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$\Upsilon(nS)$ and $\chi_b(nP)$ production at hadron colliders in nonrelativistic QCD

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$\Upsilon(nS)$ and $\chi_b(nP)$ (n=1,2,3) production at the LHC is studied at next-to-leading order in α_s in nonrelativistic QCD. Feeddown contributions from higher χ_b and Υ states are all considered for lower Υ cross sections and polarizations. The long distance matrix elements (LDMEs) are extracted from the yield data, and then used to make predictions for the $\Upsilon(nS)$ polarizations, which are found to be consistent with the measured polarization data within errors. In particular, the $\Upsilon(3S)$ polarization puzzle can be understood by a large feeddown contribution from $\chi_b(3P)$ states. Our results may provide a good description for both cross sections and polarizations of prompt $\Upsilon(nS)$ and $\chi_b(nP)$ production at the LHC.

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I. INTRODUCTION

Since the surprisingly large production rate of ψ' at large p_T was found by CDF in 1992 [1], the production of heavy quarkonium at hadron colliders has been a problem full of puzzles. While the color-octet (CO) mechanism [2] at leading order (LO) in nonrelativistic QCD (NRQCD) factorization [3] might explain the large production rates of ψ' and J/ψ at large p_T via gluon fragmentation, the predicted transverse polarizations for $J/\psi(\psi')$ were in contradiction with the measurements that the produced $J/\psi(\psi')$ were almost unpolarized (see Ref. [4] for a comprehensive review). In recent years, significant progress has been made in the next-to-leading order (NLO) QCD calculations in NRQCD. Calculations and fits for both yield and polarization in J/ψ production are performed by three groups [5–7], but the conclusions are quite different. In Ref.[6] a simultaneous description for the observed J/ψ yield and polarization can be achieved at large p_T (>7 GeV) by considering possible cancelations between contributions of S- and P-wave color-octet channels. Recently, by including leading power fragmentation corrections, which improves the convergence of α_s expansion at large p_T, a good explanation for the J/ψ polarization is also found[8].

Recently, polarizations of $\Upsilon(1S, 2S, 3S)$ have been measured by CMS at the LHC [9]. It is interesting to study the Υ production within the same framework as that for the J/ψ production and further test the interpretation for the polarization puzzle in Ref.[6]. Note that Υ should be a more suitable system than J/ψ to apply NRQCD, since both v (the relative velocity of heavy quarks in heavy quarkonium) and α_s are smaller for bottomonium than charmonium, and thus the double expansion in α_s and v should converge faster for bottomonium production. Earlier studies of Υ and χ_b production can be found in Refs. [10–13] and references therein. In Ref.[14], a NLO calculation of $\Upsilon(1S, 2S, 3S)$ polarizations is given, where the polarizations for $\Upsilon(1S, 2S)$ agree with the CMS measurements [9], but the predicted ratio of differential cross sections of $\chi_{b2}(1P)$ to $\chi_{b1}(1P)$[14] is too large and inconsistent with the CMS data[15]. Furthermore, without considering the $\chi_b(3P)$ feeddown, the polarization data of $\Upsilon(3S)$ can not be explained [14].

Recently, the radiative transition of $\chi_b(3P)$ to $\Upsilon(3S)$ was first seen by LHCb [16]. So the explanation of $\Upsilon(1S, 2S)$ and $\Upsilon(3S)$ polarizations should be reconsidered, and a proper treatment for $\chi_b(1P, 2P, 3P)$ feeddown is needed, since the treatment of $\chi_b(3P)$ and $\Upsilon(3S)$ will affect the production of $\Upsilon(1S, 2S)$ through the cascaded effects. In this work, we study the prompt production of $\Upsilon(1S, 2S, 3S)$ with both direct and feeddown contributions at NLO in α_s in NRQCD.

