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Fast collective neutrino oscillations inside the neutrino sphere in core-collapse supernovae

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Neutrinos are believed to have a key role in the explosion mechanism of core-collapse supernovae as they carry most of the energy released by the gravitational collapse of a massive star. If their flavor is converted fast inside the neutrino sphere, the supernova explosion may be influenced. This paper is reporting the results of the extended work of our previous paper. We perform a thorough survey of the ELN crossing in one of our self-consistent, realistic Boltzmann simulations in two spatial dimensions under axisymmetry for the existence of the crossings between ν_e and $\bar{\nu}_e$ angular distributions, or the electron lepton number (ELN) crossing. We report for the first time the positive detections deep inside the core of the massive star in the vicinity of neutrino sphere at $r \approx 16 - 21$ km. We find that low values of the electron fraction Y_e produced by convective motions together with the appearance of light elements are critically important to give rise to the ELN crossing by enhancing the chemical potential difference between proton and neutron, and hence by mitigating the Fermi-degeneracy of ν_e . Since the region of positive detection are sustained and, in fact, expanding with time, it may have an impact on the explosion of core-collapse supernovae, observational neutrino astronomy and nucleosynthesis of heavy nuclei.

PACS numbers:

I. INTRODUCTION

Neutrinos (ν 's) are fermions and are one of the most abundant particles in the universe [1]. They are massive particles, with their mass eigenstates being not diagonal with their flavor eigenstates [2] and, as a result, they oscillate among their three flavors (ν_e , ν_μ and ν_τ) while propagating in vacuum [3]. When neutrinos propagate through a medium, they gain effective mass due to interactions with matter and give rise to a phenomena called the Mikheyev-Smirnov-Wolfenstein (MSW) effect [3–5]. The self-energy can be also generated by interacting with other neutrinos and if the environment is neutrino-rich, the so-called collective neutrino oscillations occur in which the flavor evolution becomes nonlinear and hence very complicated [6–9].

Core-collapse supernovae (CCSNe), which are the end phase in the evolution of the massive stars with a zero-age-main-sequence (ZAMS) mass of $\gtrsim 8 M_\odot$, are one of the most energetic explosions in the universe and are environments where an enormous amount of neutrinos are produced [10]. Although the exact explosion mechanism of CCSNe is not fully understood, it is well known that

neutrinos are the key players as almost all of the energy released by the gravitational collapse of a massive star is emitted in the form of neutrinos and the kinetic energy of ejected materials in the explosion is only about one percent of the neutrino energy [11].

About a decade ago, it was pointed out by Sawyer that the neutrino flavor conversion may occur near the neutrino sphere and if true, it will have a strong impact on the explosion mechanism of CCSNe. He reported in a series of papers [12–14] a mechanism called "fast oscillation", in which the frequency is proportional to the neutrino potential $\mu \sim \sqrt{2}G_F n_\nu$. It is known that the fast flavor conversions occur when the "electron-lepton-number (ELN) crossing" exists [14–19], that is, the difference between the energy-integrated distribution functions of electron-type neutrinos and their anti-particles changes its sign as a function of propagation direction. It is believed that CCSNe are one of the best astrophysical environments for the fast flavor conversion because neutrinos with different flavors are highly populated there, having different angular distributions.

So far the fast flavor conversions in the realistic settings have been studied only in 1D under spherical sym-

metry and no sign of the ELN crossing has been found [20]. In our previous paper [21], we conducted a pilot study based on a small number of data extracted from our fully self-consistent realistic simulations of CCSNe in two spatial and three momentum dimensions with our Boltzmann-neutrino-radiation-hydrodynamics code [22]. We did not find any crossing between ν_e and $\bar{\nu}_e$ angular distributions at a specific point ($r = 44.8$ km, $\theta = 2.36$ rad) in three different time-steps after bounce (15, 190 and 275 ms). In contrast with our results, a positive detection of the ELN crossings at $r \gtrsim 50 - 70$ km was reported in [23]. It should be mentioned that their results were based on the time-independent neutrino distributions computed for some fixed matter profiles and hence were not fully self-consistent. Very recently Nagakura et al. [24] also found the crossings at similar regions in one of their latest self-consistent simulations with an updated EOS.

The goal of this paper is to conduct a more thorough survey of the ELN crossing in our fully self-consistent, realistic simulations of CCSN in two dimensions under axisymmetry with our Boltzmann-neutrino-radiation-hydrodynamics code that computes neutrino transport together with hydrodynamics. The same data as employed in the previous study [21] are again used. We report for the first time in the following a positive result found at about $r \gtrsim 16 - 21$ km, which is inside the neutrino sphere. If the fast flavor conversion occurs at this small radius indeed, it is likely to have some impact on the supernova explosion.

This paper is structured as follows. In Sec.II, we summarize the equations to give the dispersion relation for the fast collective neutrino oscillations, which will be used in the later analysis, and present briefly the numerical models adopted in this study. Sec.III presents the survey results and finally in Sec.IV we conclude the paper with some discussions.

II. METHOD AND MODELS

A. Dispersion relation

In this paper, we rely on the ELN crossing as a criterion of the fast flavor conversion and do not conduct linear analysis as in the previous paper [21] in most cases. We will still use, however, in some cases the dispersion relation and the growth rate of the fast conversion for detailed analysis. We hence give the procedure to obtain the dispersion relation. Following the previous works [19, 21, 25–30], neglecting ordinary collisions, we begin with the equation of motion for density matrix ρ as

$$(\partial_t + \mathbf{v} \cdot \nabla_{\mathbf{r}})\rho = i[\rho, H]. \quad (1)$$

The Hamiltonian H in Eq. (1) is written as

$$H = H_{\text{vacuum}} + H_{\text{matter}} + H_{\text{collective}}, \quad (2)$$

where each term is given as

$$H_{\text{vacuum}} = \frac{M^2}{2E}, \quad (3)$$

$$H_{\text{matter}} = v^\mu \Lambda_\mu \frac{\sigma_3}{2}, \quad (4)$$

$$H_{\text{collective}} = \sqrt{2}G_F \int \frac{E'^2 dE'}{2\pi^2} d\Upsilon' v^\mu v'_\mu \rho'. \quad (5)$$

In the above expressions, $M^2 = \text{diag}(m_1^2, m_2^2, m_3^2)$ in the mass eigenstates is the mass-squared matrix in vacuum. In H_{matter} , we define the four velocity of neutrino as $v^\mu = (1, \mathbf{v})$; $\Lambda^\mu = \sqrt{2}G_F(n_e - n_{e^+})$ u^μ is the matter potential given with four velocity of matter u^μ ; σ_3 is the third Pauli matrix $\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. In this paper the

metric is assumed to be $\text{diag}(1, -1, -1, -1)$. In $H_{\text{collective}}$, which is responsible for the collective neutrino oscillation, G_F is the Fermi coupling constant, ρ' denotes the density matrix for the neutrinos with energy E' , which moves at the four velocity of $v'^\mu = (1, \mathbf{v}')$; $d\Upsilon' = d\mathbf{v}'/4\pi$ is the infinitesimal solid angle in momentum space normalized by 4π .

