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Generalized Hidden \mathcal{Z}_2 Symmetry of Neutrino Mixing

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We explore the consequences of the neutrino mass matrix having a hidden \mathcal{Z}_2 symmetry and one zero eigenvalue. When implemented, these two conditions give relations among the mixing angles. In addition, fitting these relations to the existing oscillation data allows limits to be placed on the parameter of the symmetry.

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1. INTRODUCTION

Neutrino physics can anticipate an era of higher precision measurements with the upcoming generation of neutrino experiments. In the past, measurements have shown that the mixing pattern of lepton sector is quite different from that of quark sector. In the Pontecorvo-Maki-Nakagawa-Sakata (PMNS) parameterization [1] for lepton mixing there are two large mixing angles. The atmospheric mixing angle $\theta_a \equiv \theta_{23}$ is almost maximal while the solar mixing angle $\theta_s \equiv \theta_{12}$ is also large and the reactor mixing angle $\theta_x \equiv \theta_{31}$ nearly vanishes. The recent results are summarized in Table I. We can see that the uncertainties in mixing angles are not particularly small. In most measurements there is roughly a 3° deviation at 1σ confidence level. The mixing matrix which incorporates these angles and diagonalizes

	$\Delta_s (10^{-5} \text{eV}^2)$	$\Delta_a (10^{-3} \text{eV}^2)$	$\sin^2 \theta_s (\theta_s)$	$\sin^2 \theta_a (\theta_a)$	$\sin^2 \theta_x (\theta_x)$
Central Value	7.67	2.39	0.312 (34.0°)	0.466 (43.0°)	0.016 (7.3°)
1σ Range	7.48 – 7.83	2.31 – 2.50	0.294 – 0.331 (32.8 – 35.1°)	0.408 – 0.539 (39.7 – 47.2°)	0.006 – 0.026 (4.4 – 9.3°)

TABLE I: The global 3ν fit [2] for the neutrino mass-squared differences and mixing angles including the available data from solar, atmospheric, reactor (KamLAND and Chooz) and accelerator (K2K and MINOS) experiments [3].

the mass matrix M_ν via $U^T M_\nu U = M_\nu^{\text{diag}}$ is given by

$$U_\nu = \begin{pmatrix} c_s c_x & -s_s c_x & -s_x e^{i\delta_D} \\ s_s c_a - c_s s_a s_x e^{-i\delta_D} & c_s c_a + s_s s_a s_x e^{-i\delta_D} & -s_a c_x \\ s_s s_a + c_s c_a s_x e^{-i\delta_D} & c_s s_a - s_s c_a s_x e^{-i\delta_D} & c_a c_x \end{pmatrix} \quad (1)$$

with $(s_\alpha, c_\alpha) \equiv (\sin \theta_\alpha, \cos \theta_\alpha)$ for $\alpha = s, a, x$. δ_D is the Dirac phase and we have neglected Majorana phases. From Table I we see that a good first approximation is to take $\theta_x = 0$, $\theta_a = 45^\circ$ which gives

$$U_\nu(\theta_s) = \begin{pmatrix} \cos \theta_s & -\sin \theta_s & 0 \\ \sqrt{\frac{1}{2}} \sin \theta_s & \sqrt{\frac{1}{2}} \cos \theta_s & -\sqrt{\frac{1}{2}} \\ \sqrt{\frac{1}{2}} \sin \theta_s & \sqrt{\frac{1}{2}} \cos \theta_s & \sqrt{\frac{1}{2}} \end{pmatrix} \equiv (v_1 \ v_2 \ v_3). \quad (2)$$