The polarized cross section for a bottomonium H can be factorized as [3]

$$d\sigma_{s_z,s_z} = \sum_{i,j,n} \int dx_1 dx_2\, G_{i/p} G_{j/p} \langle \mathcal{O}^H_n \rangle\, d\hat{\sigma}^{i,j,n}_{s_z,s_z}, \quad (1)$$

where p denotes either proton or anti-proton, $G_{i,j/p}$ are the parton distribution functions (PDFs) of p, and the indices i, j run over all the partonic species. $\langle \mathcal{O}^H_n \rangle$ is the long distance matrix element (LDME), with "n" denotes the color, spin and angular momentum of the intermediate $b\bar{b}$ pair, which can be ${}^3S_1^{[1,8]}$, ${}^1S_0^{[8]}$ and ${}^3P_J^{[8]}$ for Υ, and ${}^3P_J^{[1]}$ and ${}^3S_1^{[8]}$ for χ_b. The yield can be obtained by summing the polarized cross sections over the spin quantum number s_z. The virtual corrections are calculated by using our Mathematica code [6, 17, 18], and the real corrections are obtained by using the HELAC-Onia program [19]. We further use the CTEQ6L1 and CTEQ6M PDFs [20] respectively for LO and NLO calculations. The bottom quark mass is set to be $m_b = 4.75$ GeV, the renormalization, factorization, and NRQCD scales are $\mu_r = \mu_f = \sqrt{p_T^2 + 4m_b^2}$ and $\mu_\Lambda = m_b$.

II. FEEDDOWN AND $\chi_b(nP)$

For Υ the polarization observable λ_θ can be expressed as $\lambda_\theta = \frac{\mathrm{d}\sigma_{11}-\mathrm{d}\sigma_{00}}{\mathrm{d}\sigma_{11}+\mathrm{d}\sigma_{00}}$, where σ_{00} and σ_{11} are polarized prompt cross sections, including both direct production and feeddown contributions from higher $\Upsilon(nS)$ and $\chi_b(nP)$ states. Since the transitions between $\Upsilon(nS)$ are dominated by the S-wave dipion modes, the feeddown of higher $\Upsilon(nS)$ will inherit the spin index of the mother particles. While for the $\chi_b(nP)$ feeddown, which proceeds mainly through $\chi_b(nP) \to \Upsilon(mS)\gamma$, the general inheritance relations of polarizations are given in Ref. [21]:

$$\begin{aligned} \lambda_\theta^{\chi_{b0}\to\Upsilon} &= 0, \\ \lambda_\theta^{\chi_{b1}\to\Upsilon} &= \frac{\mathrm{d}\sigma_{00}^{\chi_{b1}} - \mathrm{d}\sigma_{11}^{\chi_{b1}}}{3\mathrm{d}\sigma_{11}^{\chi_{b1}} + \mathrm{d}\sigma_{00}^{\chi_{b1}}}, \\ \lambda_\theta^{\chi_{b2}\to\Upsilon} &= \frac{6\mathrm{d}\sigma_{22}^{\chi_{b2}} - 3\mathrm{d}\sigma_{11}^{\chi_{b2}} - 3\mathrm{d}\sigma_{00}^{\chi_{b2}}}{6\mathrm{d}\sigma_{22}^{\chi_{b2}} + 9\mathrm{d}\sigma_{11}^{\chi_{b2}} + 5\mathrm{d}\sigma_{00}^{\chi_{b2}}}. \end{aligned} \tag{2}$$

Similar to χ_{cJ}[31], at NLO in α_s the χ_{bJ} production is determined by the color-octet (CO) ${}^3S_1^{[8]}$ and color-singlet (CS) ${}^3P_J^{[1]}$ contributions. If CO ${}^3S_1^{[8]}$ is dominant, which leads to transverse polarization at large p_T, the ratios of polarized cross sections become $\mathrm{d}\sigma_{00}^{\chi_{b1}} : \mathrm{d}\sigma_{11}^{\chi_{b1}} = 2:1$ and $\mathrm{d}\sigma_{00}^{\chi_{b2}} : \mathrm{d}\sigma_{11}^{\chi_{b2}} : \mathrm{d}\sigma_{22}^{\chi_{b2}} = 1/3 : 1/2 : 1$, and the feeddown polarization parameters in Eq. (2) are 0.20 for χ_{b1} and 0.29 for χ_{b2}. Further including the CS ${}^3P_J^{[1]}$ contribution only slightly change the overall polarization of χ_{bJ} feeddown. This shows that the χ_b feeddown contribute a modest transverse polarization for Υ at large p_T.