As in our previous paper [21], we work in the two-flavor (ν_e and ν_x) approximation as a common practice for simplicity, where ν_x stands for ν_μ and ν_τ collectively. The density matrix in Eq. (1) is expressed as

$$\rho = \frac{f_{\nu_e} + f_{\nu_x}}{2} + \frac{f_{\nu_e} - f_{\nu_x}}{2} \begin{pmatrix} s_{\mathbf{p}} & S_{\mathbf{p}} \\ S_{\mathbf{p}}^* & -s_{\mathbf{p}} \end{pmatrix}, \quad (6)$$

where the off-diagonal element in the matrix, $S_{\mathbf{p}}(t, \mathbf{r})$, is the complex scalar field and represents the flavor coherence between ν_e and $\bar{\nu}_e$ for momentum \mathbf{p} ; the diagonal element, $s_{\mathbf{p}}(t, \mathbf{r})$, is the real field and obeys $|S_{\mathbf{p}}|^2 + s_{\mathbf{p}}^2 = 1$; f_{ν_e} and f_{ν_x} are the neutrino distribution functions for ν_e and ν_x , respectively. When the neutrino is in one of the flavor eigenstates, $s_{\mathbf{p}}$ and $S_{\mathbf{p}}$ are 1 and 0, respectively, and Eq. (1) can be written as

$$\rho = \frac{f_{\nu_e} + f_{\nu_x}}{2} + \frac{f_{\nu_e} - f_{\nu_x}}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (7)$$

We assume ν_x and $\bar{\nu}_x$ have the same distributions in this paper. As we focus on the fast modes here, we ignore H_{vacuum} in the following. Then the equations no longer include E explicitly and we will deal with the energy-integrated quantities, e.g., $S_{\mathbf{v}} = \int E^2 dE S_{\mathbf{p}}$ from this point on. By assuming that $S_{\mathbf{v}} \ll 1$, we linearize Eq. (1) integrated over energy and obtain the following EOM for $S_{\mathbf{v}}$:

$$i(\partial_t + \mathbf{v} \cdot \nabla_{\mathbf{r}})S_{\mathbf{v}} = v^\mu (\Lambda_\mu + \Phi_\mu)S_{\mathbf{v}} - \int d\Upsilon' v^\mu v'_\mu G_{\mathbf{v}'} S_{\mathbf{v}'}, \quad (8)$$

where $G_{\mathbf{v}}$ is the electron-lepton number (ELN) angular distribution defined as

$$G_{\mathbf{v}} = \sqrt{2}G_F \int_0^\infty \frac{dE E^2}{2\pi^2} [f_{\nu_e}(E, \mathbf{v}) - f_{\bar{\nu}_e}(E, \mathbf{v})]. \quad (9)$$

The corresponding ELN current is $\Phi^\mu \equiv \int \frac{d\mathbf{v}}{4\pi} G_{\mathbf{v}} v^\mu$. If the solution of Eq. (8) is assumed to take the form of $S_{\mathbf{v}} = Q_{\mathbf{v}} e^{-i(\Omega t - \mathbf{K} \cdot \mathbf{r})}$, the equation for the amplitude $Q_{\mathbf{v}}$ is given as

$$v^\mu k_\mu Q_{\mathbf{v}} = - \int d\Upsilon' v^\mu v'_\mu G_{\mathbf{v}'} Q_{\mathbf{v}'}, \quad (10)$$

where $k^\mu = K^\mu - \Lambda^\mu - \Phi^\mu$ with $k^\mu = (\omega, \mathbf{k})$ and $K^\mu = (\Omega, \mathbf{K})$.

Equation (10) can be recast into

$$Q_{\mathbf{v}} = v^\mu a_\mu / v^\mu k_\mu \quad (11)$$

with $a_\mu \equiv - \int \frac{d\mathbf{v}'}{4\pi} v'_\mu G_{\mathbf{v}'} Q_{\mathbf{v}'}$. Putting Eq. (11) into Eq. (10), we obtain

$$v_\mu \Pi^{\mu\nu}(\omega, \mathbf{k}) a_\nu = 0, \quad (12)$$

where the polarization tensor $\Pi^{\mu\nu}$ is given as

$$\Pi^{\mu\nu} = \eta^{\mu\nu} + \int \frac{d\mathbf{v}}{4\pi} G_{\mathbf{v}} \frac{v^\mu v^\nu}{\omega - \mathbf{v} \cdot \mathbf{k}} \quad (13)$$

with the Minkowsky metric $\eta^{\mu\nu} = \text{diag}(1, -1, -1, -1)$. Equation (12) has non-trivial solutions if and only if

$$\det \Pi = 0. \quad (14)$$

This equation gives us the dispersion relation between ω and \mathbf{k} . As shown in our previous study [21], it depends on the direction of the wave vector \mathbf{k} in general and the radial direction is not necessarily the most important direction.

B. Models

As in our previous work [21], we conduct our analysis to the results of the realistic two-dimensional (2D) fully self-consistent Boltzmann-neutrino-radiation-hydrodynamics simulations for the non-rotating progenitor model of $11.2 M_\odot$ [31], which were performed on the Japanese K supercomputers [22]. In these simulations, three neutrino species, ν_e , $\bar{\nu}_e$ and ν_x are considered and their distributions are computed on spherical coordinates (r, θ) under spatial axisymmetry; we also employ spherical coordinates in momentum space $(E, \theta_\nu, \phi_\nu)$, in which the two angles are measured from the local radial direction. The computational domain covers $0 \leq r \leq 5000$ km, $0 \leq \theta \leq \pi$, $0 \leq E \leq 300$ MeV, $0 \leq \theta_\nu \leq \pi$ and $0 \leq \phi_\nu \leq 2\pi$ with $384(r)$, $128(\theta)$, $20(E)$, $10(\theta_\nu)$ and $6(\phi_\nu)$ mesh cells. The results of the Furusawa-Shen equation of state (FSEOS), which is based on the relativistic mean field theory for nuclear matter [32, 33], are adopted also in this paper, simply for continuity from the previous pilot study [21].

