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Characterization of blood velocity in arteries using a combined analytical and Doppler imaging approach

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| 6 7 | Characterization of blood velocity in arteries using a combined analytical and doppler imaging approach |
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Abstract

We report a novel experimental and analytical approach to characterize the pulsatile blood flow field based on Doppler ultrasound imaging of the carotid and brachial arteries. The diameter-averaged (DA) velocity, obtained from the instantaneous velocity histograms extracted from the Doppler waveform, was adapted to the solution of a pulsatile flow in a pipe; from which the instantaneous velocity profiles were predicted and compared to local velocity measurements in the carotid and brachial arteries of four healthy human subjects. Very good agreement as demonstrated by the regression slope of 0.97 and near-zero intercept was observed between the spatiotemporal flow field predictions and local velocity measurements at specific distances from the vessel wall. Near-real-time in vivo measurements statistically demonstrate that, the novel analytical and experimental approach presented herein precisely captures the pulsatile blood flow behavior in large blood vessels.

- **Key Words:** Wall Shear Stress, Ultrasound Imaging, Pulsatile flow

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1. Introduction

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Extensive studies have shown that the wall shear stress (WSS) is a major mechanical modulator of many functions in the cardiovascular system. With the assistance of duplex Doppler ultrasound imaging, we aim to develop a comprehensive experimental and analytical approach based on a pipe pulsatile flow to determine the instantaneous velocity field at specific streamwise locations in arteries, as it is the only factor, aside from blood viscosity, that determines WSS.

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1.1. Significance of WSS

Endothelial function relies on the mechanical signals from WSS maintaining cell 57 homeostasis and in adapting to the environmental changes around the endothelial cells 58 of blood vessels. The adaptive response of endothelial cells to the changing shear and 59 stretch load imposed by the flow in blood vessels was examined and it was observed 60 that living cells react to changes in their micro-environments to maintain their "well-61 being", also referred to as "the wisdom of the cell" [1]. The mechanism underlying 62 mechano-transduction by the endothelial glycocalyx layer, a negatively charged layer of 63 proteoglycan, glycoprotein and glycosaminoglycans, was investigated where it was 64 shown that the flexural rigidity of the core proteins is small enough to be compressed 65 easily and yet big enough to make those proteins adequate transducers of the WSS 66 [2,3]. Smooth muscle cell marker genes have been found to be modulated by the WSS 67 following a vascular injury [4]. Complex spatiotemporal WSS in regions where the flow 68 is disturbed near arterial bifurcations lead to atherosclerosis susceptibility [5]. Several 69 70 vascular pathology studies [6-9] have shown low average WSS to be a plaquemodulating factor. 71

The laboratory observations have been supported by several clinical studies that 72 emphasize the importance of monitoring WSS in predicting, characterizing, and treating 73 numerous cardiovascular diseases including aortic dilation or valvular stenosis [10]. 74 dilation and dissection in arteries [11,12], aortic stenosis [13], asymptomatic carotid 75 plaque [14], coronary atherosclerosis [15], arteriovenous malformation (AVM) [16]. 76 Additionally, cognitive impairment in old patients has been found to be correlated with 77 WSS in the common carotid artery [17]. Altered WSS caused by prolonged sitting has 78 been shown to be underlying endothelial dysfunction and impairment of the flow-79 80 mediated dilation of the popliteal artery in legs [18]. A differential regulation of flow and WSS was observed in the carotid and brachial arteries, in response to water immersion 81 of the human body [19]. A new family of cardiovascular risk indicators for the 82 assessment of WSS spatiotemporal patterns, was introduced [20] and proved to be 83 more appropriate than the oscillatory shear index (OSI). 84

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1.2. WSS Assessment

In the context of the discussion above, accurate and quantitative assessment of 86 WSS is of high priority. WSS can be either evaluated based on time-resolved 3D phase 87 contrast Magnetic Resonance Imaging (better known as 4D-MRI) or calculated through 88 image-based Computational Fluid Dynamics (CFD) simulation. In 4-Dimensional 89 Magnetic Resonance Imaging (4D-MRI) [10,12,14,16], the artery is first segmented, 90 local velocity components are measured within the set spatial and temporal resolutions, 91 from which local WSS distributions are estimated. However, the MRI approach may 92 93 underestimate the magnitude of WSS due to the limited spatial resolution [14, 21]. One may argue that increased spatial resolution of MRI can potentially improve the accuracy 94

of WSS measurement, while this would make the approach even more expensive and
 requires long imaging time.

97 As an alternative approach, the CFD methods [14,15] use computerized axial tomography (CAT scans) or Magnetic Resonance Angiography (MRA) to construct the 98 artery geometry specific to each patient and compute the flow field by solving the 99 100 Navier-Stokes equations. However, it also has several limitations in the assessment of pulsatile flow and WSS. The first crucial aspect concerns the quality of reconstructing 101 the geometry based on the image data (CT or MRI). Many studies applied smoothing 102 103 filters to the reconstructed geometries to ensure the quality of volume mesh and computational efficiency [22] which inevitably introduces mismatches between the 104 reconstructed geometry and the true geometry. The second issue is the setup of 105 boundary conditions in the numerical simulation. It is not realistic to simulate the whole 106 cardiovascular system, and thus the regional simulation of blood flow requires 107 appropriate physiological conditions at inlet and outlets [23] Because of this, the 108 physiological parameters applied to the boundaries (pressure difference, volume flow 109 rate, etc.) might be different from the real flow conditions. Besides, the image-based 110 111 CFD is a time-consuming approach, it may take hours of labor to convert medical images to 3D geometries [24] and require a huge amount of computer time to conduct 112 the simulation. 113

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These aforementioned limitations will be mitigated in the current study.

