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# Measurement of the forward-backward asymmetries in the production of $\Xi$ and $\Omega$ baryons in $p \bar{p}$ collisions 

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We measure the forward-backward asymmetries $A_{\mathrm{FB}}$ of charged $\Xi$ and $\Omega$ baryons produced in $p \bar{p}$ collisions recorded by the D0 detector at the Fermilab Tevatron collider at $\sqrt{s}=1.96 \mathrm{TeV}$ as a function of the baryon rapidity $y$. We find that the asymmetries $A_{\text {FB }}$ for charged $\Xi$ and $\Omega$ baryons are consistent with zero within statistical uncertainties.

## I. INTRODUCTION

We present a study of the forward-backward asymmetries $A_{\mathrm{FB}}$ for the production of charged $\Xi$ and $\Omega$ baryons in $p \bar{p}$ collisions at a center of mass energy $\sqrt{s}=1.96 \mathrm{TeV}$, recorded by the D0 detector at the Fermilab Tevatron collider.

We previously performed a study of $A_{\mathrm{FB}}$ for $\Lambda$ and $\bar{\Lambda}$ production [1], where $A_{\mathrm{FB}}$ is defined as the relative excess of $\Lambda(\bar{\Lambda})$ baryons produced with longitudinal momentum $p_{z}$ in the $p(\bar{p})$ direction. These results are in agreement with the observations in a wide range of proton collision experiments that the $\bar{\Lambda} / \Lambda$ production ratio follows a universal function of the "rapidity loss" $y_{p}-y$ between the beam proton and the produced $\bar{\Lambda}$ or $\Lambda$ baryon which does not depend significantly on $\sqrt{s}$ or on the nature of the target $p, \bar{p}, \mathrm{Be}$, or Pb (see Ref. [1] and references therein). These results support the view that a strange quark produced directly in the hard scattering of pointlike partons, or indirectly in the subsequent showering, can coalesce with a diquark remnant of the beam particle to produce a $\Lambda$ baryon with a probability that increases

[^1]as the rapidity difference between the incoming proton and outgoing $\Lambda$ baryon decreases.

If this hypothesis is correct, we also expect $A_{\mathrm{FB}}>0$ for $\Lambda_{c}\left(\bar{\Lambda}_{c}\right)$, and $\Lambda_{b}\left(\bar{\Lambda}_{b}\right)$ production in which a $c$ or $b$ quark can coalesce with a diquark form the proton. For the $B$ mesons and $\Xi$ and $\Omega$ baryons, we expect $A_{F B} \approx 0$ since these particles do not share a diquark with the proton. Previous D0 measurements include $A_{\mathrm{FB}}\left(B^{-}, B^{+}\right)[2]$ and $A_{\mathrm{FB}}\left(\Lambda_{b}, \bar{\Lambda}_{b}\right)[3]$.

In this article, we present measurements of the forward-backward asymmetries of $\Xi^{\mp}$ and $\Omega^{\mp}$ production, where we use the notation $\Xi^{+} \equiv \overline{\Xi^{-}}$and $\Omega^{+} \equiv \overline{\Omega^{-}}$. The $\Xi^{-}$and $\Xi^{+}$baryons are defined as "forward" if their $p_{z}$ points in the $p$ or $\bar{p}$ direction, respectively. The asymmetry $A_{\mathrm{FB}}$ is defined as

$$
\begin{equation*}
A_{\mathrm{FB}} \equiv \frac{\sigma_{\mathrm{F}}\left(\Xi^{-}\right)-\sigma_{\mathrm{B}}\left(\Xi^{-}\right)+\sigma_{\mathrm{F}}\left(\Xi^{+}\right)-\sigma_{\mathrm{B}}\left(\Xi^{+}\right)}{\sigma_{\mathrm{F}}\left(\Xi^{-}\right)+\sigma_{\mathrm{B}}\left(\Xi^{-}\right)+\sigma_{\mathrm{F}}\left(\Xi^{+}\right)+\sigma_{\mathrm{B}}\left(\Xi^{+}\right)} \tag{1}
\end{equation*}
$$