The CS LDMEs for $\chi_{bJ}(nP)$ can be related to the derivatives of radial wave functions at the origin by

$$\langle \mathcal{O}^{\chi_{bJ}(nP)}({}^3P_J^{[1]})\rangle = (2J+1)\frac{3}{4\pi}|R'_{nP}(0)|^2, \tag{3}$$

where $|R'_{nP}(0)|^2$ can be estimated in potential models. E. g. the B-T potential model[22] gives $|R'_{1P,2P,3P}(0)|^2 = (1.417,\ 1.653,\ 1.794)$ GeV^5. In fact, various potentials in Refs.[22] and [23] all indicate $|R'_{1P}(0)|^2 \approx |R'_{2P}(0)|^2 \approx |R'_{3P}(0)|^2$. So, as a balanced approximation, we use

$$|R'_{nP}(0)|^2 \approx 1.653\ \text{GeV}^5, \quad n = 1,2,3, \tag{4}$$

as input. The CO LDMEs are introduced via the ratio

$$r_{nP} = m_b^2 \langle \mathcal{O}^{\chi_{bJ}(nP)}({}^3S_1^{[8]})\rangle / \langle \mathcal{O}^{\chi_{bJ}(nP)}({}^3P_J^{[1]})\rangle, \tag{5}$$

which is independent of J since $\langle \mathcal{O}^{\chi_{bJ}(nP)}({}^3S_1^{[8]})\rangle = (2J+1)\langle \mathcal{O}^{\chi_{b0}(nP)}({}^3S_1^{[8]})\rangle$. Unlike the CS LDMEs, r_{nP} can not be estimated from potential models, but should be extracted from experimental data.

We also assume that the total decay widths of $\chi_{bJ}(nP)$, which are related to $|R'_{nP}(0)|^2$, are approximately independent of n. Then, taking the partial decay widths of $\chi_{bJ}(nP) \to \Upsilon(mS)\gamma$ calculated in Ref.[23] and the PDG values of $\mathrm{Br}(\chi_{bJ}(1P) \to \Upsilon(1S)\gamma)$ [25] as inputs, we can calculate the branching ratios $\mathrm{Br}(\chi_{bJ}(2P) \to \Upsilon(2S)\gamma)$ and $\mathrm{Br}(\chi_{bJ}(2P) \to \Upsilon(1S)\gamma)$, which are found to be close to their PDG values[25], as shown in Tab. I. This implies that it may be a good approximation that the total widths of $\chi_b(nP)$ are independent of n. The above approximation is also roughly consistent with the recent calculations based on the potential model in [24]. With this approximation we further calculate $\mathrm{Br}(\chi_{bJ}(3P) \to \Upsilon(1S,2S,3S)\gamma)$, which are listed in Tab. II.

Br	$theory$	$Experiment$[25]
$\chi_{b1}(2P) \to \Upsilon(2S)$	15.6%	$19.9 \pm 1.9\%$
$\chi_{b1}(2P) \to \Upsilon(1S)$	9.7%	$9.2 \pm 0.8\%$
$\chi_{b2}(2P) \to \Upsilon(2S)$	8.3%	$10.6 \pm 2.6\%$
$\chi_{b2}(2P) \to \Upsilon(1S)$	7.3%	$7.0 \pm 0.7\%$

TABLE I: Predicted branching ratios $\mathrm{Br}(\chi_{b1,b2}(2P) \to \Upsilon(1S,2S)\gamma)$ by assuming the total decay widths of $\chi_{bJ}(nP)$ are independent of n, as compared with experiments[25].

