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# Novel inclusive search for the Higgs boson in the four-lepton final state at CDF

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An inclusive search for the standard model Higgs boson using the four-lepton final state in proton-antiproton collisions produced by the Tevatron at  $\sqrt{s} = 1.96$  TeV is conducted. The data are recorded by the CDF II detector and correspond to an integrated luminosity of  $9.7 \text{ fb}^{-1}$ . Three distinct Higgs decay modes, namely  $ZZ$ ,  $WW$ , and  $\tau\tau$ , are simultaneously probed. Nine potential signal events are selected and found to be consistent with the background expectation. We set a 95% credibility limit on the production cross section times the branching ratio and subsequent decay to the four lepton final state for hypothetical Higgs boson masses between  $120 \text{ GeV}/c^2$  and  $300 \text{ GeV}/c^2$ .

The vector gauge bosons mediating the weak force, the  $W$  and  $Z$ , are massive. Within the standard model of

particle physics (SM) their masses arise through spontaneous symmetry breaking (SSB) of the electroweak symmetry [1–3] through the introduction of a scalar field that maintains the gauge invariance of the theory [4–7]. Fundamental fermions may acquire mass through Yukawa couplings with this field. Quantization of the field necessitates the existence of an associated spin-0 particle, known as the Higgs boson, the discovery of which would confirm the SM mechanism for electroweak SSB.

Direct searches at LEP [8], combined with recent search results from the Tevatron [9] and LHC experiments [10, 11], exclude all potential SM Higgs masses outside the ranges  $116.6 - 119.4 \text{ GeV}/c^2$  and  $122.1 - 127 \text{ GeV}/c^2$ . Recently both LHC experiments observed a resonance at a mass of about  $125 \text{ GeV}/c^2$  with properties consistent with those of the SM Higgs boson [12, 13]. Although signal processes leading to four-lepton final states are not the most dominant within these mass ranges the inclusion of additional channels improves overall search sensitivity. Moreover, using the new techniques presented here, the search is simultaneously sensitive to multiple production and decay modes. Also based on this inclusive approach, information regarding Higgs boson couplings to both fermions and bosons can be extracted, which could be useful to probe the SM or beyond SM nature of a hypothetical signal.

Searches for the SM Higgs boson to four leptons

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through the decay  $ZZ^{(*)}$  have been published by ATLAS [14] and CMS [15] using proton-proton collision data at a center of mass energy of 7 TeV. In this Letter a search for the Higgs boson produced inclusively and decaying into final states containing either four electrons ( $4e$ ), four muons ( $4\mu$ ) or two electrons and two muons ( $2e2\mu$ ) is reported. Data from proton-antiproton collisions at center of mass energy 1.96 TeV, collected with the CDF II detector [16] at the Tevatron and corresponding to  $9.7\text{ fb}^{-1}$  of integrated luminosity, are used. Higgs boson decays to  $Z$  boson pairs ( $H \rightarrow ZZ$ ) are the dominant contribution for most of the considered Higgs boson mass range. Owing to the full reconstruction of the final state, the  $ZZ$  channel provides good search sensitivity in spite of the small potential signal yields because the resonance structure of the signal can be exploited to distinguish it from non-resonant backgrounds. Along with the direct Higgs production mechanisms of gluon-fusion ( $ggH$ ) and fusion of vector bosons emitted from the incoming partons (VBF), the search is sensitive to associated Higgs boson production processes ( $VH$ ). In the case of  $ZH$  production, additional potential four-lepton event contributions originate from Higgs boson decays to  $W$  boson pairs ( $ZH \rightarrow ZWW \rightarrow ll\nu\nu$ ) and  $\tau$ -lepton pairs ( $ZH \rightarrow Z\tau\tau \rightarrow ll\nu\nu$ ). The detection of four leptons provides one of the cleanest signatures available at a hadron collider. Because of the small probability associated with mimicking the signature of isolated lepton candidates in the detector, background contributions from ubiquitous multi-jet production processes are negligible in this final state.

Separation of potential signal and background event contributions is obtained from the observed four-lepton invariant mass ( $m_{4\ell}$ ) and missing transverse energy spectra [38]. Sensitivity is significantly enhanced through the inclusion of  $\cancel{E}_T$  information since  $H \rightarrow WW$  and  $H \rightarrow \tau\tau$  decays to leptons result in final state neutrinos, while none are expected in background events. The inclusion of  $\cancel{E}_T$  significantly improves the search sensitivity at Higgs masses below  $150\text{ GeV}/c^2$ , and represents the main advancement in analysis technique over previously published searches in this final state.