In this work, unlike in our previous paper [21], which we analyzed a single point ($r = 44.8$ km, $\theta = 2.36$ rad)

alone, we make a survey for all the grid points inside the shock wave, searching for the crossings in the angular distributions between the electron-type neutrinos ν_e and anti-electron neutrinos $\bar{\nu}_e$, i.e., the change of sign in the ELN defined as

$$ELN(\theta_\nu, \phi_\nu) = \int \frac{E^2 dE}{2\pi^2} (f_{\bar{\nu}_e}(E, \theta_\nu, \phi_\nu) - f_{\nu_e}(E, \theta_\nu, \phi_\nu)), \quad (15)$$

where $f_{\bar{\nu}_e}$ and f_{ν_e} are the neutrino distribution functions of $\bar{\nu}_e$ and ν_e , respectively. We also look into the dispersion relation at some representative points.

Figure 1 is the contour plots for density ρ , electron fraction Y_e and temperature T , respectively. White arrows represent the matter velocity, in the central portion ($r \lesssim 30$ km) of the core, on which we will focus in this paper. These are the data from the same snapshots at $t_{pb} = 190$ and 275 ms as used in our previous paper. Although we have explored various post-bounce times this time, we will mainly adopt the results at these times for reasons that will become clear.

III. RESULTS

We report here the positive detection of the crossing in the region of $r \sim 16 - 18$ km at $t_{pb} = 190$ ms. These crossings are sustained for later times to $t_{pb} = 275$ ms, at which time they occur in even wider region of $r \sim 16 - 21$ km. It is important that regions are located in the vicinity of the neutrino sphere and, as we mentioned already, the fast flavor conversion there may have an influence not only on the terrestrial observations of supernova neutrinos and nucleosyntheses but also on the supernova explosion itself.

In order to do further analysis, we pick up a point at the radius $r = 16.5$ km and the spatial zenith $\theta = 2.1$ rad in the crossing region for $t_{pb} = 190$ ms as a representative. We show in Fig. 2 the energy-integrated angular distributions of ν_e (bottom row), $\bar{\nu}_e$ (middle row) and their difference (top row) in the meridian sections of momentum space for different values of the azimuthal angle ϕ_ν . The colors indicate neutrino species: red and blue arrows represent ν_e and $\bar{\nu}_e$, respectively for the bottom and middle rows; they indicate which species is dominant for the top row. The length of each arrow indicates the value of the distribution function for the bottom and middle rows and its difference between the two species for the top row in the direction that the arrow specifies. The vertical direction corresponds to the local radial direction. Each column shows the meridian section which corresponds to a pair of azimuthal angles in momentum space: panels A, B and C have $\phi_\nu[\text{radian}] = (0.35, 3.49)$, $(1.57, 4.71)$ and $(2.78, 5.92)$, respectively. Note that the scales are almost the same in these panels, which is in sharp contrast with our previous results [21], where ν_e overwhelms $\bar{\nu}_e$ by a large factor. One can see that both distributions are almost isotropic. However, some differences, albeit small,