Br	$n=1$	$n=2$	$n=3$
$\chi_{b0}(3P) \to \Upsilon(nS)$	0.24%	0.22%	0.50%
$\chi_{b1}(3P) \to \Upsilon(nS)$	3.81%	3.68%	10.44%
$\chi_{b2}(3P) \to \Upsilon(nS)$	1.92%	1.91%	6.11%

TABLE II: Predicted branching ratios $\mathrm{Br}(\chi_{bJ}(3P) \to \Upsilon(1S,2S,3S)\gamma)$ by assuming the total decay widths of $\chi_{bJ}(nP)$ are independent of n.

III. PROMPT $\Upsilon(nS)$ PRODUCTION

Having clarified how to treat the feeddown contributions, we now extract LDMEs of $\Upsilon(nS)$ and $r(nP)$ defined in (5) by fitting the yield data at LHC, and leave polarizations as our prediction. Data in our fit includes: (1) Differential cross sections of $\Upsilon(nS)$ measured by ATLAS[26] and CMS[27]; (2) Fractions of $\Upsilon(nS)$ production originating from $\chi_b(nP)(n=1,2,3)$ freedown contributions measured by LHCb [16] which are denoted as $R^{\chi_b(nP)}_{\Upsilon(mS)}$ (values for $m \neq n$ are not included in the fit but predicted by using the branching ratios in TABLEs 1 and 2 and compared with data, as shown in Fig.2); (3) Cross section ratio of $\chi_{b2}(1P)$ to $\chi_{b1}(1P)$ measured by CMS [15]. To avoid potential non-perturbative effects in the sense that only the first two powers in the $1/p_T^2$ expansion of cross sections are proven to be factorizable [28], we need to introduce a relatively large p_T cutoff for the data (for the similar case in the production of $\psi^{(\prime)}$, see Ref. [17, 18, 29]). In our fit, we only use data in the region $p_T > 15$ GeV because the $\chi^2/d.o.f.$ will increase quickly when the p_T cutoff becomes smaller than 15 GeV. For example, by choosing the p_T cutoff to be $7, 9, 11, 13, 15$, and 17 GeV, the corresponding $\chi^2/d.o.f.$ in fitting $\Upsilon(3S)$ data are $4.2, 4.0, 2.5, 1.9, 1.3$, and 1.0, respectively.