Our search is performed using the CDF II detector consisting of a solenoidal spectrometer with a silicon tracker and an open-cell drift chamber (COT) surrounded by calorimeters and muon detectors [16]. The geometry is characterized using the azimuthal angle  $\phi$  and the pseudorapidity  $\eta$ . Transverse energy,  $E_T$ , is defined as  $E \sin \theta$ , where  $E$  is the energy of an electromagnetic or hadronic calorimeter energy cluster. Transverse momentum,  $p_T$ , is the track momentum component transverse to the beam line.

Electron candidates are identified by matching a central or forward track to energy deposited within the calorimeter. Muon candidates are formed from charged particle tracks matched to minimum ionizing energy de-

position in the calorimeter, which may or may not be matched to track segments (*stubs*) in the muon chambers situated behind the calorimeters. Lepton reconstruction algorithms are well validated and described in detail elsewhere [16]. Taus are included in this search only if they decay to electrons or muons.

Candidate leptons are separated into various categories: electrons, muons, and isolated tracks that project to detector regions with insufficient calorimeter coverage for energy measurements. Electron candidates are distinguished by whether they are found in the central or forward calorimeters ( $|\eta| > 1.1$ ) where only silicon tracking information is available. The electron selection relies on track quality, track-calorimeter matching, calorimeter energy, calorimeter profile shape, and isolation information. Most muon candidates rely on direct detection in the muon chambers, which are distinguished by their acceptance in pseudorapidity: central muon detectors ( $|\eta| < 0.6$ ), central muon extension detectors ( $0.6 < |\eta| < 1.0$ ), and the intermediate muon detector ( $1.0 < |\eta| < 1.5$ ). Remaining muon candidates rely on track matches to energy deposits consistent with a minimum ionizing charged particle in the central and forward electromagnetic calorimeters respectively, failing to have an associated stub in the muon sub-detectors. Isolated tracks pointing to uninstrumented regions of the detector and satisfying high-quality requirements are reconstructed as leptons of unidentified flavor. All leptons are required to be isolated by imposing the condition that the sum of the transverse energy of the calorimeter towers in a cone of  $\Delta R \equiv \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = 0.4$  around the lepton is less than 10% of the electron  $E_T$  (muon  $p_T$ ). A similar isolation requirement is applied based on the reconstructed tracks in a  $\Delta R = 0.4$  cone around the lepton candidate track.

The probability that a jet will be misidentified as a lepton is measured using jet-enriched samples and is corrected for the contributions of leptons from  $W$  and  $Z$  boson decays. An average value is obtained from samples collected with different single jet  $E_T$  threshold requirements (20, 50, 70, and 100 GeV) and an uncertainty is assigned based on the spread within the individual measurements. The range of measured misidentification probabilities for the lepton categories, which vary according to  $E_T$  or  $p_T$ , is 0.5% – 3% (central electrons), 2% – 6% (forward electrons), 0.5% – 4% (central muons), 0.5% – 2% (extension muons), 0.5% – 2% (intermediate muons), 0.5% – 6% (calorimeter only muons), and 0.5% – 3% (isolated tracks).

Events for the analysis are collected using online event selection (trigger) requirements corresponding to the presence of a single high- $E_T$  electron or high- $p_T$  muon. The electron trigger requires an electromagnetic energy cluster in the central calorimeter with  $E_T > 18\text{ GeV}$  geometrically matching the direction of a charged particle reconstructed in the COT with  $p_T > 8\text{ GeV}/c$ . Muon

triggers are based on track segments in the muon chambers geometrically matching the direction of a charged particle reconstructed in the COT with  $p_T > 18$  GeV/c. Trigger efficiencies are measured using samples of leptonic  $Z$  decays [16]. To ensure a uniform trigger efficiency over the lepton momentum spectra, an offline selection is applied to the lepton matched to the trigger object, requiring  $E_T > 20$  GeV ( $p_T > 20$  GeV/c) for electrons (muons).