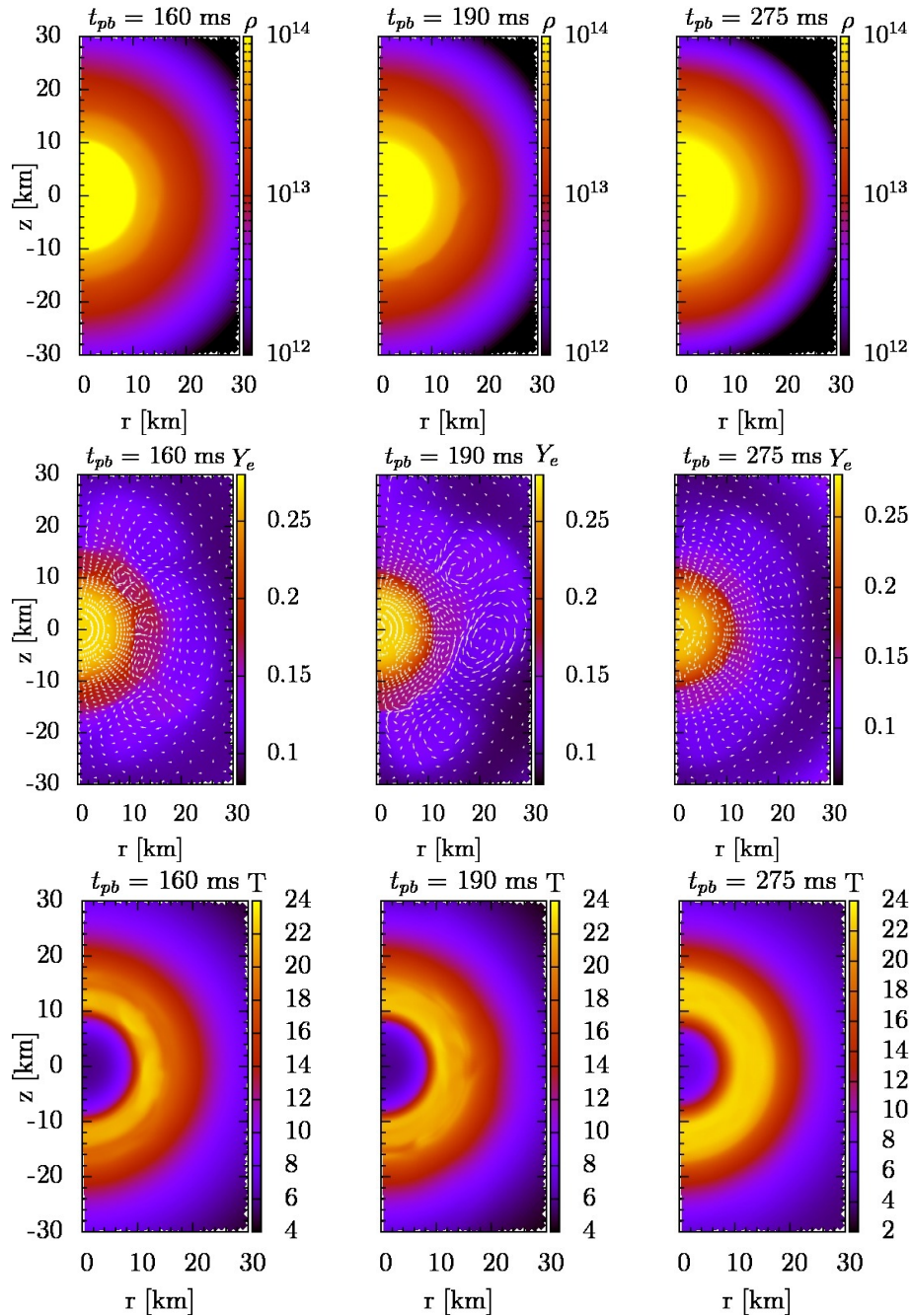


FIG. 1: The density (top), Y_e (middle) and temperature (bottom) distributions in the half meridian section at three different post-bounce times: $t_{pb} = 160$ ms (left), 190 ms (center) and 275 ms (right). In the middle row the arrows show the matter velocities.

do exist, which give rise to the crossing. Indeed, ν_e is still dominant in the radially outward direction, $\bar{\nu}_e$ prevails in the opposite direction as shown on the top row. It is also recognized that the distributions are not axisymmetric.

We think that the small differences in the angular distributions between ν_e and $\bar{\nu}_e$ are not numerical artifacts. Firstly, the neutrino distributions in momentum space are almost isotropic without fine features on small angular scales in the region of our concern as mentioned. This

implies that many angular mesh points are not necessary. Secondly, we already studied the accuracy of our simulations with the same angular resolution in other papers (see, for example, [34] and [35]) and found that it is very good in this region.

The angular distribution differences between ν_e and $\bar{\nu}_e$ are also exhibited in the upper row of Fig. 3. The surfaces in these 3D plots are the ELN defined in Eq. (14) as a function of the propagation direction. The blue part cor-

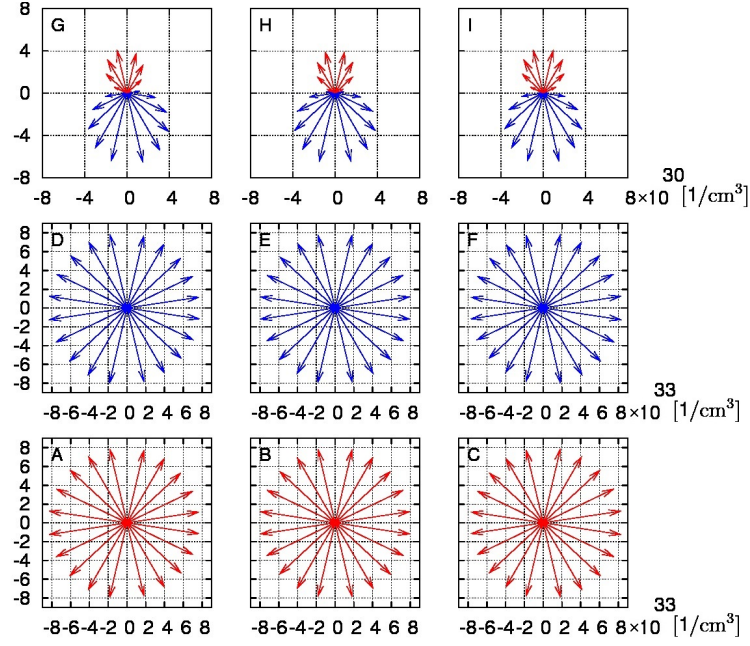


FIG. 2: Energy-integrated angular distributions of ν_e (bottom row) and $\bar{\nu}_e$ (middle row) and their difference (top row) on three meridian sections in momentum space at $t_{pb} = 190$ ms. The red (blue) color represent ν_e ($\bar{\nu}_e$) for the bottom and middle rows whereas it implies that ν_e ($\bar{\nu}_e$) is dominant on the top row. The left, center and right columns display the meridian sections for the pairs of ϕ_ν [radian] = (0.35, 3.49), (1.57, 4.71) and (2.78, 5.92), respectively.

responds to the directions, in which $\bar{\nu}_e$ is dominant over ν_e and the opposite is true for the red section. The z-axis is aligned with the local radial direction and the x-axis is chosen so that it is contained in the spatial meridian section through that point. The four panels give the same angular distribution but viewed from different directions. It is clear that $\bar{\nu}_e$ overwhelms ν_e in the radially inward directions.

We now apply the linear analysis to this ELN. The dispersion relations (DR) between the absolute value of the wave number k and the frequency ω of perturbations are displayed in the bottom panels of Fig. 3 for different directions of \mathbf{k} as indicated by arrows in the upper panels at $t_{pb} = 190$ ms. All the patterns in these DRs given here were actually observed in our previous paper [21] for artificially modified ELNs. This time the ELN is real, being obtained in our realistic self-consistent simulations. From the results in [21] we can say that this flavor state is unstable to the fast pairwise conversion and the growth rate will be largest for \mathbf{k} directed radially inward, which is confirmed indeed in Fig. 4. The instability is found for a much broader range of directions of \mathbf{k} with similar growth rates, though, as understood from the same figure. The maximum growth rate is 1.3 cm^{-1} and the conversion occurs really fast compared with the hydrodynamical time scale. Note that the actual direction, in which perturbations grows, is not the direction of \mathbf{k} but given by the group velocity.

Now we come to the core question: why the crossing occurs in this region of the core. In our previous pa-

per we contended that the crossing is not easy to obtain, since the angular distributions of ν_e and $\bar{\nu}_e$ become more different as one goes deeper, where neutrinos are coupled with matter and their distributions are affected by matter motions, but the population of $\bar{\nu}_e$ will be suppressed by stronger Fermi-degeneracy of ν_e and vice versa at larger radii, where $\bar{\nu}_e$ is more abundant but the angular distributions become not much different between ν_e and $\bar{\nu}_e$. We find that in the regions, where the crossings occur, $\bar{\nu}_e$ is not suppressed and populated almost in comparable numbers to ν_e in fact.