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1.3. Doppler Ultrasound Imaging with Applications to WSS Estimation

With typical spatial and temporal resolutions ranging from 0.5 to 1 mm and 5 to 116 10 ms, respectively, Doppler ultrasound allows monitoring of systolic velocities to detect 117 vascular diseases such as large-artery occlusive disease and fetal compromise [25,26]. 118 Vector ultrasound technique has been used for a better visualization of the blood flow 119 complexity that directly affects WSS [27]. But it is more expensive than the traditional 120 spectral Doppler imaging (still substantially cheaper than MRI). On the other hand, in 121 Ref. [28] C-plane Doppler was used to estimate the blood flow in the lower abdominal 122 vessels of a porcine model. The measurements were performed using a matrix array 123 transducer system. The results were compared to readings from a surgically implanted 124 ultrasonic transit-time flow probe. Good agreement was found between the two 125 techniques. Because the C-plane Doppler is both non-invasive and relatively cheaper, it 126 is adopted in the current study. 127

As a reliable blood flow data acquisition technique, Doppler ultrasound has been 128 coupled with Womersley's analytical solution [29] to a pulsatile flow in a rigid pipe, as 129 well as CFD methods to estimate the velocity profile, WSS, and blood flow rate. For 130 example, Maximum Doppler velocity, based on the maximum spectral velocity detected. 131 has been combined with the Womersley's equations to predict the velocity profile [30-132 31]. Some correction factor was suggested [26] to this method for a better flow rate 133 estimation. However, the correction did not cover WSS estimation. On the other hand, 134 135 computer simulated ultrasound experiments were performed using computational blood flow data to assess the accuracy of Doppler imaging [33]. The computation was done by 136 introducing particles in the flow field to mimic red blood cells, and then computationally 137 generating the Doppler signal that would result from bouncing ultrasound waves off 138

these particles. It was found that mean velocity tracking is more successful than peak 139 velocity tracking. Nevertheless, this method relies heavily on the computational 140 methods. The same limitation is found in other studies [34-36]. Alternatively, vascular 141 phantoms, an apparatus consisting of components that mimic blood vessels embedded 142 under human tissue, have been employed where the Doppler detected centerline 143 144 velocity was used in Womersley's model to obtain the flow rate and shear stress [35-37]. The problem with this method is that it infers the velocity profile based on velocity 145 readings from a very narrow central region in the blood vessel, making the prediction 146 more likely to be thrown off by any fluctuation in the profile shape towards asymmetry. 147

The main limitations to all the above-mentioned studies about Doppler ultrasound WSS estimation, are that they either rely on synthetic Doppler data based on CFD methods, or vascular phantoms. That is on top of the reliance on the centerline velocity as an input to Womersley's model, making both the flow rate and WSS prone to overestimation [36].

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1.4. Focus of the current study

In the present work, these limitations are treated by in vivo, near-real-time 154 validation of the estimated velocity profile that was based on a more appropriate and 155 carefully interpreted input from Doppler data to Womersley's equations. The study 156 integrates the spectral Doppler velocity measurements with the classical Womersley's 157 model to estimate the velocity field and the WSS in arteries with two major differences. 158 First, the calculations are based on the measured instantaneous mean velocity 159 160 averaged across the entire diameter of the artery (DA). That is, instantaneous velocity histograms derived from Doppler waveforms are used rather than just the peak Doppler 161

velocity (VC). The advantage of the DA velocity is that it takes into consideration the 162 entire flow behavior over the cross-section. Under a reasonable assumption that red 163 blood cells are uniformly distributed over the diameter of the vessel, the mean velocity 164 trace of spectral Doppler is translated mathematically into an expression that can be 165 incorporated in Womersley's model. Also, as explained later, the time-averaged mean 166 167 velocity (TAMn) trace was not interpreted to be the flow-averaged velocity that would yield the flow rate when multiplied by the cross-sectional area. Instead, it was treated as 168 the arithmetic mean of the velocity along the diameter, which is not the same as the 169 flow-averaged velocity for a circular cross-section. The second difference is that local 170 velocity measurements were performed to validate the predicted velocity profiles, rather 171 than using a CFD model. As will be shown later in this paper, this novel experimental 172 and analytical approach is statistically proven to be more reliable than using just the 173 centerline velocity. In most of the cases, both temporal and spatial predictions of the 174 velocity profiles agree very well with local measurements. 175

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2. Experiments

2.1. Whole-vessel spectral Doppler

The study was performed on 4 healthy subjects whose ages ranged from 27 to 68 years old. Vascular ultrasound imaging of the brachial and carotid arteries was performed with the subjects lying supine, exposing the arm or neck. Images were acquired with a Zonare ZS3 scanner (ZONARE Medical Systems, Bernardo, CA, USA) using a broadband high resolution L14-5 MHz hockey stick transducer or L14-5 MHz wide linear array transducers. The experimental protocol was approved by the institutional review committee and informed consent for the study was obtained from all human subjects in accordance with the WORLD Medical Association Declaration of
 Helsinki: Ethical principles for medical research involving human subjects, 2008.

The artery was first located using a cross-sectional view (Fig. 1a) and then, 187 switched to a chosen longitudinal plane of the artery. The diameter was measured from 188 the ultrasound image before the Doppler mode was turned on. Since the Doppler effect 189 is due to the head-on component of the velocity relative to the ultrasound beam which 190 would be zero if the angle was 90 degrees, the sample line of Doppler measurement 191 (the short line in Fig 1b) was set in an orientation that was as close to being parallel to 192 193 the flow direction as possible, while the Doppler insonation angle (the angle between the ultrasound beam and the sample line in Fig 1b) was set to be under 60 degrees. 194 195 Since the angle is "known" to the scanner, it is accounted for and the output values 196 would then correspond to the actual axial velocity that we are seeking. The sample gate 197 size was first set to be that of the vessel's diameter. Once the gate size was set, the 198 time-averaged mean velocity (TAMn), systolic peak velocity, and diastolic velocity were recorded. The resulting Doppler waveform is shown in Fig. 1c. 199