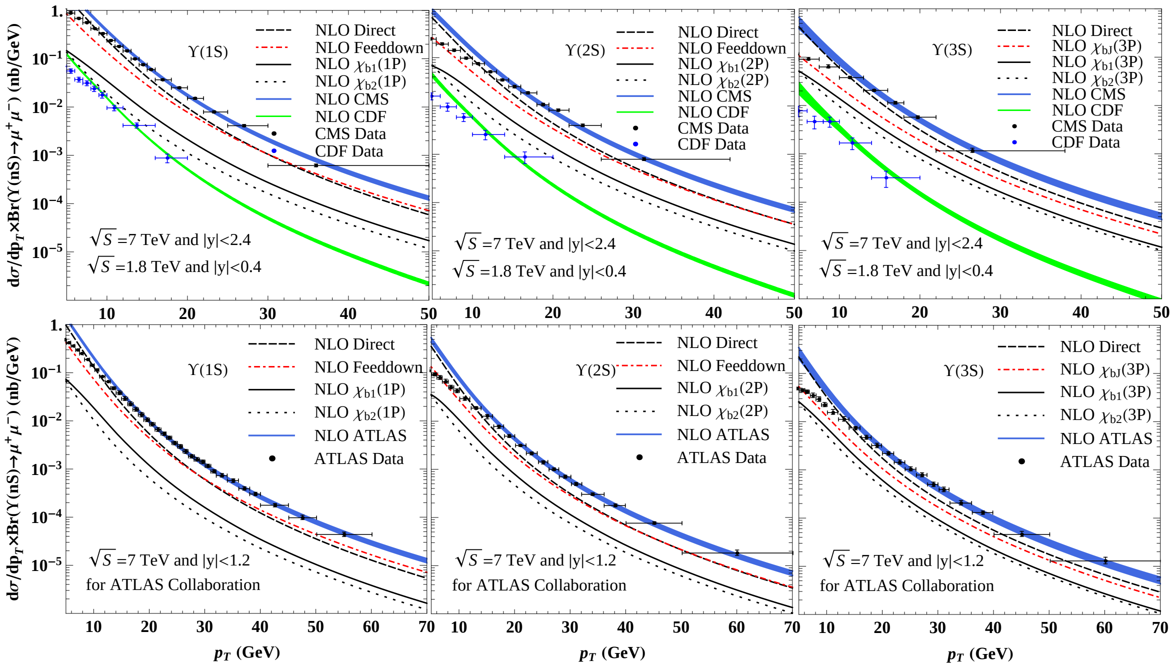

FIG. 1: Differential p_T cross sections for the experimental windows of ATLAS, CMS and CDF. From left to right: $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$. The contributions from direct production are denoted by dashed lines, while those from feeddown by dashed-dotted lines. The $\chi_{b1}(nP)-\Upsilon(nS)$ and $\chi_{b2}(nP)-\Upsilon(nS)$ feeddown contributions are denoted by the solid and dotted lines, respectively. The experimental data are taken from Refs. [26, 27, 30].

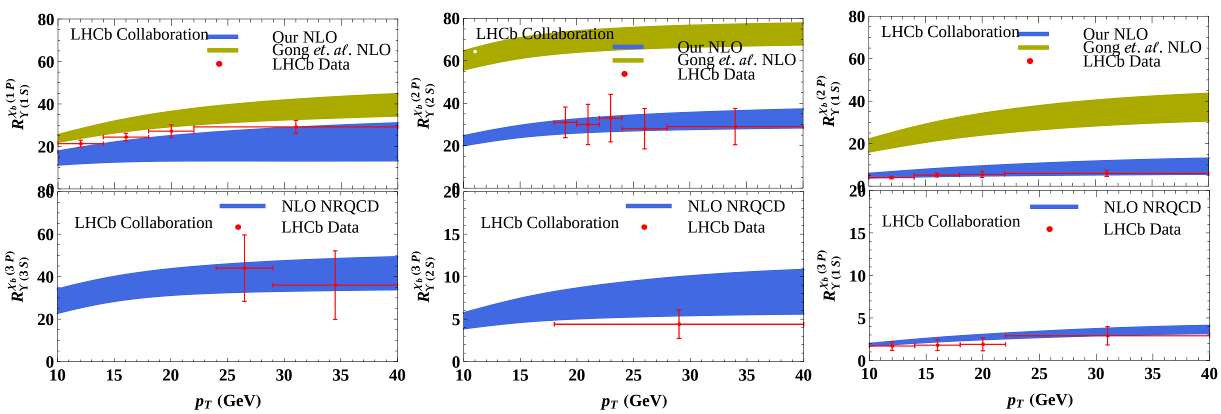

FIG. 2: The fractions of $\Upsilon(mS)$ $(m=1,2,3)$ production originating from $\chi_b(nP)$ $(n=1,2,3;\ n\geq m)$ feeddown contributions, denoted as $R^{\chi_b(nP)}_{\Upsilon(mS)}$ (in units of percentage). From left to right: $R^{\chi_b(1P)}_{\Upsilon(1S)}$, $R^{\chi_b(2P)}_{\Upsilon(2S)}$, $R^{\chi_b(2P)}_{\Upsilon(1S)}$ in the first row and $R^{\chi_b(3P)}_{\Upsilon(3S)}$, $R^{\chi_b(3P)}_{\Upsilon(2S)}$, $R^{\chi_b(3P)}_{\Upsilon(1S)}$ in the second row. Our predictions are denoted by the blue bands, while those obtained by using parameters in Ref.[14] are denoted by the yellow bands. Experimental data are taken from Ref.[16].