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Using this as a starting point, we wish to investigate whether there is an underlying symmetry G of the neutrino mass matrix. This matrix must satisfy $[G, M_\nu] = 0$, or $G^T M_\nu G = M_\nu$. Given a G , the transformation GU_ν also diagonalizes M_ν . But U_ν is unique except for phases. This can be seen by supposing that d_ν is a unitary matrix such that $U_\nu d_\nu$ also diagonalizes M_ν . For this to be true, d_ν must satisfy $d_\nu^T M_\nu^{\text{diag}} d_\nu = M_\nu^{\text{diag}}$ and this implies that $d_\nu^2 = 1$. Since GU_ν diagonalizes M_ν , it must have the form $GU_\nu = U_\nu d_\nu$, where d_ν has 1 or -1 diagonal elements. Thus

$$U_\nu^\dagger G U_\nu = d_\nu \equiv \begin{pmatrix} d_1 & & \\ & d_2 & \\ & & d_3 \end{pmatrix} \Leftrightarrow G = U_\nu d_\nu U_\nu^\dagger. \quad (3)$$

There are only eight possible combinations for the elements of d_ν . Two of these are the unit matrix and its negative, both of which define G as a multiple of the identity. Of the remaining six, three have two entries of $+1$ and one of -1 , while the other three have two entries of -1 and one of $+1$. These diagonal matrices differ by an overall minus sign, so only one of the two types is independent. If we choose the three with one entry of 1 and two entries of -1 ($\det(G) = 1$), then it is easy to see that multiplying any pair of these diagonal matrices will result in the remaining matrix. Hence, in reality, only two of these matrices are independent and both represent a \mathcal{Z}_2 symmetry. Since the independent G 's commute, the horizontal symmetry of lepton mixing is $\mathcal{Z}_2 \times \mathcal{Z}_2$ if neutrinos are Majorana fermions [4–7].

A representation of G can be obtained using

$$G = d_1 v_1 v_1^\dagger + d_2 v_2 v_2^\dagger + d_3 v_3 v_3^\dagger. \quad (4)$$

Since the eigenvalue 1 can occur in three places, there are three symmetry matrices G

$$G_1 = \begin{pmatrix} c_s^2 - s_s^2 & \sqrt{2}s_s c_s & \sqrt{2}s_s c_s \\ \sqrt{2}s_s c_s & -c_s^2 & s_s^2 \\ \sqrt{2}s_s c_s & s_s^2 & -c_s^2 \end{pmatrix} \quad G_2 = \begin{pmatrix} -(c_s^2 - s_s^2) & -\sqrt{2}s_s c_s & -\sqrt{2}s_s c_s \\ -\sqrt{2}s_s c_s & -s_s^2 & c_s^2 \\ -\sqrt{2}s_s c_s & c_s^2 & -s_s^2 \end{pmatrix},$$

and

$$G_3 = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & -1 & 0 \end{pmatrix},$$

where the subscript i on G_i denotes the component of $d_\nu^{(i)}$ that is $+1$. G_3 gives $\mu - \tau$ symmetry [8] while G_1 is symmetric and commutes with G_3 . For simplicity we can parameterize the solar mixing angle θ_s as

$$\cos \theta_s \equiv \frac{-k}{\sqrt{k^2 + 2}}, \quad \sin \theta_s \equiv \frac{\sqrt{2}}{\sqrt{k^2 + 2}}. \quad (5)$$

Then the mixing matrix (2) takes the form [7]

$$U_\nu(k) = \begin{pmatrix} \frac{-k}{\sqrt{2+k^2}} & \frac{-\sqrt{2}}{\sqrt{2+k^2}} & 0 \\ \frac{1}{\sqrt{2+k^2}} & \frac{-k}{\sqrt{2(2+k^2)}} & -\frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2+k^2}} & \frac{-k}{\sqrt{2(2+k^2)}} & \frac{1}{\sqrt{2}} \end{pmatrix} \equiv U_k. \quad (6)$$

Consequently, the symmetry transformation matrix $G_1(\theta_s)$ can be reexpressed in terms of k

$$G_1(k) = \frac{1}{2+k^2} \begin{pmatrix} 2-k^2 & 2k & 2k \\ 2k & k^2 & -2 \\ 2k & -2 & k^2 \end{pmatrix}. \quad (7)$$

Although we can “derive” a generalized form of G_1 symmetry transformation matrix (7) given the mixing matrix (6), this relationship cannot be reversed. The mixing matrix U_ν can not be uniquely determined solely by G_1 due to the fact that G_1 has degenerate eigenvalues.