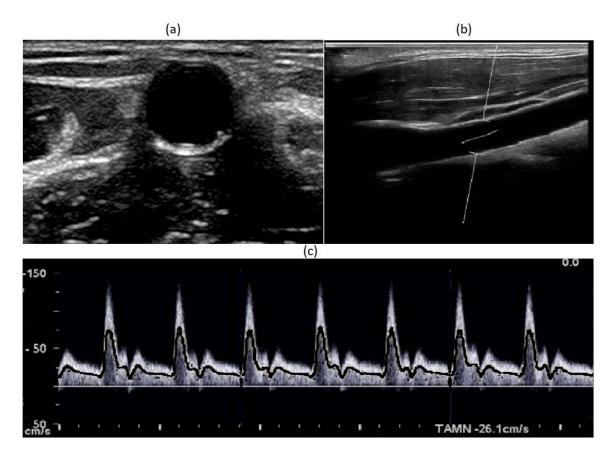


Figure 1: (a), A cross-sectional view of the artery. (b), A longitudinal view of the carotid
 artery with the measurement window spanning the entire diameter (5.8 mm). The
 sample line can be seen halfway across from either side of the measurement window.
 (c), Waveform of the velocity resulting from the measurement shown in (b). x-axis: The
 interval between two consecutive bold markers represent 1 second. y-axis: Velocity
 (cm/s).

The output waveform in Fig. 1c consists of consecutive columns of bright and dim pixels, with one column for every instant in time. The brightness of each individual pixel indicates the relative population of red blood cells (RBCs) having a velocity in the neighborhood of that pixel, as indicated by its vertical position in Fig 1c. The average velocity over the entire population of RBCs that occupy the diameter at the instant, corresponding to a column of pixels, is then computed by the scanner, collapsing that

column into a point belonging to the black trace shown in Fig 1c. Mathematically, theblack trace in Fig. 1c plots the following function of time:

$$f(t) = \frac{\int_{-R}^{R} \rho_{rbc} v_z(r, t) dr}{\int_{-R}^{R} \rho_{rbc} dr}$$
(1)

where *R* is the vessel radius, ρ_{rbc} is the linear density of the RBCs (RBC/m), and v_z is the axial component of the velocity vector. The monitored vessels in this work were the brachial and carotid arteries whose diameters ranged from 2.8 mm to 6.3 mm. At this scale, it is reasonable to assume that the RBC's are evenly distributed throughout the vessels diameter. In other words, ρ_{rbc} in Eq. (1) is a constant. Hence:

$$f(t) = \frac{\int_{-R}^{R} v_z(r, t) dr}{\int_{-R}^{R} dr}$$
⁽²⁾

Allowing for axisymmetry, Eq. (2) is reduced to

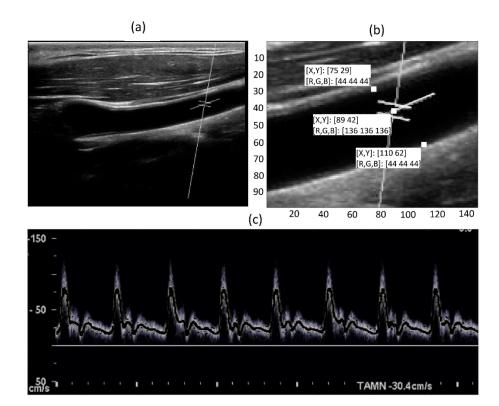
$$f(t) = \frac{1}{R} \int_{0}^{R} v_z(r, t) dr$$
(3)

It is worth noting that f(t) is not the flow-rate-averaged velocity. In other words, multiplying f(t) by the cross-sectional area does not give the flow rate. However, f(t) is the most compact form in which velocity data provided by the scanner could be adapted to the theoretical model.

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2.2. Localized Doppler ultrasound imaging

While at the same cross section, the sample gate size was subsequently reduced to the smallest size possible, and the sample line was placed at several chosen radial locations. An example is shown in Fig. 2a. The small size of the window (1mm) makes it reasonable to assume that the local velocity at that position is the same as the averagevelocity computed by the scanner.



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Figure 2: (a), A longitudinal view of the carotid artery (Diameter: 5.8 mm) with the
 measurement window narrowed down (1 mm) around the location of interest. (b), Using
 pixel coordinates to compute the radial locations of the measurement window's center.
 The figure is obtained by reading the ultrasound image in MATLAB[©]. (c), Waveform of
 the velocity resulting from the measurement shown in (a).

The radial position, was then computed using the coordinates of three aligned pixels as shown in Fig. 2b. If denotes the distance from the upper pixel to the one in the middle, and denotes the distance from the latter to the bottom pixel, then

(4)

It should be noted that the radius R was measured separately, and it is in fact 240 periodically changing with every cardiac cycle. However, the change in the radius is 241 measured to be within less than 2% from the average value. The vessels were therefore 242 treated as rigid tubes in the theoretical section of this paper. It is also noted that, any set 243 of three pixels that include the one in the middle of the local measurement window can 244 be used to deduce the radial location, as long as they are aligned. This is simply done 245 by using similar triangle ratios to work out the expression in Eq. (4). Since a set of three 246 aligned pixels was not always available in such a way that their line was perpendicular 247 to the vessel, the more general case where they are just required to be aligned was 248 dealt with and hence illustrated in Fig. 2. 249

Fig 2c. shows the local TAMn measured by the scanner. This local waveform will be used to compare with the analytically predicted velocity profile for the verification of the theoretical model.

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3. Formulation

The blood flow is modelled as a pulsatile, laminar, incompressible, Newtonian fluid [38] flowing in a rigid tube of radius, R. The flow was assumed to be periodic and fully developed, i.e. zero radial and orthoradial components of the velocity field. A pulsatile pressure gradient driving the flow, takes the following form

$$\frac{\partial p}{\partial z} = -Ce^{i\omega t} \tag{5}$$

where *C* is a constant. It is noticed that only the real part of all mentioned quantities is relevant in the equations. For this problem, the resulting axial velocity is given in Ref. [29]:

$$v_{z}(r,t) = \frac{C}{i\omega\rho} \left[1 - \frac{J_{0}\left(i^{\frac{3}{2}}\alpha \frac{r}{R}\right)}{J_{0}\left(i^{\frac{3}{2}}\alpha\right)} \right] e^{i\omega t}$$
(6)

261 where ω is the radial frequency, and α is the Womersley number given by $\alpha = R \sqrt{\frac{\omega \rho}{\mu}}$.