where $\sigma_{\mathrm{F}}$ and $\sigma_{\mathrm{B}}$ are the forward and backward cross sections of $\Xi^{-}$or $\Xi^{+}$production, and similarly for $\Omega^{\mp}$ baryons. The measurements include $\Xi^{\mp}$ and $\Omega^{\mp}$ baryons that are either directly produced or decay products of heavier hadrons. The measurement strategy for the asymmetry $A_{\mathrm{FB}}$ of $\Xi^{\mp}$ and $\Omega^{\mp}$ baryons presented here closely follows the analysis method used to determine $A_{\mathrm{FB}}$ for $\Lambda$ and $\bar{\Lambda}$ baryons in Ref. [1].

## II. DETECTOR AND DATA

The D0 detector is described in detail in Refs. [4-7]. The collision region is surrounded by a central tracking system that comprises a silicon microstrip vertex detector and a central fiber tracker, both located within a 1.9 T superconducting solenoidal magnet [4], surrounded successively by the liquid-argon/uranium calorimeters, a layer of the muon system [5], comprising drift chambers and scintillation trigger counters, the 1.8 T magnetized
iron toroids, and two additional muon detector layers after the toroids.


FIG. 1: Invariant mass distributions of reconstructed $\Xi^{-} \rightarrow$ $\Lambda \pi^{-}$(circles) and $\Xi^{+} \rightarrow \bar{\Lambda} \pi^{+}$(triangles) for $p \bar{p} \rightarrow \mu \Xi^{\mp} X$ data.


FIG. 2: Invariant mass distributions of reconstructed $\Omega^{-} \rightarrow$ $\Lambda K^{-}$(circles) and $\Omega^{+} \rightarrow \bar{\Lambda} K^{+}$(triangles) for $p \bar{p} \rightarrow \mu \Omega^{\mp} X$ data.

The longitudinal momentum $p_{z}$ and the rapidity $y \equiv$ $\ln \left[\left(E+p_{z}\right) /\left(E-p_{z}\right)\right] / 2$ are both measured with respect to the proton beam direction in the $p \bar{p}$ center of mass frame where $E$ is the energy of the baryon. We present results for the full integrated luminosity of $10.4 \mathrm{fb}^{-1}$, collected from 2002 to 2011, using two data sets (i) $p \bar{p} \rightarrow$ $\Xi^{\mp} X$, and (ii) $p \bar{p} \rightarrow \mu \Xi^{\mp} X$. The first data set is unbiased since it is collected using a pre-scaled trigger on beam crossing ("zero bias events") or with a pre-scaled trigger on energy deposited in the forward counters ("minimum bias events"). The second data set is selected with a suite of single muon triggers which implies that most events contain heavy-quark ( $b$ or $c$ ) decays. This data set is defined using the same muon triggers and muon selections as in Ref. $[8,9]$. The muon data provides a sizable data
set that adds additional statistics for the analysis. For $\Omega$ 's there are fewer events, so we only present results for the set $p \bar{p} \rightarrow \mu \Omega^{\mp} X$.