When $p_T > 15$ GeV, we find the CO P-wave ${}^3P_J^{[8]}$ contribution can be decomposed into a linear combination of ${}^1S_0^{[8]}$ and ${}^3S_1^{[8]}$ (just similar to the J/ψ case [17, 18]),

$$\mathrm{d}\hat{\sigma}({}^3P_J^{[8]}) = r_0\ \mathrm{d}\hat{\sigma}({}^1S_0^{[8]}) + r_1\ \mathrm{d}\hat{\sigma}({}^3S_1^{[8]}), \tag{6}$$

where $r_0 = 3.8$, $r_1 = -0.52$, which may slightly change with rapidity ranges. So with three CO LDMEs we can extract two linear combinations, which are denoted by

$$M_{0,r_0}^{\Upsilon(nS)} = \langle \mathcal{O}^{\Upsilon(nS)}({}^1S_0^{[8]})\rangle + \frac{r_0}{m_b^2}\langle \mathcal{O}^{\Upsilon(nS)}({}^3P_0^{[8]})\rangle, \tag{7}$$

$$M_{1,r_1}^{\Upsilon(nS)} = \langle \mathcal{O}^{\Upsilon(nS)}({}^3S_1^{[8]})\rangle + \frac{r_1}{m_b^2}\langle \mathcal{O}^{\Upsilon(nS)}({}^3P_0^{[8]})\rangle,$$

which account for $1/p_T^6$ and $1/p_T^4$ behaviors, respectively.

	$\langle\mathcal{O}(^3S_1^{[1]})\rangle$ GeV^3	M_{0,r_0} 10^{-2}GeV^3	M_{1,r_1} 10^{-2}GeV^3
$\Upsilon(1S)$	9.28	13.70 ± 1.11	1.17 ± 0.02
$\Upsilon(2S)$	4.63	6.07 ± 1.08	1.08 ± 0.20
$\Upsilon(3S)$	3.54	2.83 ± 0.07	0.83 ± 0.02

TABLE III: The LDMEs for $\Upsilon(1S, 2S, 3S)$ production. The combined LDMEs are obtained by the fit, while the CS ones are estimated by using the $B-T$ potential model in Ref.[22].

Based on the above method, we fit two linear combinations $M_{0,r_0}^{\Upsilon(nS)}$ and $M_{1,r_1}^{\Upsilon(nS)}$ for $\Upsilon(1S, 2S, 3S)$ with $\chi^2/\text{d.o.f} = 0.99, 2.07, 1.25$, together with CS LDMEs that are estimated by using the B-T potential model [22](see Tab.III). As for $r(nP)$, the results are listed in Tab.IV, with those obtained in Ref. [14] for comparison. In

$r(nP)$	$n=1$	$n=2$	$n=3$
This work	0.42 ± 0.05	0.62 ± 0.08	0.83 ± 0.22
Ref. [14]	0.85 ± 0.11	1.58 ± 0.38	

TABLE IV: The values of $r(nP)$ for $n = 1, 2, 3$ in this work and in Ref. [14].

Tab. III, we find that the central value of $M_{0,r_0}^{\Upsilon(nS)}$ decrease more quickly than that of $M_{1,r_1}^{\Upsilon(nS)}$ as n increases, while the values of $M_{1,r_1}^{\Upsilon(nS)}$ almost have no changes. This explains why a higher $\Upsilon(nS)$ tends to have a less steep p_T cross sections.