For any random, but periodic pressure gradient of period *T*, one may decompose it into a Fourier series to fit to the actual pressure gradient. In complex form:

$$-\frac{\partial p}{\partial z}(t) = A_0 + \sum_{n=1}^{N} D_n e^{i\theta_n} e^{i\omega_n t}$$
(7)

264 where $\omega_n = \frac{2\pi n}{T}$.

Making use of the linearity of the problem, and the fact that the A_0 term in Eq. (7) corresponds to the steady component of the pressure gradient, the respective solutions to A_0 and each of the terms inside the summation can be superposed as follows:

$$v_{z}(r,t) = \frac{A_{0}R^{2}}{4\mu} \left(1 - \frac{r^{2}}{R^{2}}\right) + \sum_{n=1}^{N} \frac{C_{n}}{i\omega_{n}\rho} \left[1 - \frac{J_{0}\left(i^{\frac{3}{2}}\alpha_{n}\frac{r}{R}\right)}{J_{0}\left(i^{\frac{3}{2}}\alpha_{n}\right)}\right] e^{i\omega_{n}t}$$
(9)

where:

$$C_n = D_n e^{i\theta_n} \qquad \alpha_n = R_{\sqrt{\frac{\rho\omega_n}{\mu}}}$$
(10)

The first term on the right-hand side of Eq. (9) is the parabolic solution to the Poiseuille's flow driven by a constant pressure gradient:

$$\frac{\partial p}{\partial z} = -A_0 \tag{11}$$

The data provided by the scanner was the average velocity over the span of the measurement window as shown in Figs. 1b and 2a. When the window size is set to be that of the vessel's diameter, the black trace in Fig. 1c represents the f(t) in Eq. (3). Before proceeding, all experimental velocity signals were exactly reproduced with the first 60 harmonics (N = 60) of their Fourier expansions.

276 Since the location where the velocity measurement was performed is far enough 277 downstream from the heart, it is reasonable to assume full hydrodynamic development 278 of the flow. Substituting Eq. (9) into Eq. (3) yields

$$f(t) = \frac{A_0 R^2}{6\mu} + \sum_{n=1}^{N} \frac{C_n}{i\omega_n \rho} (1 - F_n) e^{i\omega_n t}$$
(12)

where:

$$F_n = \int_0^R \left[\frac{J_0\left(i^{\frac{3}{2}}\alpha_n \frac{r}{R}\right)}{RJ_0\left(i^{\frac{3}{2}}\alpha_n\right)} \right] dr$$
(13)

The experimentally obtained f(t) is then broken into a Fourier series:

$$f(t) = A_{f0} + \sum_{n=1}^{N} D_{fn} e^{i\theta_{fn}} e^{i\omega_n t}$$
(14)

Finally, by matching Eq. (12) with Eq. (14), the unknown A_0 and C_n can be obtained:

$$A_{0} = \frac{6\mu A_{f0}}{R^{2}} \qquad C_{n} = \frac{i\omega_{n}\rho D_{fn}e^{i\theta_{fn}}}{1 - F_{n}}$$
(15)

Eq. (15) is then substituted into Eq. (9) to obtain the theoretical prediction of the velocity field. Different from our method outlined in Eqs. (12-15), Refs. [35-37] used the measured centerline velocity to back-calculate the velocity field. Substituting r = 0 in Eq. (9) yields the theoretical centerline velocity $V_c(t)$ as:

$$V_{c}(t) = \frac{A_{0}R^{2}}{4\mu} + \sum_{n=1}^{N} \frac{C_{n}}{i\omega_{n}\rho} \left[1 - \frac{1}{J_{0}\left(i^{\frac{3}{2}}\alpha_{n}\right)} \right] e^{i\omega_{n}t}$$
(16)

The measured centerline velocity $V_{exp}(t)$ is again, broken into a Fourier series:

$$V_{exp}(t) = A_{c0} + \sum_{n=1}^{N} D_{cn} e^{i\theta_{cn}} e^{i\omega_n t}$$
(17)

Similarly, by matching Eq. (16) with Eq. (17) A_0 and C_n can be obtained:

$$A_{0} = \frac{4\mu A_{c0}}{R^{2}} \qquad C_{n} = \frac{i\omega_{n}\rho D_{cn}e^{i\theta_{cn}}}{1 - \frac{1}{J_{0}\left(i^{\frac{3}{2}}\alpha_{n}\right)}}$$
(18)

The problem with this method is that it only accounts for the flow behavior at a very narrow central region in the blood vessel. As evidenced by the comparison performed later, the DA velocity as an input to Womersley's model, leads to a better agreement between predicted and measured local velocities.

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4. Results and discussion

The black trace f(t) in Fig. 1c was broken into a Fourier series according to Eq. (14). Then, the coefficients A_0 and C_n were computed using Eq. (18), which were further substituted into Eq. (9) to obtain the analytical predictions of the velocity profile. The radial positions, at which the local measurements were made, were obtained based on the pixel coordinates as shown in Fig. 2b and calculated using Eq. (4). Each of the calculated radial positions, r, were then substituted into Eq. (9) to obtain the theoretical prediction of the local velocity, which was later compared to the experimental data shown as the black trace of Fig. 2c. A similar procedure is followed to obtain the analytical velocity profiles from the measured centerline velocity which was broken into a Fourier series (Eq. (17)). A_0 and C_n which are obtained from Eq. (18), were substituted in Eq. (9) to get the analytical velocity profile.