We observe $\Xi$ baryons through their decays $\Xi^{-} \rightarrow \Lambda \pi^{-}$ and $\Xi^{+} \rightarrow \bar{\Lambda} \pi^{+}$, and $\Omega$ baryons through their decays $\Omega^{-} \rightarrow \Lambda K^{-}$and $\Omega^{+} \rightarrow \bar{\Lambda} K^{+}$, with $\Lambda \rightarrow p \pi^{-}$and $\bar{\Lambda} \rightarrow \bar{p} \pi^{+}$in both cases. The $\Lambda$ and $\bar{\Lambda}$ candidates are reconstructed from pairs of oppositely curved tracks with a common vertex $\left(V^{0}\right)$. Each track is required to have a non-zero impact parameter in the transverse plane (IP) with respect to the $p \bar{p}$ interaction vertex with a significance of at least two standard deviations. The proton (pion) mass is assigned to the daughter track with larger (smaller) total momentum since the decay $\Lambda \rightarrow p \pi$ is just above threshold. The invariant mass of the $\left(p, \pi^{-}\right)$or $\left(\bar{p}, \pi^{+}\right)$pair is required to be in the interval $1.105<M(p \pi)<1.125 \mathrm{GeV}[1]$. We require $\Lambda$ and $\bar{\Lambda}$ candidates with $1.5<p_{T}<25 \mathrm{GeV}$ and pseudorapidity $|\eta|<2.2$ [10], and their IP must be non-zero with a significance greater than two standard deviations.

The $\Lambda(\bar{\Lambda})$ candidate is combined with a negatively (positively) charged-particle track with separation in the transverse plane from the primary vertex with significance greater than three standard deviations and a good vertex with the $\Lambda(\bar{\Lambda})$ candidate. This track is assigned the pion mass for $\Xi$ 's or the kaon mass for $\Omega$ 's. The $\Xi^{\mp}$ or $\Omega^{\mp}$ candidates are required to have an IP consistent with zero within three standard deviations. The observed decay lengths in the transverse plane of the $\Lambda$ and $\Xi^{-}$or $\Omega^{-}$(or $\bar{\Lambda}$ and $\Xi^{+}$or $\Omega^{+}$) are required to be greater than 4 mm . The invariant mass for the $\Xi^{\mp}$ candidate is required to be in the interval $1.2<M(\Lambda \pi)<1.5 \mathrm{GeV}$ and $1.55<M(\Lambda K)<1.85 \mathrm{GeV}$ for $\Omega^{\mp}$ candidates. The kinematic selections for the $\Xi^{\mp}$ and $\Omega^{\mp}$ candidates are $p_{T}>2.0 \mathrm{GeV}$ and $|\eta|<2.2$. The pion or kaon track and the two daughter tracks of the $\Lambda$ baryon are required to be different from any track associated to a muon. The invariant mass distributions for the decays $\Xi^{-} \rightarrow \Lambda \pi^{-}$ and $\Xi^{+} \rightarrow \bar{\Lambda} \pi^{+}$are shown in Fig. 1 and for the decays $\Omega^{-} \rightarrow \Lambda K^{-}$and $\Omega^{+} \rightarrow \bar{\Lambda} K^{+}$in Fig. 2.

## III. RAW ASYMMETRIES AND DETECTOR EFFECTS

We obtain the numbers $N_{\mathrm{F}}\left(\Xi^{\mp}\right)$ and $N_{\mathrm{B}}\left(\Xi^{\mp}\right)$ of reconstructed $\Xi^{\mp}$ baryons in the forward and backward categories in each bin of $|y|$ by counting $\Xi^{\mp}$ candidates in the signal region, $1.305<M(\Lambda \pi)<1.335 \mathrm{GeV}$, and subtracting the counts in two sideband regions, defined by $1.2775<M(\Lambda \pi)<1.2925 \mathrm{GeV}$ and $1.3475<$ $M(\Lambda \pi)<1.3625 \mathrm{GeV}$. The signal region for $\Omega^{\mp}$ candidates is $1.6575<M(\Lambda K)<1.6875 \mathrm{GeV}$, and the sideband regions are $1.630<M(\Lambda K)<1.645 \mathrm{GeV}$ and $1.700<M(\Lambda K)<1.715 \mathrm{GeV}$.