In order to see this more quantitatively, we define Γ as the ratio of the number density of $\bar{\nu}_e$, $n_{\bar{\nu}_e} = \int \frac{d^3p}{(2\pi)^3} f_{\bar{\nu}_e}$ to that of ν_e , $n_{\nu_e} = \int \frac{d^3p}{(2\pi)^3} f_{\nu_e}$:

$$\Gamma \equiv \frac{n_{\bar{\nu}_e}}{n_{\nu_e}}. \quad (16)$$

Although we have investigated many snapshots at different post-bounce times, we show the color contour plots of Γ for only three selected times ($t_{pb} = 160, 190$ and 275 ms) in Fig. 5. We find that there are indeed some regions, where Γ becomes close to unity near the PNS surface as early as $t_{pb} = 160$ ms. It is clear the maximum value of Γ for $t_{pb} = 160$ ms, is 0.9 but it reaches to 1.2 and 2.5 at $t_{pb} = 190$ and 275 ms, respectively. As we mentioned, we found the crossing for the first time at $t_{pb} \sim 190$ ms and we confirmed that the unstable regions continue to exist at later times. It seems indeed that the crossing occurs when Γ reaches unity for the first time

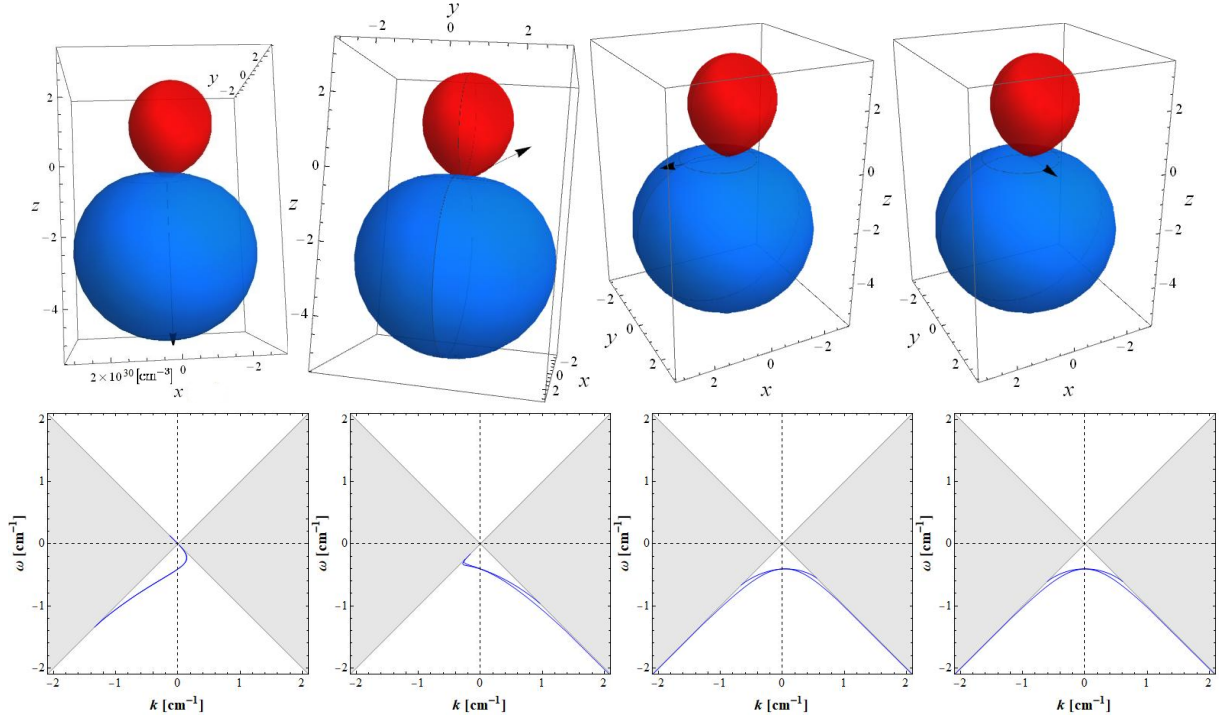


FIG. 3: Angular distribution differences between ν_e and $\bar{\nu}_e$ (upper panels) and the DRs (lower panels) for different wave vectors \mathbf{k} of perturbations given in the upper panels at $t_{pb} = 190$ ms for $r = 16.5$ km and $\theta = 2.1$ rad. In the upper panels, ELN defined in Eq. (14) is shown as a surface with the red (blue) color indicating the dominance of ν_e ($\bar{\nu}_e$). The z -axis corresponds to the local radial direction. Black arrows indicate the direction of wave vector \mathbf{k} . In the bottom panels, the blue lines represent different branches of stable modes.

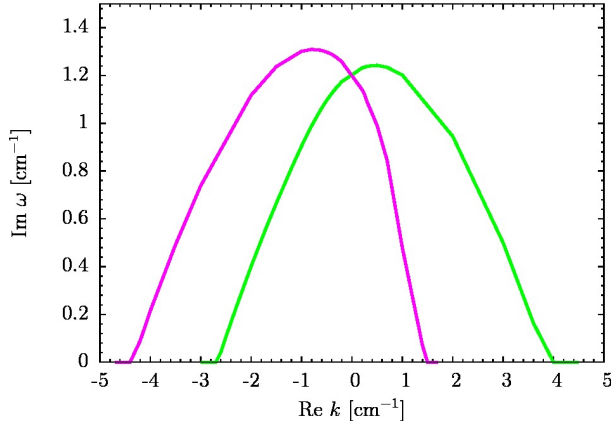


FIG. 4: Growth rates of unstable modes as a function of k for the radially inward direction (magenta) and the direction perpendicular to it (green).

(see the middle panel of Fig. 5). Note that the contour lines in the second and third panels ($t_{pb} = 190$ and 275 ms) in Fig. 5 show $\Gamma = 0.9, 1.0$ and, 1.1 respectively. It is seen that the crossing occurs (black cross signs) when Γ is very close to unity. This is understandable. Inside the neutrino sphere, the angular distributions of ν_e and $\bar{\nu}_e$ are both nearly isotropic (see Fig. 2) and, in order to get the crossing, their number densities should be almost the

same. As mentioned above, Γ reaches unity for the first time at about $t_{pb} = 190$ ms, when we find the crossings to happen. Although the maximum value of Γ exceeds unity at later times, the ELN crossings still occur in the regions, where the value of Γ is ≈ 1 at these times.