Fig. 3 shows the velocity behavior at a specific radial location of the brachial 305 artery, $\frac{r}{R} = 0.047$, throughout a complete cardiac cycle. Different possible values of 306 human blood viscosity were considered, µ= 1 cP, 2.5 cP and 4 cP for computation. It 307 308 shows that the analytical predictions with different values of blood viscosity are very close to each other, and the analytical model agrees well with the experimental data. 309 The observed low sensitivity of the local velocity to the blood viscosity is attributed to 310 the fact that, the measured mean velocity has the highest impact on the velocity 311 distribution; with the same mean velocity, the velocity profile is defined. The good 312 agreement between the analytical model and the experimental data demonstrates the 313 validity of the theory. The slight difference between the analytical model and the 314 experimental measurements is attributed to the fact that cardiac cycles are not perfectly 315 316 rhythmic. In the current study, we assume the blood to be a Newtonian fluid. This assumption is supported by extensive literatures. For example, the Newtonian model of 317 blood was numerically compared with four non-Newtonian models for a pulsatile flow in 318 319 the presence of arterial stenosis [38]. The results showed that the deviation of all four non-Newtonian models from the Newtonian behavior was negligible for the case of 320 pulsatile flow and thus provide justification for the Newtonian assumption. 321

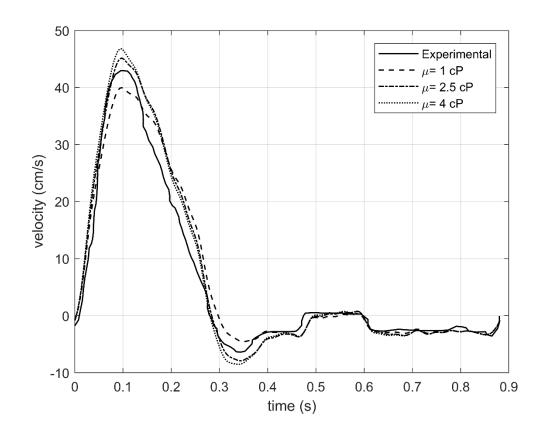
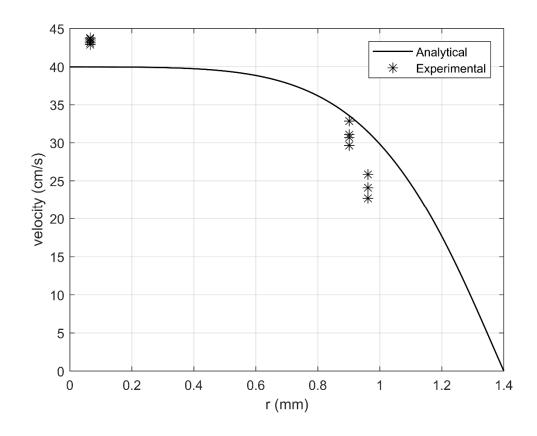


Figure 3: Comparison of the analytical solution for different values of viscosity, with experimental data for the time dependent velocity in the brachial artery at r/R=0.047 where R=1.4 mm

Fig. 4 shows representative results for the comparison between the analytical 326 velocity profile and the experimental data at three radial locations, $\frac{r}{p} = 0.047$, 0.64, and 327 0.69, respectively, in the brachial artery at the peak systole of the cardiac cycle. The 328 329 blood viscosity, μ = 1 cP. The variation of the experimental data at the same location confirms that, cardiac cycles are not completely identical. The analytical model shows 330 that the blood velocity reaches its maximum at the center of the blood vessel, r=0, and 331 gradually decreases with the increase of distance away from the center. The rate of 332 change in the velocity increases sharply near the vessel wall and becomes zero when 333 r= R. Fig. 4 shows that the analytical prediction agrees very well with the experimental 334

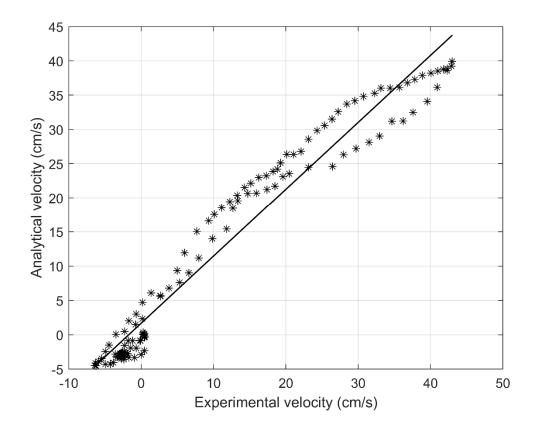
data. The observed slight difference between them might arise from the non synchronized measurement of average and local velocities, as shown in Figs 1 and 2,
 respectively.



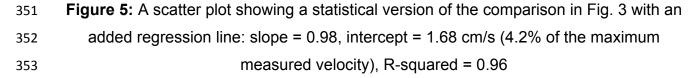
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Figure 4: Comparison of the analytical solution with experimental data at the peak
 systole of the cardiac cycle for 3 radial locations in the brachial artery: r/R = 0.047, 0.64,
 and 0.69 where R=1.4 mm

To further examine the accuracy of the model, a scatter plot of the analytical vs. experimental velocities corresponding to the data of Fig. 3, is shown in Fig. 5. Since the effect of viscosity on the predicted time-dependent velocity was shown to be minimal, the analytical curve corresponding to a viscosity of 1 cP in Fig. 3, was chosen for the comparison in Fig. 5. The slope of 0.98 of the regression line, as well as its small intercept value, 1.68 cm/s (4.2% of maximum measured velocity), show a reliable
 prediction of the experimental velocities with a very slight bias. The tightness of the
 scatter around the regression line is indicated by a high R-squared value of 0.96.