The normalization factor $N$ and the three raw asymmetries $A_{\mathrm{FB}}^{\prime}, A_{\mathrm{NS}}^{\prime}$, and $A_{\Xi}^{\prime}$ are defined by

$$
N_{\mathrm{F}}\left(\Xi^{-}\right) \equiv N\left(1+A_{\mathrm{FB}}^{\prime}\right)\left(1-A_{\mathrm{NS}}^{\prime}\right)\left(1+A_{\Xi}^{\prime}\right)
$$

$$
\begin{align*}
& N_{\mathrm{B}}\left(\Xi^{-}\right) \equiv N\left(1-A_{\mathrm{FB}}^{\prime}\right)\left(1+A_{\mathrm{NS}}^{\prime}\right)\left(1+A_{\Xi}^{\prime}\right), \\
& N_{\mathrm{F}}\left(\Xi^{+}\right) \equiv N\left(1+A_{\mathrm{FB}}^{\prime}\right)\left(1+A_{\mathrm{NS}}^{\prime}\right)\left(1-A_{\Xi}^{\prime}\right), \\
& N_{\mathrm{B}}\left(\Xi^{+}\right) \equiv N\left(1-A_{\mathrm{FB}}^{\prime}\right)\left(1-A_{\mathrm{NS}}^{\prime}\right)\left(1-A_{\Xi}^{\prime}\right), \tag{2}
\end{align*}
$$

and similarly for $\Omega$. The raw asymmetries $A_{\mathrm{FB}}^{\prime}, A_{\mathrm{NS}}^{\prime}$, and $A_{\Xi}^{\prime}$ have contributions from the physical asymmetries $A_{\mathrm{FB}}, A_{\mathrm{NS}}$, and $A_{\Xi}$, and from detector effects. The forward-backward asymmetry $A_{\mathrm{FB}}$ measures the relative excess of $\Xi^{-}\left(\Xi^{+}\right)$baryons with $p_{z}$ in the $p(\bar{p})$ direction. The asymmetry $A_{\mathrm{NS}}$ is given by the relative excess of the sum of $\Xi^{-}$and $\Xi^{+}$baryons with $p_{z}$ in the $\bar{p}$ beam direction (north) with respect to the $p$ beam direction (south). The asymmetry $A_{\Xi}$ is the relative excess of negatively charged over positively charged baryons.

The initial $p \bar{p}$ state is invariant with respect to CP conjugation, which changes the sign of $A_{N S}$ and $A_{\Xi}$, while $A_{\mathrm{FB}}$ remains unchanged. A non-zero value of $A_{\mathrm{NS}}$ or $A_{\Xi}$ would indicate CP violation.


FIG. 3: Distributions of $p_{T}, p_{z}$, and $y$ of reconstructed $\Xi^{-}$ (circles) and $\Xi^{+}$candidates (triangles) with $p_{T}>2 \mathrm{GeV}$, for the minimum bias data sample $p \bar{p} \rightarrow \Xi^{\mp} X$.

The asymmetry $A_{\mathrm{NS}}^{\prime}$ is mainly due to differences in the product of the acceptance and efficiency between the northern hemisphere of the $\mathrm{D} \emptyset$ detector with respect to the southern hemisphere. The difference in reconstruction efficiencies of $\Xi^{-}$and $\Xi^{+}$baryons caused by the different inelastic interaction cross-sections of $p$ and $\bar{p}$ with
the detector material creates the additional asymmetry $A_{\Xi}^{\prime}[1]$.


FIG. 4: Asymmetries $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}, A_{\mathrm{NS}}^{\prime}$ and $A_{\Xi}^{\prime}$ of reconstructed $\Xi^{-}$and $\Xi^{+}$candidates with $p_{T}>2 \mathrm{GeV}$, as a function of $|y|$, for the minimum bias data sample $p \bar{p} \rightarrow \Xi^{\mp} X$. The uncertainties are statistical.