The widths of the unstable regions look very narrow but are still much wider than the conversion length scale, in fact. This is understood from Fig. 5 as follows: the difference between the values of Γ is $\Delta\Gamma = 0.1$ for the adjacent contour lines that are separated spatially by $\sim 0.5 \times 10^3$ m; since the ELN is a fraction $\sim 10^{-3}$ of the individual number densities of ν_e and $\bar{\nu}_e$ in the region, where the ELN crossing occurs, the width of that region may be estimated as ~ 5 m, which should be compared with the length of the fast conversion, ~ 0.01 m (see Fig. 4).

In Fig. 6 we present, instead of their ratios, the number densities of ν_e (left panel) and $\bar{\nu}_e$ (right panel) themselves at $t_{pb} = 190$ ms as color contours. On top of them are the flux vectors of each species of neutrinos. In each panel, the left half is for the laboratory frame and the right half is for the fluid-rest frame. It is observed that ν_e becomes more abundant at smaller radii whereas $\bar{\nu}_e$ is almost absent at $r \lesssim 10$ km. This is due to the strong Fermi-degeneracy of ν_e in this region. It is also obvious that $\bar{\nu}_e$ is most abundant off center at $r \sim 15$ km (see the yellow half-circles in the right panel of Fig. 6). This cor-

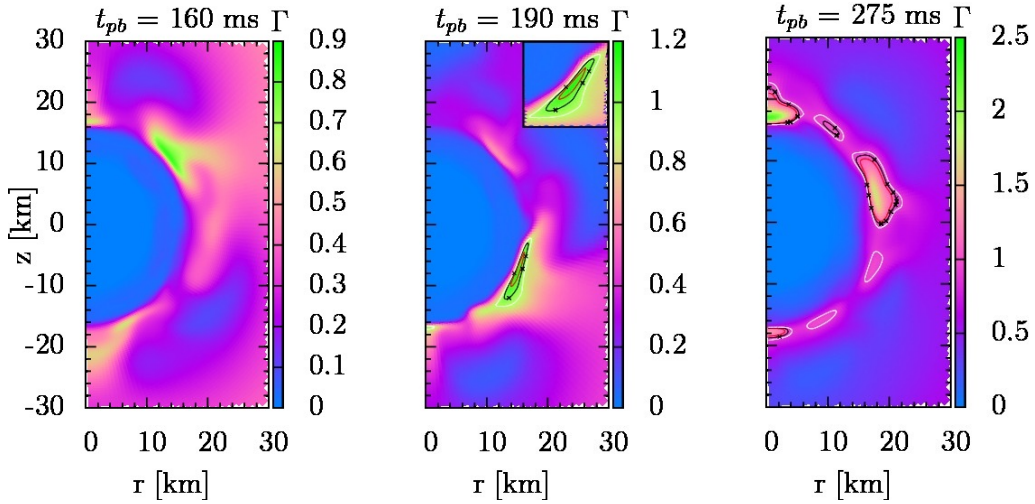


FIG. 5: The contour plots of the ratio of the number density of $\bar{\nu}_e$ to that of ν_e number densities, Γ , for the three different times. Black crosses for $t_{pb} = 190$ and 275 ms indicate the places where the crossings are observed. Three contour lines correspond to $\Gamma = 0.9$ (white), 1.0 (black), and 1.1 (red). It is evident that the crossings occur when the value of Γ is very close to 1.

responds to the region, where the ratio Γ is largest. The neutrino fluxes are trailing the convective matter motions irrespective of species in the laboratory frame as can be seen in the left half of each panel. In the fluid-rest frame, on the other hand, the fluxes are determined mainly by diffusions. As a result, the ν_e flux is mostly directed outwards, reflecting the fact that their number density is decreasing with radius rather monotonically. Some non-radial flows are recognized, though, in the regions, where the convective matter motions produce fluctuations of Y_e (see Fig. 2) as well as of n_{ν_e} and $n_{\bar{\nu}_e}$; on the other hand, $\bar{\nu}_e$'s are diffusing inward at $r \lesssim 15$ km and outward at $r \gtrsim 15$ km. This is just as expected, since their number densities are peaked at $r \sim 15$ km as mentioned above. As a consequence of this fact, ν_e and $\bar{\nu}_e$ are flowing in the opposite directions at $r \lesssim 15$ km. This gives the crossings we found above.

We now consider why $\bar{\nu}_e$ has a peak in the number density at $r \sim 15$ km. The crucial quantity is the chemical potential of ν_e defined as

$$\mu_{\nu_e} = (\mu_e + \mu_p) - \mu_n, \quad (17)$$

where μ_e , μ_p and μ_n are the chemical potentials of electrons, protons and neutrons, respectively. It is true that this has a rigorous thermodynamical meaning only when neutrinos are in chemical equilibrium, but the weak-equilibrium is almost established in this region in fact and hence μ_{ν_e} serves us as a measure to gauge whether $\bar{\nu}_e$ is suppressed by Fermi-degeneracy or not. In Fig. 7 we give the color maps of $\eta_{\nu_e} \equiv \mu_{\nu_e}/T$ at the same three post-bounce times. Crosses in this figure indicate again the places, at which we find the ELN crossings. One quickly recognizes that there appear a pocket of regions, where $\eta_{\nu_e} \lesssim 0$ (light-bluish regions), at $t_{pb} \sim 190$ ms and that the boundaries of these regions coincide with the re-

gions we see the ELN crossings. Deeper inside, $\eta_{\nu_e} \gg 1$ is satisfied and it is consistent with the previous observation that the presence of $\bar{\nu}_e$ is strongly suppressed.

Fig. 8 shows the radial profiles of temperature T and Y_e as a function of density both for the present model and the 1D counterpart. Note that for the 2D case, we plot them for different radial directions. The results are a single line for 1D but a band (a superposition of many lines actually) for 2D with the width reflecting the fluctuations among different angles, for T and Y_e individually. It is seen that the values of Y_e in 1D is larger than the band for 2D around the region of our concern ($r \sim 16.5$ km and $\rho \sim 2.4 \times 10^{13} \text{ g/cm}^3$); in particular, the values of Y_e in the direction, in which the ELN crossing occurs in 2D, are running along the bottom of the band except for $\log(\rho[\text{g/cm}^3]) \gtrsim 13.5$. Note that it is known that the PNS convection facilitates the deleptonization of PNS [36, 37]. It should be mentioned that the above 1D model was computed not with the Furusawa-Shen EOS (FSEOS) as for the 2D model but with the Furusawa-Togashi EOS (FTEOS)[38] based on the variational method with realistic nuclear potentials, since that is the only model for the same progenitor available to us at the moment. As we will show later at the end of this section, this will not be a concern.