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To demonstrate the better suitability of the DA velocity as a more appropriate input to Womersley's [29] equations, the same procedure described above to obtain the instantaneous local velocity predictions from the acquired DA velocity (Eqs. (12-15)), was applied to the centerline velocity (VC) (Eqs. (16-18)). Figs. 6a and 6b show a sample comparison (The solid straight line is the first bisector (y = x)) of the two

approaches in the brachial artery referred to in Figs. 3-5, where the difference in the 359 error is clearly illustrated by the scatter plots. The DA approach led to an RMS error of 360 9.87 % (Fig 6a), whereas the error corresponding to the VC approach was 17.4% (Fig 361 6b). Fig. 6c shows the wall shear rate, resulting from both methods over one cardiac 362 cycle. An almost 50% difference in the systolic (typically of the most clinical relevance) 363 wall shear rate is observed. In view of the better agreement with experimental data, 364 shown in Fig. 6a, the systolic WSS resulting from the DA method would therefore be 365 considered closer to the actual value, and thus more reliable. 366

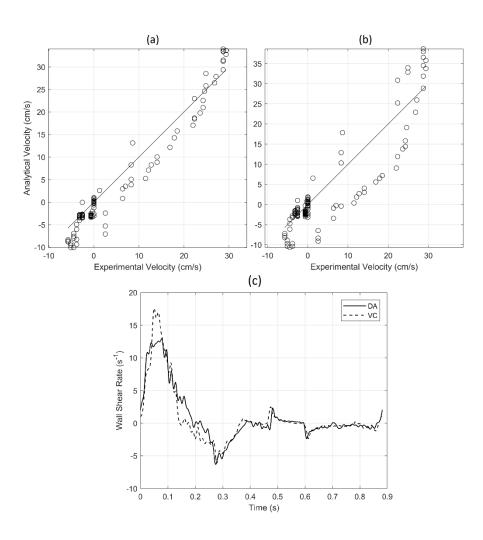


Figure 6: Predicted local velocity versus the one measured experimentally, in the brachial artery at r/R = 0.64 where R = 1.4mm.

(a): using the DA velocity, RMS error = 9.87%, (b): using VC, RMS error = 17.4%370 (c): Wall shear rate where the systolic difference between the two methods is 49.8% 371 The Root Mean Square (RMS) errors and R-squared values for both approaches 372 are compared in Table 1. In most of the cases, the proposed DA approach 373 outperformed the traditional VC approach, in regard to both criteria of RMS errors and 374 R-squared values. Out of the three cases (Br-1, Ca-3, and Ca-6) where the error 375 corresponding to the current DA approach exceeded that of the VC approach, only one 376 (Ca-6) was substantial. For all the cases shown in Table 1, the R-squared values 377 corresponding to the DA approach either exceeded or tied with (in the Br-1 case) the 378 VC approach. This is interpreted as the DA velocity approach having a better 379 predictability for the deviation from experimental results, even when the error is higher 380 than that of the VC approach. An important point needs to be made at this stage. The 381 uncertainties involved in this study apply to both approaches, which validates the 382 conclusion that the presented approach yield a better performance in predicting the 383 velocity profile. 384

Table I. RMS errors and R-Squared values with their corresponding differences,
 resulting from the presented DA velocity approach, and the VC approach.

| | RMS error (%) | | Difference from VC (%) | R-Squared | | Difference from VC (%) |
|------|------------------|-------|---------------------------|-----------|------|---------------------------|
| | DA | VC | | DA | VC | |
| Br-1 | 12.86 | 11.15 | 15.34 | 0.95 | 0.95 | 0.00 |
| Br-2 | 17.09 | 24.3 | -29.67 | 0.96 | 0.8 | 20.00 |
| Br-3 | 9.87 | 17.4 | -43.28 | 0.95 | 0.82 | 15.85 |
| Br-4 | 13.8 | 20.41 | -32.39 | 0.94 | 0.81 | 16.05 |
| Ca-1 | 29.64 | 49.6 | -40.24 | 0.6 | 0.32 | 87.50 |
| Ca-2 | 24.73 | 28.74 | -13.95 | 0.74 | 0.51 | 45.10 |
| Ca-3 | 28.72 | 25.09 | 14.47 | 0.76 | 0.53 | 43.40 |
| Ca-4 | 23.31 | 40.77 | -42.83 | 0.95 | 0.86 | 10.47 |
| Ca-5 | 17.41 | 16.16 | 7.74 | 0.93 | 0.77 | 20.78 |

Ca-6 33.98 22.86 48.64 0.96 0.92 4.35

387

In Fig. 7, the wall shear stress, based on the DA approach applied to the brachial 388 artery of Fig. 3, is plotted versus time throughout a cardiac cycle. The WSS, τ_w profiles 389 were calculated based on the constitutive equation of the blood, $\tau_w = \mu (\partial v_z / \partial r)_{r=R}$ at 390 391 two representative viscosity values, μ = 2.5 cP and 4 cP. The time-averaged absolute 392 WSS corresponding to a viscosity of 2.5 cP, is about 8.51 dyn/cm2; while that corresponding to 4 cP is about 13.62 dyn/cm2. It is important to note that the peak 393 394 systolic WSS determined by our approach in Fig. 7 is consistent with previously reported vascular systolic WSS values of 5.03 N/m2 (50.3 dyn/cm2) for men, and 5.31 395 N/m2 (53.1 dyn/cm) for women [39]. 396

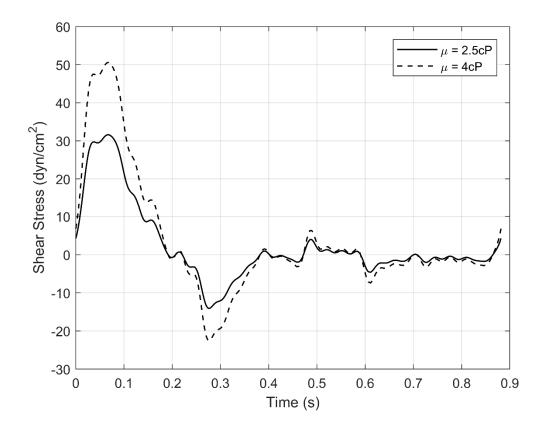
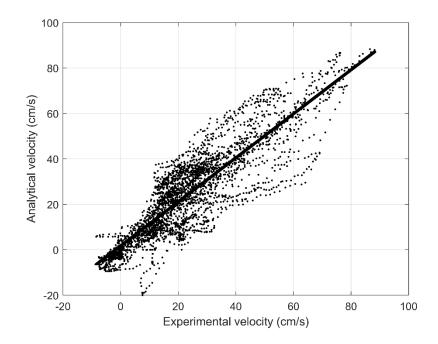