The raw asymmetries including terms up to secondorder in the asymmetries are given by

$$
\begin{align*}
& A_{\mathrm{FB}}^{\prime}=A_{\mathrm{NS}}^{\prime} A_{\Xi}^{\prime} \\
& \quad+\frac{N_{\mathrm{F}}\left(\Xi^{-}\right)-N_{\mathrm{B}}\left(\Xi^{-}\right)+N_{\mathrm{F}}\left(\Xi^{+}\right)-N_{\mathrm{B}}\left(\Xi^{+}\right)}{N_{\mathrm{F}}\left(\Xi^{-}\right)+N_{\mathrm{B}}\left(\Xi^{-}\right)+N_{\mathrm{F}}\left(\Xi^{+}\right)+N_{\mathrm{B}}\left(\Xi^{+}\right)},  \tag{3}\\
& A_{\mathrm{NS}}^{\prime}=A_{\mathrm{FB}}^{\prime} A_{\Xi}^{\prime} \\
& \quad+\frac{-N_{\mathrm{F}}\left(\Xi^{-}\right)+N_{\mathrm{B}}\left(\Xi^{-}\right)+N_{\mathrm{F}}\left(\Xi^{+}\right)-N_{\mathrm{B}}\left(\Xi^{+}\right)}{N_{\mathrm{F}}\left(\Xi^{-}\right)+N_{\mathrm{B}}\left(\Xi^{-}\right)+N_{\mathrm{F}}\left(\Xi^{+}\right)+N_{\mathrm{B}}\left(\Xi^{+}\right)},  \tag{4}\\
& A_{\Xi}^{\prime}=A_{\mathrm{FB}}^{\prime} A_{\mathrm{NS}}^{\prime} \\
& \quad+\frac{N_{\mathrm{F}}\left(\Xi^{-}\right)+N_{\mathrm{B}}\left(\Xi^{-}\right)-N_{\mathrm{F}}\left(\Xi^{+}\right)-N_{\mathrm{B}}\left(\Xi^{+}\right)}{N_{\mathrm{F}}\left(\Xi^{-}\right)+N_{\mathrm{B}}\left(\Xi^{-}\right)+N_{\mathrm{F}}\left(\Xi^{+}\right)+N_{\mathrm{B}}\left(\Xi^{+}\right)} . \tag{5}
\end{align*}
$$

The polarities of the solenoid and toroid magnets were reversed about once every two weeks during data-taking to collect approximately the same number of events for each of the four solenoid-toroid polarity combinations. We apply weights to equalize the sums of $\Xi^{-}$and $\Xi^{+}$candidates reconstructed for each of the four polarity combinations. This averaging over magnet polarities cancels contributions from the detector geometry to $A_{\mathrm{FB}}^{\prime}$ and $A_{\Xi}^{\prime}$, but not to $A_{\text {NS }}^{\prime}[1]$.

TABLE I: Forward-backward asymmetry $A_{\mathrm{FB}}$ of $\Xi^{\mp}$ baryons with $p_{T}>2 \mathrm{GeV}$ in minimum bias events, $p \bar{p} \rightarrow \Xi^{\mp} X$, and muon events $p \bar{p} \rightarrow \mu \Xi^{\mp} X$, and $A_{\mathrm{FB}}$ of $\Omega^{-}$and $\Omega^{+}$baryons with $p_{T}>2 \mathrm{GeV}$ in muon events $p \bar{p} \rightarrow \mu \Omega^{\mp} X$. The first uncertainty is statistical, the second is systematic due to the detector asymmetry $A_{\mathrm{NS}}^{\prime} A_{\Xi}^{\prime}$.