The fact that ν_e and $\bar{\nu}_e$ are flowing in the opposite directions at small radii ($r \lesssim 10$ km) is well known in the proto-neutron star (PNS) cooling. Since this happens in 1D already, the multi-dimensional effect should be important for the ELN crossing, which is not observed in 1D. In Fig. 9, we present the degeneracy parameter of ν_e again as a function of density for both 1D and 2D. We find that the degeneracy of ν_e is not resolved ($\eta_{\nu_e} \sim 1$) in 1D although it is lowered substantially in the region of our concern ($r \sim 16.5$ km and $\rho \sim 2.4 \times 10^{13} \text{ g/cm}^3$) unlike in 2D; as a result, the ratio Γ never comes close to

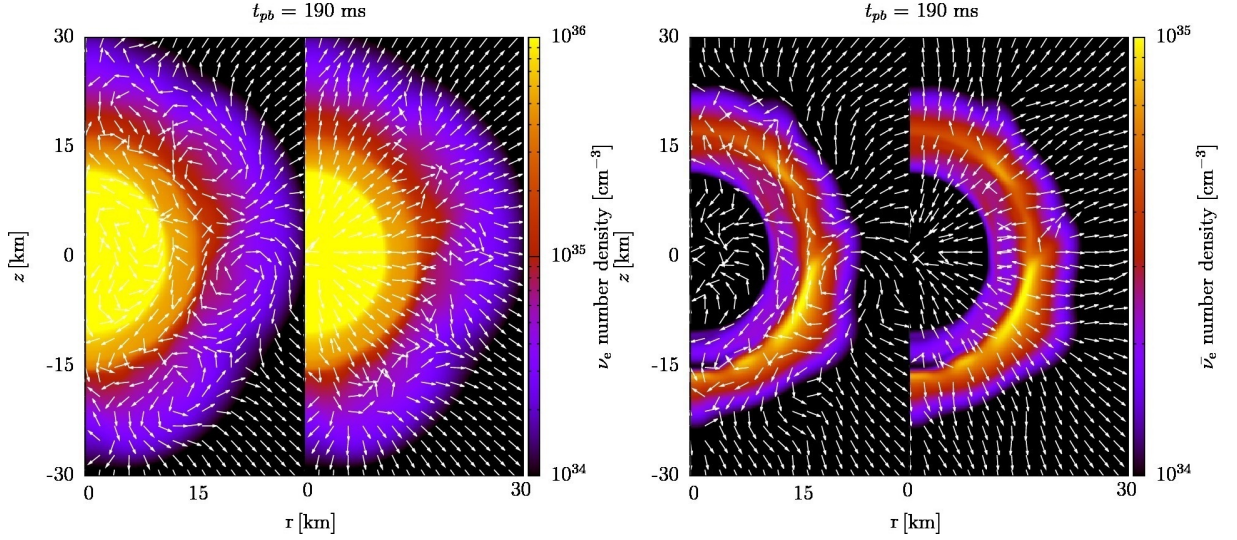


FIG. 6: The number fluxes (arrows) and color contours of the number densities of ν_e (left two panels) and $\bar{\nu}_e$ (right two panels), respectively. The left and right halves of each panel are for the laboratory and fluid-rest frames, respectively.

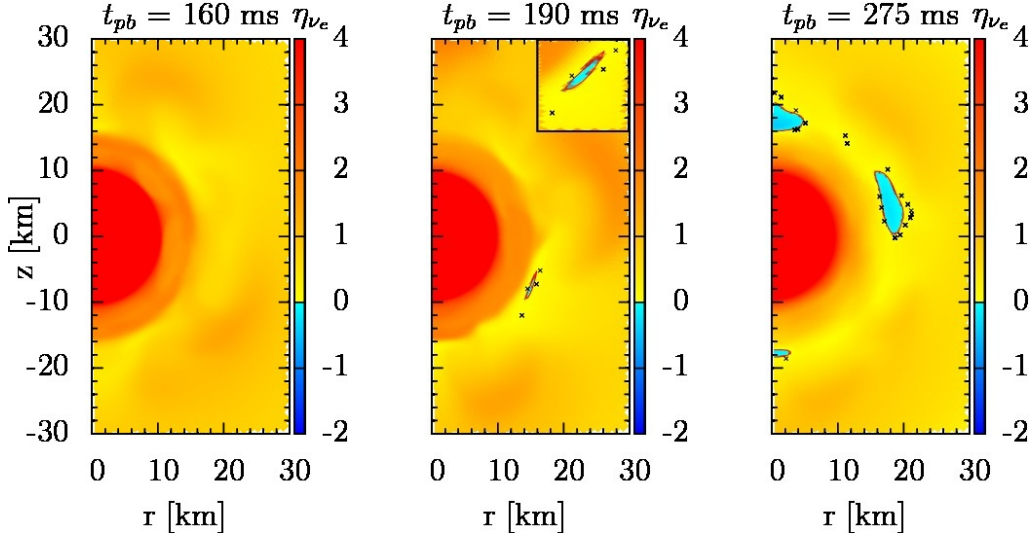


FIG. 7: The color contour plots of $\eta_{\nu_e} = \mu_{\nu_e}/T$ at different times. Blue regions indicate the places, where η_{ν_e} is slightly smaller than 0 and the contour line indicates the points, at which $\eta_{\nu_e} = 0$. The crosses show the points of ELN crossings.

unity in 1D, implying no chance for ELN crossing.

To give some numbers of relevance for one of the representative points ($r = 16.5$ km and $\theta = 2.1$ rad), where detailed analyses are done so far, the values of density, temperature and electron fraction are $\rho = 2.4 \times 10^{13}$ g/cm³, $T = 20.4$ MeV and $Y_e = 0.13$, respectively, and the chemical potentials of electron, proton and neutron including the rest mass are $\mu_e = 56.9$ MeV, $\mu_p = 848$ MeV and $\mu_n = 903$ MeV, respectively. The electron is hence strongly degenerate here as expected. It is also found, however, that this large chemical potential of electron is almost canceled by the chemical potential difference between proton and neutron.

