are consistent with the current HQE predictions, giving further confidence in this theoretical framework, and also provide an accurate test for future lattice QCD calculations. For the Λ^0_b we measure $\tau(\Lambda^0_b) = [1.537 \pm \text{(stat)}0.045 \pm \text{(syst)}0.014]$ ps and $R_\Lambda = [1.020 \pm 0.030\text{(stat)} \pm 0.008\text{(syst)}]$. This measurement is the most precise measurement of $\tau(\Lambda^0_b)$ and is consistent with the previous CDF measurement in this decay channel of $\tau(\Lambda^0_b) = [1.593^{+0.083}_{-0.078}\text{(stat)} \pm 0.033\text{(syst)}]$ ps [6] but is more than 2σ larger than the world average of $1.383^{+0.049}_{-0.048}$ ps and the previous CDF measurement [7], performed on a different decay channel ($\Lambda_c^\pm \pi^\mp$): $[1.401 \pm 0.046\text{(stat)} \pm 0.035\text{(syst)}]$ ps. The ratio is also higher than the predicted values of 0.88±0.95. In summary, we report the most precise determination of $\tau(B^+)/\tau(B^0)$. It is consistent with other measurements and the predicted value which gives confidence in the HQE framework for flavor observables. We also report the most precise measurement of $\tau(\Lambda^0_b)$, which supports a higher value than the world average and theory predictions.

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[1] M. Beneke *et al.*, Nucl. Phys. B **639**, 389 (2002); E. Franco, V. Lubicz, F. Mescia, and C. Tarantino, Nucl. Phys. B **633**, 212 (2002); C. Tarantino, Eur. Phys. J. C **33**, S895 (2004); A. J. Lenz, AIP Conf. Proc. **1026**, 36 (2008).

[2] K. Nakamura *et al.* (Particle Data Group), J. Phys. G **37**, 075021 (2010).

[3] I. Bigi *et al.*, in *B Decays*, 2nd ed., S. Stone (ed.), World Scientific, Singapore (1994).

[4] F. Gabbiani, A. Onischenko, and A. Petrov, Phys. Rev. D **68**, 114006 (2003) and *ibidem* D **70**, 094031 (2004).

[5] N.G. Uraltsev, Phys. Lett. B **376**, 303 (1996); D. Pirjol and N.G. Uraltsev, Phys. Rev. D **59**, 034012 (1999); P. Colangelo and F. De Fazio, Phys. Lett. B **387**, 371 (1996); M. Di Pierro, C.T. Sachrajda, and C. Michael, Phys. Lett. B **468**, 143 (1999).

[6] A. Abulencia *et al.* (CDF Collaboration), Phys. Rev. Lett. **98**, 122001 (2007).

[7] T. Aaltonen *et al.* (CDF Collaboration), Phys. Rev. Lett. **104**, 102002 (2010).

[8] A. Sill *et al.*, Nucl. Instrum. Methods A **447**, 1 (2000).

[9] T. Affolder*et al.*, Nucl. Instrum. Methods A **526**, 249 (2004).

[10] CDF II uses a cylindrical coordinate system in which ϕ is the azimuthal angle, r is the radius from the nominal beamline, and z points in the proton beam direction with the origin at the center of the detector. The transverse plane is the plane perpendicular to the z axis. The $\hat{x}$ axis points outward from the accelerator ring, and the $\hat{y}$ axis points upwards.

[11] G. Ascoli *et al.*, Nucl. Instrum. Methods A**268**, 33 (1988).

[12] T. Dorigo *et al.*, Nucl. Instrum. Methods A **461**, 560 (2001).

[13] E. J. Thomson *et al.*, IEEE Trans. Nucl. Sci. **49**, 1063 (2002).

[14] G. Punzi, arXiv:physics/0401045v1 [physics.data-an] (2004).

[15] K. Abe *et al.* (Belle Collaboration), Phys. Rev. D **71**, 072003 (2005). Erratum-ibid. **71**, 079903 (2005).