Comparisons between our fit and data are shown in Figs. 1, 2 and 3, along with our postdiction for the CDF cross section [30]. It is interesting to see that the yield, fractions of $\Upsilon(mS)$ production from $\chi_b(nP)$ decays, and cross section ratios for $\Upsilon(1S, 2S, 3S)$ can be well described simultaneously. In particular, good agreement with $R_{\Upsilon(nS)}^{\chi_b(3P)}$ is achieved explicitly by a relatively large feeddown contribution from $\chi_b(3P)$, as indicated by the large value of $r(3P)$ in Tab. IV. For comparison, we also present the fractions $R_{\Upsilon(mS)}^{\chi_b(nP)}$ using the parameters in Ref. [14], which are shown in Fig. 2 as the yellow bands. From Fig. 2, one see that the $\chi_b(1P, 2P)$ production rates predicted by Ref. [14] are too large compared with data, whereas our predictions of the production rates of $\chi_b(1P, 2P)$ and $\chi_b(3P)$, denoted by the blue bands in Fig. 2, are roughly consistent with data. In Fig. 3, with the extracted value of r_{1P} in Tab. IV we can well describe the measured ratio of differential cross sections of χ_{b2} to χ_{b1} by CMS [15], clearly better than that in Ref.[14].

With the LDMEs extracted from yield data, we can calculate the $\Upsilon(nS)$ polarizations. The predicted λ_θ of $\Upsilon(1S, 2S, 3S)$ are the weighted averages of the direct production and feeddown contributions. This can be seen directly from Fig. 4, where the results for the CMS window at $\sqrt{S} = 7$ GeV are shown. The predictions for prompt $\Upsilon(1S, 2S, 3S)$ polarizations are roughly consistent with data. Note that the $\Upsilon(3S)$ polarization is obtained with a relatively large feeddown contribution from $\chi_b(3P)$ (see the feeddown fraction $R_{\Upsilon(3S)}^{\chi_b(3P)}$ shown in Fig. 2), which reduces the value of λ_θ of direct production and leads to a smaller total polarization λ_θ of prompt $\Upsilon(3S)$. The feeddown contributions also affect the $\Upsilon(1S, 2S)$ polarizations and lead to better agreement with data.

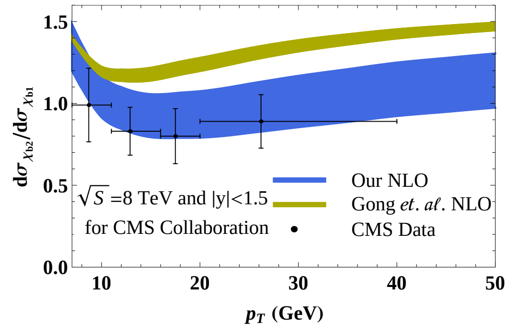

FIG. 3: The ratio of differential cross sections of $\chi_{b2}(1P)$ to $\chi_{b1}(1P)$ for the experimental windows of CMS. The blue band is our NLO results with the extracted value of r_{1P} in Tab. IV and the yellow band is obtained by using parameters in Ref.[14]. Experimental data are taken from Ref.[15]

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In fact, the predicted λ_θ's of the prompt $\Upsilon(1, 2, 3S)$ are the weighted averages of the contributions from direct production and feeddown processes. This can be seen from Fig. 4. In particular, for the λ_θ of $\Upsilon(3S)$, the weight of feeddown contribution is just the fraction $R_{\Upsilon(3S)}^{\chi_b(3P)}$ shown in Fig. 2, which is as large as about 40%, as observed by LHCb [16]. Since the fraction $R_{\Upsilon(3S)}^{\chi_b(3P)}$ is determined by the product of the $\chi_b(3P)$ production cross section and the branching ratio of $\chi_b(3P) \to \Upsilon(3S)\gamma$, a change of the branching ratio will cause a change of $\chi_b(3P)$ production cross section but keep the fitted fraction $R_{\Upsilon(3S)}^{\chi_b(3P)}$ unchanged. Namely, the uncertainty in the predicted branching ratio in Tab. II will affect the predicted value of $\chi_b(3P)$ cross section but not $R_{\Upsilon(3S)}^{\chi_b(3P)}$. As a result, the predicted polarization value λ_θ of the prompt $\Upsilon(3S)$ is insensitive to the input branching ratio of $\chi_b(3P) \to \Upsilon(3S)\gamma$ but sensitive to the observed feeddown fraction $R_{\Upsilon(3S)}^{\chi_b(3P)}$.