Invariance under G_3 requires $\theta_x = 0^\circ$, $\theta_a = 45^\circ$, but invariance under G_1 does not, so below we assume that the neutrino mass matrix is invariant under G_1 , not only in the approximation $\theta_x = 0^\circ$, $\theta_a = 45^\circ$, but for general values of all the mixing angles. In the next section we use this assumption in the form of Eq.(7) with general values of k to derive relations among the mixing angles. In Sec. 3 we compare our results with the experimental values and in Sec. 4 we summarize.

2. INVARIANCE UNDER THE Z_2 SYMMETRY G_1

In this section we show explicitly the consequences of generalized G_1 symmetry. Only two mass square differences have been measured and the neutrino's mass scale has not been determined by experiments. It is possible that one of the mass eigenvalues vanishes. This is also theoretically motivated by minimal seesaw model [9]. We will explore the joint consequences of one vanishing mass eigenvalue and G_1 invariance. For simplicity we will postpone discussion of CP phases to a later article [10].

2.1. Constraints on Mass Matrix Elements

If the neutrinos are Majorana fermions, their mass matrix must be symmetric. We will consider the case that there are three generations of light neutrinos. Then, the most general form of the neutrino mass matrix can be parameterized as

$$M_\nu = \begin{pmatrix} A & B_1 & B_2 \\ B_1 & C_1 & D \\ B_2 & D & C_2 \end{pmatrix}, \quad (8)$$

which has six independent matrix elements. We assume M_ν is invariant under the G_1 symmetry transformation,

$$G_1^T M_\nu G_1 = M_\nu. \quad (9)$$

With the help of (7) and (8), Eq.(9) gives two conditions on the neutrino mass matrix elements of (8) [11],

$$\frac{B_1 + B_2}{C_1 + C_2 + 2D - 2A} = \frac{k}{k^2 - 2}, \quad (10a)$$

$$\frac{B_1 - B_2}{C_1 - C_2} = \frac{1}{k}. \quad (10b)$$

2.2. Eigenvalues and eigenstates

If there is a vanishing mass eigenvalue $m_i = 0$ the corresponding mass eigenstate, which can be denoted as $v \equiv (\alpha, \beta, \gamma)^T$, must satisfy

$$\begin{pmatrix} A & B_1 & B_2 \\ B_1 & C_1 & D \\ B_2 & D & C_2 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \\ \gamma \end{pmatrix} = 0. \quad (11)$$

If we assume $\alpha \neq 0$ then we get three equations

$$A = -\rho B_1 - \sigma B_2, \quad (12a)$$

$$B_1 = -\rho C_1 - \sigma D, \quad (12b)$$

$$B_2 = -\rho D - \sigma C_2, \quad (12c)$$

where $\rho \equiv \beta/\alpha$, $\sigma \equiv \gamma/\alpha$. Thus we have two sets of conditions, (10) from G_1 invariance, and (12) from the vanishing mass eigenvalue.

From the relations (12) we can express the matrix element A in terms of C_1 , C_2 and D

$$A = \rho^2 C_1 + \sigma^2 C_2 + 2\rho\sigma D. \quad (13)$$

Now let us use these in the \mathcal{Z}_2 relations. Eq. (10a) and Eq. (10b) give

$$(\sigma k + 1)C_2 - (\rho k + 1)C_1 + k(\rho - \sigma)D = 0, \quad (14a)$$

$$(\rho k + 1)(2\rho - k)C_1 + (\sigma k + 1)(2\sigma - k)C_2 + [(2 - k^2)(\rho + \sigma) - 2k(1 - 2\rho\sigma)]D = 0. \quad (14b)$$