| $\|y\|$ | $A_{\mathrm{FB}} \times 100\left(\Xi^{\mp}\right.$, min. bias $)$ | $A_{\mathrm{FB}} \times 100\left(\Xi^{\mp}\right.$, with $\left.\mu\right)$ | $A_{\mathrm{FB}} \times 100\left(\Omega^{\mp}\right.$, with $\left.\mu\right)$ |
| :---: | :---: | :---: | :---: |
| 0.0 to 0.5 | $-2.78 \pm 3.20 \pm 0.34$ | $-0.20 \pm 0.72 \pm 0.01$ | $-3.43 \pm 2.90 \pm 0.13$ |
| 0.5 to 1.0 | $5.23 \pm 2.85 \pm 0.55$ | $-0.13 \pm 0.66 \pm 0.03$ | $3.25 \pm 2.78 \pm 0.10$ |
| 1.0 to 1.5 | $2.61 \pm 3.75 \pm 0.45$ | $1.55 \pm 0.77 \pm 0.05$ | $0.46 \pm 3.52 \pm 0.14$ |
| 1.5 to 2.0 | $5.09 \pm 9.00 \pm 1.64$ | $-1.14 \pm 2.05 \pm 0.27$ | $5.75 \pm 10.86 \pm 5.70$ |



FIG. 5: Distributions of rapidity $y$ for reconstructed $\Xi^{-}$(circles) and $\Xi^{+}$candidates (triangles) in events with a (a) positively or (b) negatively charged muon for $\Xi^{\mp}$ candidates with $p_{T}>2 \mathrm{GeV}$.

The raw asymmetry $A_{\mathrm{FB}}^{\prime}$ has negligible contributions from detector effects after averaging over solenoid and toroid magnet polarities. The raw asymmetries $A_{\mathrm{NS}}^{\prime}$ and $A_{\Xi}^{\prime}$ are dominated by detector effects [1]. The quadratic term $A_{\mathrm{NS}}^{\prime} A_{\Xi}^{\prime}$ in Eq. (3) corrects $A_{\mathrm{FB}}^{\prime}$ for the detector effects $A_{\mathrm{NS}}^{\prime}$ and $A_{\Xi}^{\prime}$ on the particle counts $N_{\mathrm{F}}\left(\Xi^{\mp}\right)$ and $N_{\mathrm{B}}\left(\Xi^{\mp}\right)$. We can therefore set $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}$ where $A_{\mathrm{FB}}$ is defined in Eq. (1).

## IV. MINIMUM BIAS SAMPLE EVENTS $p \bar{p} \rightarrow \Xi^{\mp} X$

The minimum bias sample contains $3.7 \times 10^{3}$ reconstructed $\Xi^{\mp}$ candidates with $p_{T}>2 \mathrm{GeV}$. Distributions of $p_{T}, p_{z}$, and $y$ for the $\Xi^{\mp}$ candidates are shown in Fig. 3 and the corresponding raw asymmetries $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}$, $A_{\mathrm{NS}}^{\prime}$ and $A_{\Xi}^{\prime}$ in Fig. 4. These asymmetries are calculated using Eqs. 3-5, neglecting the quadratic terms since they are small compared to the statistical uncertainties. The correction $A_{\mathrm{NS}}^{\prime} A_{\Xi}^{\prime}$ needed to obtain $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}$ is mea-


FIG. 6: Asymmetries $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}, A_{\mathrm{NS}}^{\prime}$ and $A_{\Xi}^{\prime}$ of reconstructed $\Xi^{-}$and $\Xi^{+}$candidates with $p_{T}>2 \mathrm{GeV}$, as a function of $|y|$, for $p \bar{p} \rightarrow \mu \Xi^{\mp} X$ events. The uncertainties are statistical.
sured to be consistent with zero within statistical uncertainties, see Figs. 4 (b) and (c). Thus, we choose not to apply this correction, but rather take the full measured detector asymmetry $A_{\mathrm{NS}}^{\prime} A_{\Xi}^{\prime}$ as the systematic uncertainty on the measurement of $A_{\mathrm{FB}}$. The results are summarized in Table I.


FIG. 7: Asymmetry $A_{\mathrm{FB}}^{\prime}=A_{\mathrm{FB}}$ as a function of $|y|$ for $p \bar{p} \rightarrow \mu \Xi^{\mp} X$ events with (a) $2.0<p_{T}<4.0 \mathrm{GeV}$, (b) $4.0<p_{T}<6.0 \mathrm{GeV}$, and (c) $p_{T}>6.0 \mathrm{GeV}$. The uncertainties are statistical.