IV. SUMMARY

At NLO in NRQCD, we study the $\Upsilon(nS)$ and $\chi_b(nP)$ (n=1,2,3) production at the LHC. We extract the LDMEs of $\Upsilon(nS)$ and $\chi_b(nP)$ production from the LHC large p_T yield data [15, 16, 26, 27], and then with these LDMEs make predictions for the $\Upsilon(nS)$ polarizations. We find

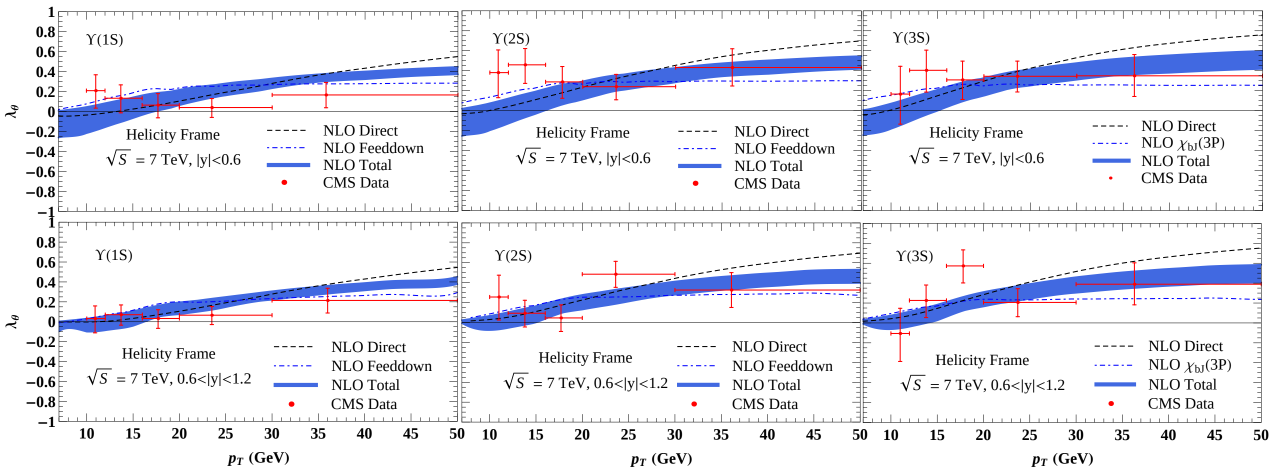

FIG. 4: The polarization parameter λ_θ in the helicity frame for the experimental widows at the LHC. From left to right: $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$. The contributions from direct production are denoted by dashed lines, while those from feeddown by dashed-dotted lines. The total results are denoted by the blue bands. The experimental data are taken from Ref.[9].

that for large p_T (>15 GeV) while the observed $\Upsilon(nS)$ differential p_T cross sections, the fractions of $\Upsilon(mS)$ production from $\chi_b(nP)$ decays, and the differential cross section ratio of $\chi_{b2}(1P)$ to $\chi_{b1}(1P)$) can be rather well described, the predicted $\Upsilon(1S, 2S, 3S)$ polarizations also agree with the recent measurements by CMS [9] within errors. As a result, a simultaneously good description for the large p_T cross sections and polarizations of $\Upsilon(1S, 2S, 3S)$ is achieved at NLO in NRQCD. In particular, the prompt $\Upsilon(3S)$ polarization puzzle can be understood with a large feeddown contribution from $\chi_b(3P)$ states.

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