The above two relations can be reexpressed in terms of only two matrix elements, D and C_2 or C_1 respectively

$$(\rho + \sigma - k)[(\sigma k + 1)C_2 + (\rho k + 1)D] = 0, \quad (15a)$$

$$(\rho + \sigma - k)[(\rho k + 1)C_1 + (\sigma k + 1)D] = 0. \quad (15b)$$

The mass eigenvalues that are nonzero are given by

$$m_{\pm} = \frac{1}{2} \left[A + C_1 + C_2 \pm \sqrt{(A + C_1 + C_2)^2 + 4(\rho^2 + \sigma^2 + 1)(D^2 - C_1 C_2)} \right], \quad (16)$$

where we have used (12b) and (12c) and (13). From (15a) and (15b) it is obvious that one possible solution to the equations for C_1, C_2, D is

$$C_1 = -\frac{\sigma k + 1}{\rho k + 1}D, \quad (17a)$$

$$C_2 = -\frac{\rho k + 1}{\sigma k + 1}D. \quad (17b)$$

This makes $D^2 = C_1 C_2$ and consequently m_- given above would also be zero. Since the experimental data shows that two mass square differences between the three neutrino mass eigenvalues are nonzero, we need at least two masses nonzero in order to have two oscillation lengths.

A second solution of (15a), (15b) is $\rho = \sigma = -1/k$ but then the three relations in (12) simply reproduce the conditions (10). So the conclusion is that we must have

$$\rho = k - \sigma \quad (18)$$

and the conditions (14a) or (14b) reduce to an equation for σ

$$\sigma = \frac{(1 + k^2)C_1 - C_2 - k^2 D}{k(C_1 + C_2 - 2D)}. \quad (19)$$

This relation represents the constraint from G_1 invariance which was originally expressed as (10) where there were two independent relations. Using (18), these two relations are satisfied simultaneously and reduce to a single constraint (19).

The condition (18) can also be substituted into (12a), (12b), and (12c) to give

$$\sigma = \frac{A + kB_1}{B_1 - B_2}, \quad (20a)$$

$$\sigma = \frac{B_1 + kC_1}{C_1 - D}, \quad (20b)$$

$$\sigma = \frac{B_2 + kD}{D - C_2}, \quad (20c)$$

respectively. These three relations are a manifestation of vanishing mass eigenvalue. We can set these equations for σ equal to get relations among the matrix elements A, \dots, D in terms of the parameter k . Not all of these equations are independent but two different relations are possible:

$$(A + kB_1)(C_1 - D) - (B_1 + kC_1)(B_1 - B_2) = 0, \quad (21a)$$

$$(A + kB_1)(D - C_2) - (B_2 + kD)(B_1 - B_2) = 0. \quad (21b)$$

In the next subsection we will write A, \dots, D in terms of the mixing angles and thereby get two relations among the mixing angles, again involving k .

2.3. Reconstruction of Neutrino Mass Matrix

Using U_ν from Eq. (1) in $M_\nu = U^* M_\nu^{\text{diag}} U^\dagger$ and comparing with (8) we get [12]

$$A = c_x^2 c_s^2 m_1 + c_x^2 s_s^2 m_2 + s_x^2 m_3 \quad (22a)$$

$$B_1 = c_x [s_s c_s c_a - s_x s_a c_s^2] m_1 - c_x [s_s c_s c_a + s_x s_a s_s^2] m_2 + c_x s_x s_a m_3 \quad (22b)$$

$$B_2 = c_x [s_s c_s s_a + s_x c_a c_s^2] m_1 - c_x [s_s c_s s_a - s_x c_a s_s^2] m_2 - s_x c_x c_a m_3 \quad (22c)$$

$$C_1 = (s_s c_a - s_x c_s s_a)^2 m_1 + (c_s c_a + s_x s_s s_a)^2 m_2 + c_x^2 s_a^2 m_3 \quad (22d)$$

$$C_2 = (s_s s_a + s_x c_s c_a)^2 m_1 + (c_s s_a - s_x s_s c_a)^2 m_2 + c_x^2 c_a^2 m_3 \quad (22e)$$

$$D = (s_s s_a + s_x c_s c_a)(s_s c_a - s_x c_s s_a) m_1 + (c_s s_a - s_x s_s c_a)(c_s c_a + s_x s_s s_a) m_2 - c_x^2 s_a c_a m_3 \quad (22f)$$

where, as mentioned above, we have deferred consideration of CP violation to a later article. The mass eigenvalues can be further parameterized in terms of experimentally measured mass square differences: $m_1 = m_0$, $m_2 = m_0 \sqrt{1+r}$ and $m_3 = 0$ for inverted mass hierarchy and $m_1 = 0$, $m_2 = m_0 \sqrt{r}$ and $m_3 = m_0$ for normal mass hierarchy where $m_0 \equiv \sqrt{\Delta_a}$ and $r \equiv \Delta_s/\Delta_a$ which is positive.