Figure 7: The WSS in the brachial artery referred to in Fig. 3, throughout a cardiac 398 cycle, for two limiting values of the normal range of human blood viscosity 399 Altogether, forty local measurements were performed that cover both brachial 400 and carotid arteries on four healthy subjects. Consistent results similar to the ones 401 presented in Figs. 3, 4 and 5 were obtained. In the appendix, we have included 402 representative results for carotid arteries. It is instructive to present a complete and 403 comprehensive scatter plot which includes the data of Fig. 5 as part of a larger sample 404 covering all the subjects and totaling 3278 data points. The result is shown in Fig. 8. It is 405 clearly shown in this figure that, the data are centered on the regression line with a 406 slope of about 0.97, and a relatively very small intercept value of about 1.71 cm/s which 407 is 1.89 % of maximum measured velocity thereby providing strong evidence in support 408 of the proposed approach. 409



411 **Figure 8:** A scatter plot showing a more comprehensive version of the comparison in

Fig. 5 with an added regression line: slope = 0.97, intercept = 1.71 cm/s (1.89 % of maximum measured velocity), R-squared = 0.79

It is worth noting that, the main purpose behind seeking an accurate prediction of 414 the velocity profile is the subsequent accurate prediction of local WSS, based on the 415 constitutive equation of the blood. Hence, the most important outcome of this paper can 416 be inferred from Fig. 8, where data points that cover all 4 subjects and all considered 417 distances from the vessel wall, are especially tight at very low velocities. In view of the 418 fact that the regions near the wall are low velocity regions, the high robustness of the 419 420 presented technical approach in predicting the flow field near the wall and consequently the WSS, becomes noticeable statistically. 421

The limitations of this study primarily arise from sources of uncertainty in the 422 measurements. The most significant one is the unsynchronized measurements of both 423 average and local velocities. Because the cardiac cycles are not completely identical, 424 matching the local velocity measured at one time with the one theoretically predicted 425 426 based on the average velocity measured at a different time, will inevitably introduce errors. Another source of error arises from the local velocity measurements near the 427 428 vessel wall. Due to the slight periodic change of the diameter some of the surrounding 429 tissue is inevitably included in that window used to measure a local velocity. Other factors that influence the results are non-perfect circular cross-section of the vessels 430 431 (Fig. 1a), as well as the viscosity fluctuations during cardiac cycles since blood is a shear thinning fluid. It should be emphasized that the presented method applies only to 432 relatively straight sections of blood vessels with minimal curvature, where Womersley's 433 solution can be reasonably assumed to capture the flow field. Evidently, there are cases 434

in which the vessel where the WSS needs to be estimated cannot be reliably scanned 435 using this technique. Coronaries for instance, which are continuously moving with the 436 heart's muscle make it impossible to take a steady image even if both the amplitude and 437 frequency of the ultrasound waves are set to reach that depth. Despite of these 438 limitations, the non-invasive, experimental and analytical approach presented herein 439 440 provides an accurate and reliable prediction of the blood velocity field, which would then be useful for predicting the WSS if the blood viscosity is measured separately. WSS has 441 been demonstrated to be a potential marker to identify various cardiovascular diseases. 442 The paper presented herein, precisely capturing the fundamental flow physics in 443 arteries, could be readily translated to clinical applications where Doppler ultrasound 444 imaging is used to estimate WSS. 445

446

5. Conclusion

To recap, in this study, the carotid and brachial arteries of healthy subjects were 447 monitored using the technique of Doppler ultrasound imaging. The diameter of the 448 arteries ranged from 2.8 mm to 6.3 mm. The average velocity over the entire diameter 449 was acquired experimentally, to which the classical Womersley's solution for a pulsatile 450 flow in a rigid pipe was adapted to obtain the theoretical prediction of the velocity profile. 451 Localized Doppler ultrasound imaging was used to obtain the detailed velocity 452 measurements at various radial locations. Very good agreement was observed between 453 the experimental results and the theoretical predictions, especially evidenced by the 454 scatter plot where one finds that, most of the data points are tightly close to the 455 regression line of a slope of 0.97, with a slight bias indicated by the relatively small 456 intercept value of 1.71 cm/s. The same procedure was then repeated with the centerline 457

velocity being the input to Womersley's equations. It was found that using the DA
velocity, rather than the center-line velocity, yields a better prediction for the velocity
field, and consequently the WSS.

461 The novelty and features of this paper are outlined as follows:

• The experimental validation of theoretical velocity predictions was made in nearreal-time, where *in vivo* local velocity measurements were compared to the analytical local velocities.

• The DA velocity was used as an input to the theoretical model to obtain the instantaneous pressure gradient, rather than using the centerline velocity that inevitably brings noise/errors in measurements.

• The TAMn trace was more appropriately interpreted as the DA velocity, rather than the flow averaged velocity that would give the flow rate when multiplied by the vessel's cross-sectional area; which explains the better performance of the presented approach.

The main contributions of this paper are summarized as follows:

The combined analytical and experimental approach provides a novel, much
 more reliable, cheap, and non-invasive method to precisely capture the
 instantaneous velocity profile in large blood vessels, hence allowing for a more
 accurate estimation of WSS.

• The *in vivo* near-real-time local velocity measurements provides a much more realistic description of the blood velocity, as compared to the ones using flow phantoms.