Additional charged electrons (muons) are required to have  $E_T > 10$  GeV ( $p_T > 10$  GeV/c). Exactly four leptons are required, each separated from the others by  $\Delta R \geq 0.1$ . This analysis evolved from a CDF measurement of the  $ZZ$  production cross section in the four lepton final state [17], where constraints on the invariant mass of opposite-sign same-flavor dilepton pairs were imposed in order to explicitly reconstruct  $Z$  bosons. For Higgs boson masses smaller than  $180$  GeV/ $c^2$ , at least one of the  $Z$  bosons is off-shell and such requirements on the mass become inefficient. We therefore require dilepton pair masses to be between 20 and 140 GeV/ $c^2$ . In the final state where all leptons are of the same flavor, opposite-sign pairings are assigned based on the the smallest deviation of the reconstructed masses from the known  $Z$  boson mass. Because the backgrounds are modest, loose constraints on the mass improve the sensitivity of the search. Higgs boson production can also involve hadronic jets originating from the associated vector boson in  $WH$  or  $ZH$  production, the additional quarks in the forward detector from VBF production, or initial state gluon radiation. Therefore, we place no restriction on the number of reconstructed jets contained within each event.

The selected events consist primarily of the background from non-resonant diboson production of  $Z^*/Z$ -boson pairs ( $ZZ$ ). Smaller contributions originate from non-fully leptonic  $ZZ$  decays and  $Z\gamma$  production, which lead to events with two or three real reconstructed leptons and one or two falsely reconstructed leptons associated with the photon or additional hadronic jet activity. The background from top-quark pair production is found to be negligible ( $< 0.01$  events).

The acceptances, efficiencies and kinematic properties of the signal and background processes are determined from simulation. PYTHIA [18] is used to model all Higgs boson production processes and the non-resonant  $ZZ$  background. The modeling of  $Z\gamma$  is based on the simulation framework developed by Bauer and Berger [19]. For all samples CTEQ5L parton distribution functions (PDFs) are used to model the momentum distribution of the initial-state partons [20]. The cross sections for each process are normalized to next-to-next-to-leading order (NNLO) calculations with logarithmic resummation for  $ggH$  [21, 22], NNLO for  $VH$  [23–25], and next-to-leading order (NLO) calculations for VBF [23, 26],  $ZZ$  [27], and  $Z\gamma$  [28].

The response of the CDF II detector is modeled with

a GEANT-based simulation [29]. Efficiency corrections for the simulated CDF II detector response for leptons and photon conversions were determined using independent data samples.

The overall normalization of the estimated background from false leptons is derived using a data-driven approach while the shape of the distribution is derived from MC simulation. The total contribution from false lepton events is estimated using events with two or three leptons and additional jets that may be falsely reconstructed as leptons weighted by the measured jet-to-lepton misidentification probabilities. Since the number of available three-lepton events in data is small, we model kinematic distributions using a weighted sum of those obtained from the simulated  $ZZ$  and  $Z\gamma$  MC samples.

The kinematic distributions of  $ZZ$  events with falsely reconstructed leptons is assumed to be the same as that for correctly reconstructed  $ZZ$  events. We apply the same procedure to the simulated  $Z\gamma$  sample to obtain the invariant mass distribution for events from this process, which are found to be well modeled by a Landau function. We model  $\cancel{E}_T$  in  $Z\gamma$  events using the distribution from simulated  $ZZ$  events with falsely reconstructed leptons, which was found to agree.

In the ranges  $50 < m_{4l} < 600$  GeV/ $c^2$  and  $0 < \cancel{E}_T < 200$  GeV, we estimate background contributions of  $10.59 \pm 1.34$   $ZZ$  events and  $0.39 \pm 0.19$  events with falsely reconstructed leptons. For a SM Higgs boson with a mass of 125 GeV/ $c^2$ , contributions of 0.053 ( $ggH$ ), 0.003 (VBF), 0.006 ( $WH$ ), and 0.089 ( $ZH$ ) events, yielding a total of  $0.15 \pm 0.01$  events, are expected. The indicated uncertainties include statistical and systematic contributions added in quadrature. Systematic uncertainties are described below.

We observe a total of nine events, which is consistent with the rate expected from background sources only. The four-lepton invariant mass and  $\cancel{E}_T$  distributions are shown in Fig. 1(a) and Fig. 1(b) respectively, both overlaid with expected contributions from backgrounds and Higgs boson production for a mass of 125 GeV/ $c^2$ , shown separately for each production process.