2.4. Correlations between Mixing Angles

To get relations between mixing angles we can substitute (22) into (21) which gives

$$-c_a c_x [c_x (c_a - s_a) + k s_x] m_1 m_2 = 0, \quad (23a)$$

$$s_a c_x [c_x (c_a - s_a) + k s_x] m_1 m_2 = 0. \quad (23b)$$

where we have assumed the mass hierarchy is inverted with vanishing m_3 and nonzero m_1 and m_2 , while c_a , s_a and c_x are also nonzero. The only possible solution is

$$k = c_x \frac{s_a - c_a}{s_x} \approx \frac{\sqrt{2} \delta_a}{\delta_x}, \quad (24)$$

where the last factor comes from expanding the two mixing angles θ_a and θ_x around the approximations 45° and 0° ,

$$\theta_a \equiv \frac{\pi}{4} + \delta_a, \quad \theta_x \equiv \delta_x. \quad (25)$$

With this solution for k and the reconstructed mass matrix elements (22) substituted back into (20a) we find

$$\frac{1}{\sigma} = \frac{c_a - s_a}{k c_a} = -\frac{s_x}{c_x c_a} \approx -\sqrt{2} \delta_x. \quad (26)$$

Since δ_x is quite small, σ should be very large according to (26).

We still have the condition from \mathcal{Z}_2 , Eq. (19). Together with (26) and (22) as well as (24) it gives

$$\tan 2\theta_s = \frac{2(c_a^2 - s_a^2) s_x}{c_x^2 - (2 + 2s_x^2) c_a s_a} = \frac{2 \left(\frac{c_a + s_a}{c_a - s_a} s_x \right)}{1 - \left(\frac{c_a + s_a}{c_a - s_a} s_x \right)^2} = \frac{2 \left(-\frac{c_a - s_a}{c_a + s_a} \frac{1}{s_x} \right)}{1 - \left(-\frac{c_a - s_a}{c_a + s_a} \frac{1}{s_x} \right)^2}. \quad (27)$$

There are two possible solutions

$$\tan \theta_s = -\frac{c_a - s_a}{c_a + s_a} \frac{1}{s_x} = \frac{k}{c_x (s_a + c_a)} \approx \frac{\delta_a}{\delta_x}, \quad (28)$$

or

$$\tan \theta_s = \frac{c_a + s_a}{c_a - s_a} s_x = -\frac{c_x (c_a + s_a)}{k} \approx -\frac{\delta_x}{\delta_a}. \quad (29)$$

These relations between mixing angles can be used to predict the not well measured θ_x in terms of the solar and atmospheric mixing angles. For example (29) gives

$$s_x = \frac{c_a - s_a s_s}{s_a + c_a c_s} \quad \Rightarrow \quad \delta_x \approx -\tan \theta_s \delta_a \quad (30)$$

Since θ_x is the focus of next generation of neutrino experiments, we use (30) to estimate its value. The scatter plot is shown in Fig. 1. A scatter plot based on (28) would look similar with a steeper slope for the points.

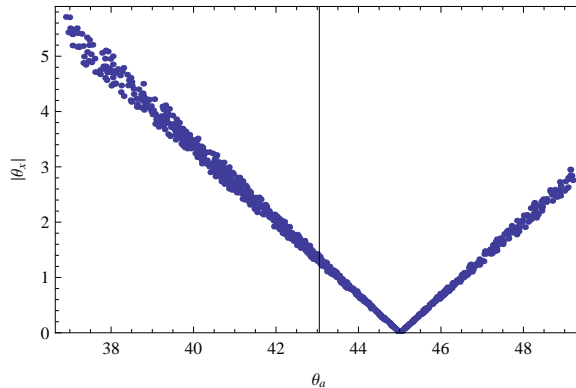


FIG. 1: Prediction of θ_x in terms of θ_a and θ_s at the 90% C.L. The vertical solid line denotes the experimentally measured central value of the atmospheric mixing angle θ_a .