## V. MUON SAMPLE EVENTS $p \bar{p} \rightarrow \mu \Xi^{\mp} X$ AND $p \bar{p} \rightarrow \mu \Omega^{\mp} X$

To study the asymmetries using a larger data set, we consider $p \bar{p} \rightarrow \mu \Xi^{\mp} X$ and $p \bar{p} \rightarrow \mu \Omega^{\mp} X$ events taken from the single muon trigger sample. Charged particles with transverse momentum in the range $1.5<p_{T}<25$ GeV and $|\eta|<2.2$ are considered as muon candidates. Muon candidates are further selected by matching central tracks with a segment reconstructed in the muon system and by applying tight quality requirements aimed at reducing false matching and background from cosmic rays and beam halo. To ensure that the muon candidate traverses the detector, including all three layers of the muon system, we require either $p_{T}>4.2 \mathrm{GeV}$ or $\left|p_{z}\right|>5.4$ $\mathrm{GeV}[8,9]$. The inclusive muon sample contains $2.2 \times 10^{9}$ reconstructed muons and $7.7 \times 10^{4}$ reconstructed $\Xi^{-}$and $\Xi^{+}$candidates with $p_{T}>2 \mathrm{GeV}$, as well as $1.4 \times 10^{4}$ reconstructed $\Omega^{-}$and $\Omega^{+}$candidates.

Rapidity distributions for reconstructed $\Xi^{-}$and $\Xi^{+}$ candidates are shown in Fig. 5. From these distribu-


FIG. 8: Distributions of $p_{T}, p_{z}$, and $y$ of reconstructed $\Omega^{-}$ (circles) and $\Omega^{+}$(triangles) with $p_{T}>2 \mathrm{GeV}$, for the data sample $p \bar{p} \rightarrow \mu \Omega^{\mp} X$.
tions we observe that (i) the detection efficiency for $\Xi^{-}$ baryons is larger than for $\Xi^{+}$baryons as explained above, and (ii) there are more $\Xi^{\mp} \mu^{ \pm}$than $\Xi^{\mp} \mu^{\mp}$ events. An example of a process with a correlated $\Xi^{-} \mu^{+}$pair is the decay $\Xi_{c}^{0} \rightarrow \Xi^{-} \mu^{+} X$. We find that the asymmetry $A_{\mathrm{FB}}^{\prime}$ obtained with events containing a $\mu^{+}$is consistent with the corresponding asymmetry with $\mu^{-}$within statistical uncertainties. We therefore combine the $\mu^{+}$and $\mu^{-}$samples to obtain the asymmetries presented in Figs. 6 and 7.

The $p_{T}, p_{z}$, and $y$ distributions for $p \bar{p} \rightarrow \mu \Omega^{\mp} X$ events are shown in Fig. 8, and the corresponding asymmetry $A_{\mathrm{FB}}$ is presented in Fig. 9. The $\Xi^{\mp}$ and $\Omega^{\mp}$ asymmetries are summarized in Table I.

## VI. CONCLUSIONS

We have measured the forward-backward asymmetries $A_{\mathrm{FB}}$ in $p \bar{p} \rightarrow \Xi^{\mp} X, p \bar{p} \rightarrow \mu \Xi^{\mp} X$, and $p \bar{p} \rightarrow \mu \Omega^{\mp} X$ events using $10.4 \mathrm{fb}^{-1}$ of integrated luminosity recorded with the D 0 detector. We find that $A_{\mathrm{FB}}$ for $\Xi^{\mp}$ and $\Omega^{\mp}$ are consistent with zero within uncertainties.


FIG. 9: Asymmetry $A_{\mathrm{FB}}$ as a function of $|y|$ for events $p \bar{p} \rightarrow$ $\mu \Omega^{\mp} X$ for $p_{T}>2 \mathrm{GeV}$. The uncertainties are statistical.

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