A variety of possible systematic effects were considered including those that affect the normalization and the shape of the kinematic distributions. The dominant systematic uncertainties are those on the theory predictions for the cross sections of signal and background processes. Systematic uncertainties associated with the MC affect Higgs signal,  $ZZ$ , and  $Z\gamma$  acceptances determined from the simulated event samples.

Uncertainties originating from lepton selection and trigger efficiency measurements are propagated through the acceptance calculation, leading to uncertainties of 3.6% and 0.5%, on the predicted signal and background event yields, respectively. In addition, all signal and background estimates obtained from simulation have an additional 5.9% uncertainty originating from the mea-

surement of the luminosity [30].

The  $gg \rightarrow H$  cross-section has been computed at NNLO and next-to-next-to-leading log (NNLL) precision varying the renormalization and factorization scales and implementing different correlated combinations of the MSTW2008 error PDFs together with allowed values of the strong coupling constant ( $\alpha_s$ ) [31, 32], yielding systematic uncertainties of 7.0% and 7.7%, respectively. Uncertainties on VBF and associated Higgs boson production, which account for about a quarter of the total Higgs boson events in our sample, are of 5% and 10%, respectively [33]. A 3% uncertainty was assigned on the branching fraction for  $H \rightarrow ZZ$  and  $H \rightarrow WW$ , which are 100% correlated, as well as a 3% uncorrelated uncertainty on  $H \rightarrow \tau\tau$  [33]. The PYTHIA MC simulation for  $ZZ$  production that is used for determining acceptances is based on an expansion at LO; MCFM [34] was used to estimate the potential difference in the acceptance arising from a full NLO simulation, estimated to be  $\pm 2.5\%$  and assigned as a systematic uncertainty. We assign a 10% uncertainty on the  $ZZ$  cross-section based on the difference between LO and NLO [35] predictions. The  $Z\gamma$  component was parametrized using a Landau distribution, where an uncorrelated 50% uncertainty was assigned to the yield prediction in each bin to fully cover the possible effects of mis-modeling. Misidentification probabilities were measured in several jet samples and the maximum spread between these measurements was assigned as a systematic uncertainty on the background estimation. Propagated through to the acceptance, this technique results in a 50% variation in the estimate of the background from fake leptons. The missing transverse energy is scaled up and down by 20% to account for potential MC mis-modeling, and the modified shapes are taken as shape uncertainties on the MC  $\cancel{E}_T$  distributions used as inputs to the limit calculation.

Only three  $\cancel{E}_T$  bins (0–15, 15–45, and 45–200 GeV) are used in the limit setting procedure. The varied spacing was chosen as a compromise between search sensitivity and ensuring sufficient number of simulated events in each bin.

As a cross-check the distribution of the number of jets per event in data was compared to the expectation from the  $ZZ$  MC simulation and found to be consistent.

Upper limits at the 95% credibility level (C.L.) are set on the Higgs boson production cross section,  $\sigma_H$ , as a function of  $m_H$ . A Bayesian technique [36] was employed, where the posterior probability density was constructed from the joint Poisson probability of observing the data in each bin of the  $m_{4l}-\cancel{E}_T$  space, integrating over the uncertainties of the normalization parameters using Gaussian priors. A non-negative constant prior in the signal rate was assumed. The expected limit and associated one and two sigma bands are shown along with the observed limit in Table I and Fig. 2. For a Higgs boson with  $m_H = 125 \text{ GeV}/c^2$  we expect a sensitivity of  $26.5 \times \sigma_{SM}$