Another way of expressing the results is to write all of the mixing angles in terms of the parameters σ and k . Using $z = 1/\sigma$, this gives

$$\sin^2 \theta_a = \frac{(1 - kz)^2}{k^2 z^2 - 2kz + 2} \quad (31a)$$

$$\sin^2 \theta_x = \frac{z^2}{(k^2 + 1)z^2 - 2kz + 2} \quad (31b)$$

and either

$$\sin^2 \theta_s = \frac{(2 - kz)^2}{k^4 z^2 - 2k^3 z + 2k^2(z^2 + 1) - 4kz + 4} \quad (31c)$$

or

$$\sin^2 \theta_s = k^2 \frac{(k^2 + 1)z^2 - 2kz + 2}{k^4 z^2 - 2k^3 z + 2k^2(z^2 + 1) - 4kz + 4}, \quad (31d)$$

Note that these equations are all unchanged under $k, z \rightarrow -k, -z$ so only the absolute value of k can be determined.

3. FIT TO EXISTING DATA

The solutions (28) or (29), which give (31c) or (31d), are identical in the following sense - oscillation experiments measure $\sin^2 2\theta$ and thus can't distinguish between θ and $\pi/2 - \theta$. Further, $\tan(\pi/2 - \theta) = 1/\tan \theta$, so a fit with (28) and θ_s assumed greater than $\pi/4$ is identical to one with (29) and θ_s assumed less than $\pi/4$. Having noted this we will proceed to fit both (31c) and (31d) with $\theta_s < \pi/4$.

Using Eqs. (31a), (31b), and (31c), the fit to the data from Ref. [2] gives $\chi_{\min}^2 = 2.10$, $|k|_{\min} = 2.09$ and $z = 0.066$. At the minimum values of $|k|$ and z , $\sin^2(\theta_a) = 0.426$ ($\theta_a = 40.7^\circ$), $\sin^2(\theta_x) = 0.0025$ ($\theta_x = 2.87^\circ$) and $\sin^2(\theta_s) = 0.313$ ($\theta_s = 34.0^\circ$). The 68.3% and 90% confidence contours are shown in Fig. 2.

The distributions of this set of mixing angles are obtained from the likelihood distribution

$$Ae^{-(\chi^2(k,z) - \chi_{\min}^2)/2}, \quad (32)$$

where A is a normalization constant, using

$$\frac{dP}{d\sin^2(\theta)} = \int dk \int dz \delta(\sin^2(\theta) - f(k, z)) Ae^{-(\chi^2(k,z) - \chi_{\min}^2)/2}, \quad (33)$$

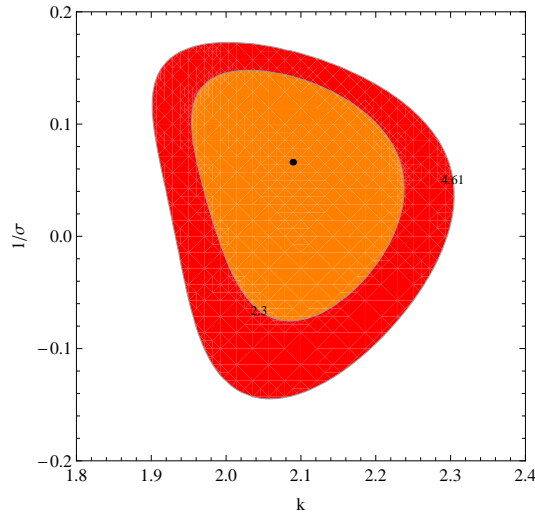


FIG. 2: The 68.3% and 90.0% confidence contours for the fit using Eqs. (31a), (31b), and (31c) are shown in red and orange, respectively. The (black) dot indicates the χ^2 minimum.

where $f(k, z)$ is one of the functions on the righthand side of Eqs. (31). The results are shown in Figs. (3).