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| 503 | |
|-----|---|
| 504 | APPENDIX |
| 505 | |
| 506 | To illustrate the generality of the cases presented in the main text, an additional |
| 507 | set of data is shown in this appendix. |
| 508 | In Fig. A.1, local velocity predictions were compared to local measurements at |
| 509 | the peak systole for 6 radial locations in the carotid artery of the subject aged 27. The |
| 510 | analytical profile compares very well with the experimental data and following the same |
| 511 | trend. |

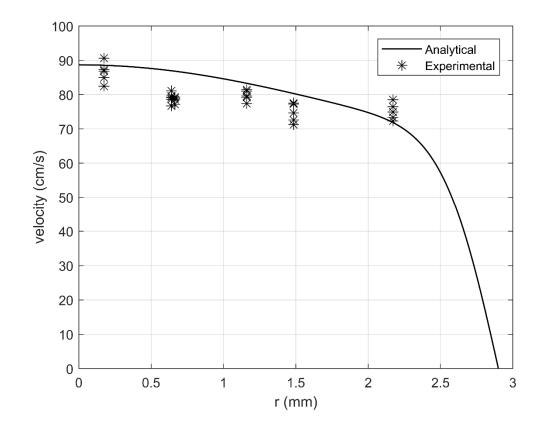
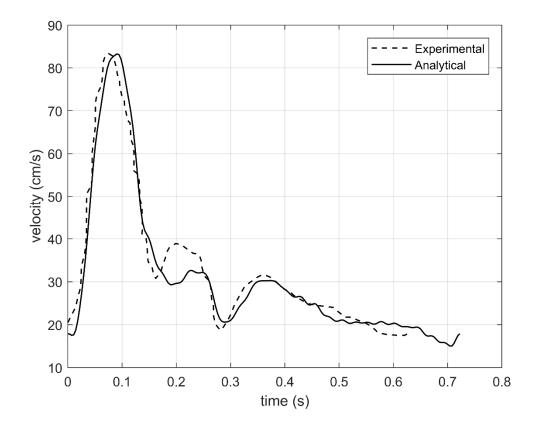


Figure A.1: Comparison of the analytical solution with experimental data at the peak systole of the cardiac cycle for 6 radial locations in the carotid artery: r/R = 0.06, 0.22,0.23, 0.4 0.51, and 0.75 where R=2.9 mm

516 Temporal predictions for the carotid artery are compared with measurements in

Fig A.2, almost half-way between the central axis and the wall. Except around t = 0.2s,

the temporal velocity prediction agrees very well with experimental measurements.



519

520 **Figure A.2:** Comparison of the analytical solution with experimental data for the time 521 dependent velocity in the carotid artery at r/R=0.4 where R=2.9 mm

The comparison in Fig. A.2 is laid on a scatter plot as shown in Fig. A.3. Similarly, to the data in Fig. 6, the regression line has a slope that is close to 1 with a very small bias indicated by the intercept value. The data points are fairly tight around the regression line which is suggested by an R-squared value of about 0.92.

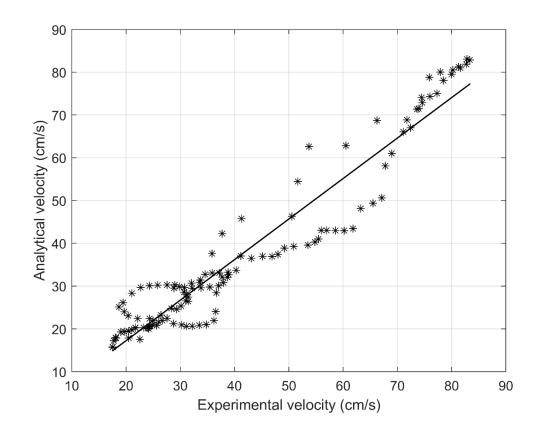


Figure A.3: A scatter plot showing a statistical version of the comparison in Fig. A2 with
 an added regression line: slope = 0.95, intercept = -1.69 cm/s (2.01 % of the maximum
 measured velocity), R-squared = 0.92

In Fig. A.4, the wall shear stress based on the predicted instantaneous velocity profile corresponding to the carotid artery of Fig. A.2, is plotted versus time throughout a cardiac cycle. The time-averaged absolute WSS corresponding to a viscosity of 2.5 cP, is about 17.78 dyn/cm²; while that corresponding to 4 cP is about 28.45 dyn/cm²

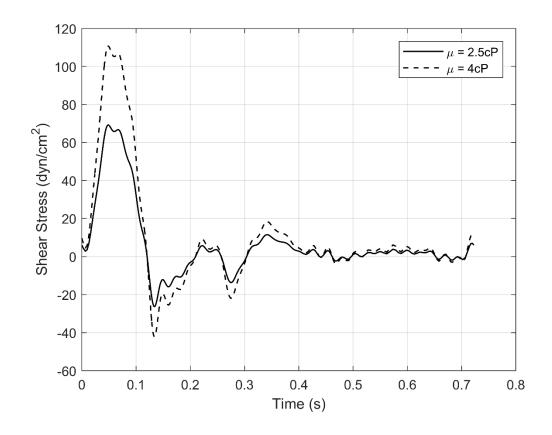


Figure A.4: The WSS in the carotid artery referred to in Fig. A2, throughout a cardiac cycle, for two limiting values of the normal range of human blood viscosity

537 For each local measurement (r/R), systolic velocity data were collected and 538 averaged. The standard deviation (SD) for each, is presented as a percentage from that 539 corresponding average. The results are shown below:

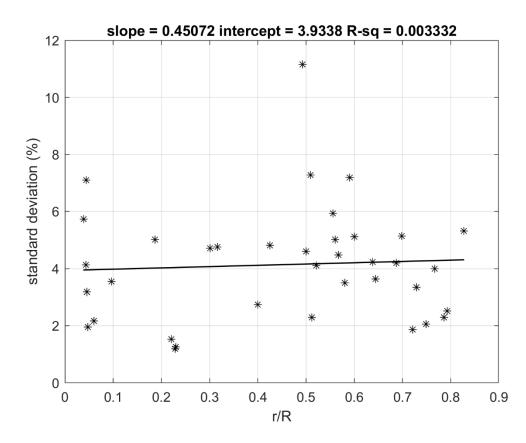


Figure A.5: Standard deviation from the average measured systolic velocities as a function of dimensionless radial position

The mean SD is about 4.13 %. The slope of the regression line shows a flat to very small sensitivity to the dimensionless radial location increasing by 0.45% from centerline to wall.

Modeling the spread distribution in SD to be Gaussian this would result in a 95%

confidence interval for the standard deviation to be, between 0.24% and 8.01%.

Applying the same interval to the entire cardiac cycle would then result in that same interval of uncertainty, for the wall shear stress.

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