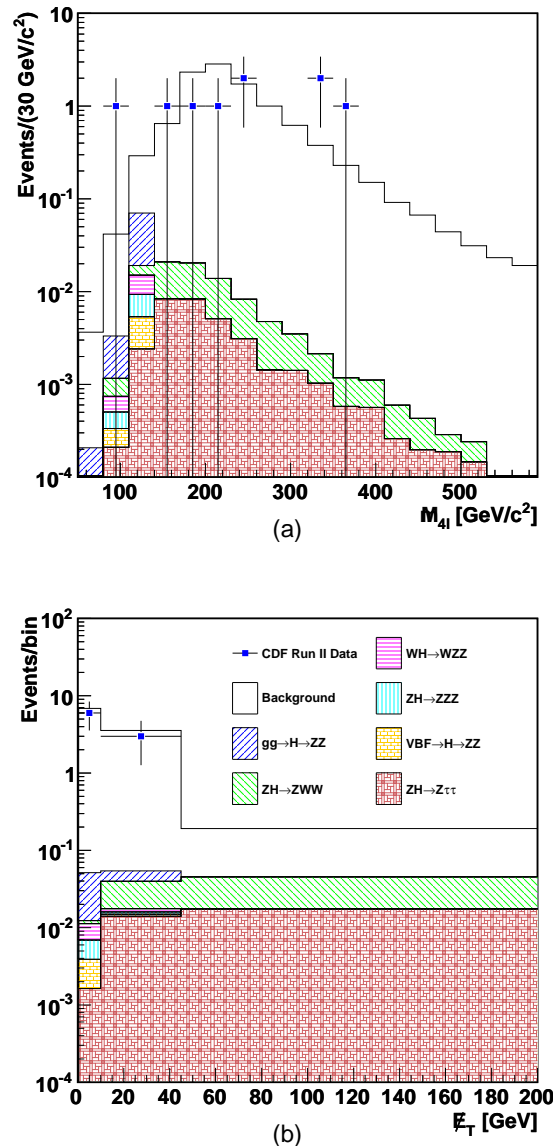


FIG. 1: Distributions of the four-lepton invariant mass (a) and the missing transverse energy (b) in data. The combined estimated contribution from non-resonant  $ZZ$  production and fake leptons, denoted as background, is overlaid. The potential contributions of the different Higgs production processes for  $m_H = 125 \text{ GeV}/c^2$  are stacked and also overlaid.

at 95% C.L. while we observe a limit of  $29.3 \times \sigma_{SM}$  at 95% C.L. . The results obtained in the investigated mass range are consistent with the observed event excess near  $325 \text{ GeV}/c^2$  in the CDF search for high-mass resonances decaying to  $ZZ$  [37]. The analysis reported here is performed using standard CDF tracking algorithms while [37] uses an alternative reconstruction.

In  $9.7 \text{ fb}^{-1}$  of integrated luminosity collected at the Tevatron, no evidence for a Higgs boson signal is found in the mass range explored. Upper limits on the Higgs bo-

TABLE I: Expected and observed upper limits on the Higgs boson production relative to the SM prediction for Higgs particle masses from 120 GeV/c<sup>2</sup> to 300 GeV/c<sup>2</sup>.

H → 4ℓ	120	130	140	150	160	170	180	190	200	210	220	230	240	250	260	270	280	290	300
Expected/σ <sub>SM</sub>	38.0	18.3	11.7	9.4	16.0	25.1	18.5	9.8	10.6	12.9	15.7	16.6	18.9	20.5	21.1	23.2	23.5	28.0	30.5
Observed/σ <sub>SM</sub>	42.4	20.5	12.6	9.5	16.8	28.5	16.3	8.2	7.2	7.9	10.3	20.5	21.1	17.4	17.3	18.2	19.9	24.1	28.6

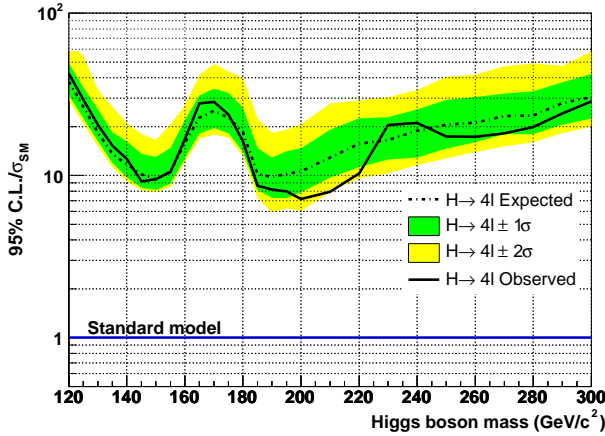


FIG. 2: Expected and observed upper limits on the Higgs boson production relative to the SM prediction as a function of Higgs boson mass.

son production cross section in the inclusive four-lepton final state relative to the SM expectation are obtained. Maximal search sensitivity is obtained not only in the high mass region where both  $Z$  bosons are produced on-shell but also in the lower mass region around 150 GeV/c<sup>2</sup> where additional signal contributions from  $ZH \rightarrow ZWW$  and  $ZH \rightarrow Z\tau\tau$  improve the sensitivity by 15%.

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