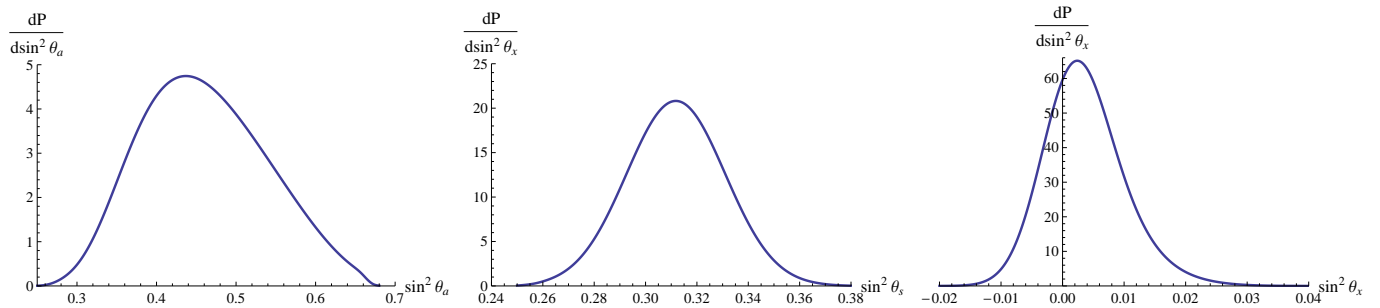


FIG. 3: The distributions of the $\sin^2 \theta_i$ obtained using Eqs. (31a), (31b), and (31c) are shown.

As would be expected, none of the distributions is exactly Gaussian. The largest contribution to the minimum χ^2 is associated with $\sin^2(\theta_a)$ and the influence of terms beyond the quadratic expansion of $\chi^2(k, z)$ can be seen in the shape of this distribution.

If we use Eqs. (31a), (31b) and (31d), the fit to the data has two local minima. The lowest of these gives $\chi^2_{\min} = 0.506$, $|k|_{\min} = 0.942$ and $z_{\min} = 0.152$. At the minimum values of $|k|$ and z , $\sin^2(\theta_a) = 0.423$ (40.5°), $\sin^2(\theta_x) = 0.013$ (6.55°) and $\sin^2(\theta_s) = 0.311$ (33.9°). At the other minimum, where $\chi^2 = 2.73$, $\sin^2 \theta_s$ and $\sin^2 \theta_x$ are slightly different, but $\sin^2 \theta_a = 0.567$. This is reflected in the individual mixing angle distributions. The confidence contours for this case are shown in Fig. 4.

The distributions of this set of mixing angles are shown in Figs. (5). Here, too, the largest contribution to the minimum χ^2 is associated with $\sin^2(\theta_a)$ and the effect of the second local minimum this is reflected in the distortion on the high side of the probability distribution.

Alternately we can fit for k using the values of $\sin^2 \theta_a$ and $\sin^2 \theta_s$ from Ref.[2] but replace $\sin^2 \theta_x$ with the value for $\sin^2(2\theta_a) \sin^2(2\theta_x)$ published by the MINOS collaboration [13]. They report $\sin^2(2\theta_a) \sin^2(2\theta_x) \simeq 0.18 \pm 0.13$ for inverted hierarchy and, for normal hierarchy, $\simeq 0.11 \pm 0.09$. From the inverted hierarchy result we get $|k| = 2.10 \pm 0.10$ with a χ^2 of 1.86 or $|k| = 0.94 \pm 0.15$ with a χ^2 of 1.35.

This was all for inverted hierarchy. Normal hierarchy, with $m1$ equal to zero, gives, after a lot of work, exactly the results of inverted hierarchy, (24), (28), (29). The parameter σ is a different function than (26),

$$\sigma = \frac{s_a(1+k^2) - c_a}{k(s_a + c_a)}, \quad (34)$$

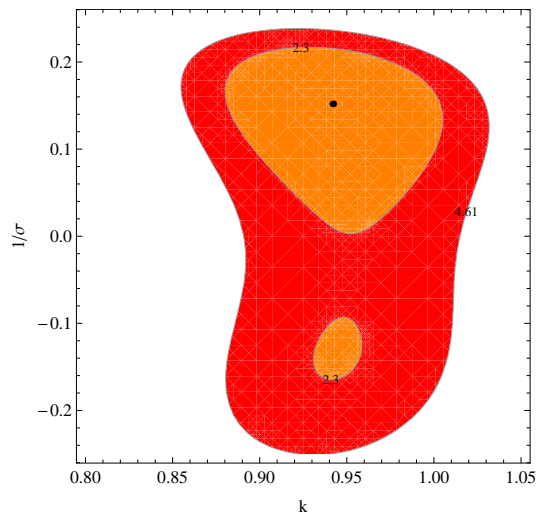


FIG. 4: The 68.3% and 90.0% confidence contours for the fit using Eqs. (31a), (31b) and (31d) are shown in red and orange, respectively. The (black) dot indicates the χ^2 minimum.

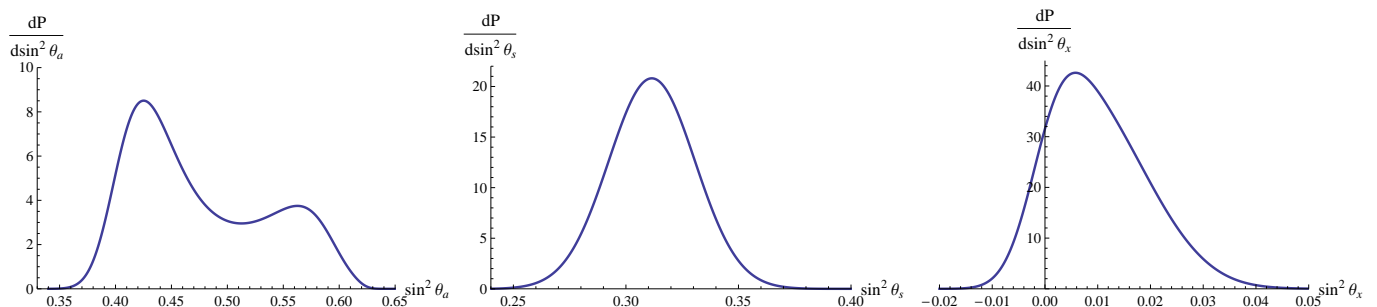


FIG. 5: The distributions of the $\sin^2 \theta_i$ obtained using Eqs. (31a), (31b) and (31d) are shown.

but this just amounts to a reparameterization of (31) with no physical consequence. Using the MINOS number for normal hierarchy we find the same values, including the errors, for $|k|$ as for the inverted hierarchy MINOS number. The χ^2 values are smaller at 1.42 or 0.80.

With either MINOS value and for either value of $|k|$ the fitted value of $\sin^2 \theta_s$ is stable at 0.312, the fitted value of $\sin^2 \theta_a$ varies only slightly from 0.46 for the larger $|k|$ to 0.42 for the smaller value, but $\sin^2 \theta_x$ is less than 0.001 for the larger $|k|$ but equal to 0.015 for the smaller.

4. SUMMARY

A hidden \mathcal{Z}_2 symmetry, as given by Eq. (9), results in only two possible sets of conditions on the neutrino mixing angles. Assuming $\theta_s < \pi/4$ then either

$$s_x = \frac{c_x}{k}(s_a - c_a), \quad (35a)$$

$$\tan \theta_s = -\frac{c_x}{k}(s_a + c_a), \quad (35b)$$

with confidence contours shown in Fig. 2, or

$$s_x = \frac{c_x}{k}(s_a - c_a), \quad (36a)$$

$$\tan \theta_s = \frac{k}{c_x(s_a + c_a)}, \quad (36b)$$

with the confidence contours shown in Fig. 4.

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