The mass fraction of free neutron and proton are X_n

$= 0.781$ and $X_p = 0.069$, respectively. If they were non-interacting ideal Boltzmann gases, they would have smaller chemical potential difference by ~ 7 MeV than those given above. Nuclear interactions are hence playing roles here. What is more important is the existence of light elements in the region. In fact, the mass fraction of light nuclei ($Z < 6$) is $X_l = 0.138$ whereas that of heavy elements ($Z \geq 6$) is $X_h = 1.16 \times 10^{-2}$ at the point of our current concern. Since free protons are largely absorbed in the light elements, they become less abundant and the chemical potential difference between neutron and proton is enhanced.

The existence of the light elements in the region of our concern is observed also in the 1D model, however.

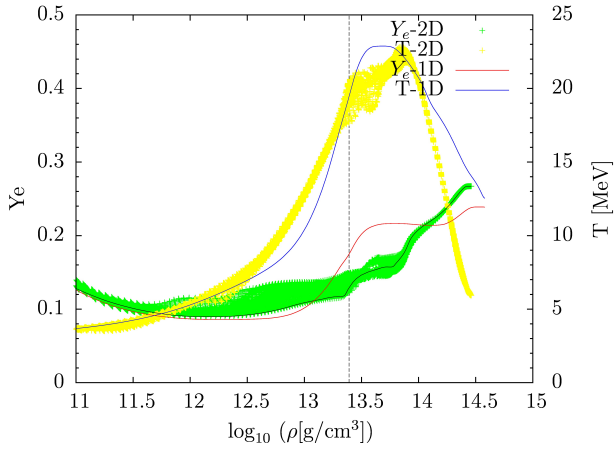


FIG. 8: Radial profiles of T and Y_e as a function of density for 1D (blue and brown lines for T and Y_e , respectively.) and 2D (yellow and green points, respectively). The black line inside the 2D plots for Y_e indicates the Y_e profile along the radial ray in the direction of the ELN crossing. Black dashed line indicates the density, at which the ELN crossing is observed in 2D.

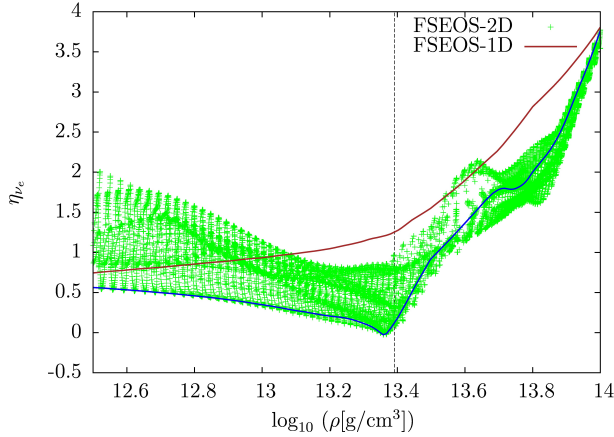


FIG. 9: Radial profile of the degeneracy parameter of electron-type neutrino η_{ν_e} as a function of density for 1D and 2D. Blue solid line corresponds to the profile of η_{ν_e} along the radial ray in the direction of the ELN crossing. Black dashed line indicates the density of our focus.

In fact, the mass fractions of relevance are $X_n = 0.726$, $X_p = 0.098$, $X_l = 0.173$ and $X_h = 2.1 \times 10^{-4}$. It is noted again that this 1D model was computed with the Furusawa-Togashi EOS (FTEOS) [38] but the mass fractions are recalculated in post-process with the FSEOS for the density, temperature and Y_e obtained in the simulation. As will be shown shortly in the next paragraph, this inconsistency is not a serious problem. It turns out that these numbers do not give $\eta_{\nu_e} \sim 0$ (see Fig. 9). The critical difference comes from Y_e . As mentioned above, the 2D model realizes smaller values of Y_e in the direction of the ELN crossing owing to convective matter motions. In the 1D model, the higher values of Y_e owing to the absence of such non-radial motions give higher val-

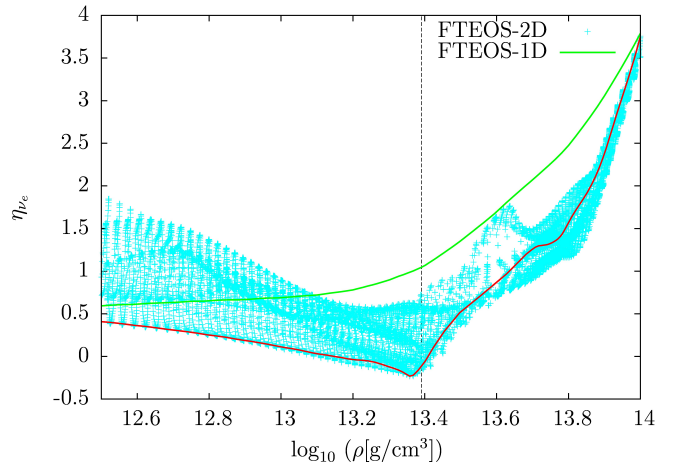


FIG. 10: Same as Fig. 9 but for FTEOS.

ues of X_p and lower values of X_n , both of which tend to reduce the chemical potential difference. The result is the larger value of the chemical potential of ν_e and the smaller value of Γ , which suppress the ELN crossing, as mentioned earlier. It is now apparent that the combination of the smaller values of Y_e induced by convective matter motions and the emergence of the light elements is the ultimate cause of the ELN crossing we found in this paper.

Finally, in Fig. 10 we show η_{ν_e} obtained with the Furusawa-Togashi EOS (FTEOS) [38], which is based on the variational method with realistic nuclear potentials and satisfies various constraints for the hadronic EOS well and is hence favored nowadays. It is clear that the values of η_{ν_e} are not much different from those derived with the FSEOS. Indeed, we can see $\eta_{\nu_e} \sim 0$ at almost the same point. This suggests that the EOS-dependence is rather weak although more work is certainly needed.

IV. SUMMARY AND DISCUSSIONS

We have extended our previous pilot study and conducted a full survey of the post-bounce core of $11.2 M_\odot$ progenitor, searching for the electron lepton number (ELN) crossing. We have reported positive detections of the crossings at $r \sim 16 - 21$ km from the post-bounce time of $t_{pb} = 190$ ms onward. These regions are located inside the neutrino sphere and may have an impact on the CCSNe explosion. Conducting linear analysis for the representative points and times, we have confirmed that the crossing really induces the fast flavor conversion at least in the linear level and that the conversion occurs very quickly. As a result, although the regions, where we have found the crossings, are confined to narrow shells, they are still wide enough for the conversion to occur.

We have then studied rather in detail why the crossing is obtained at these regions alone. We have observed

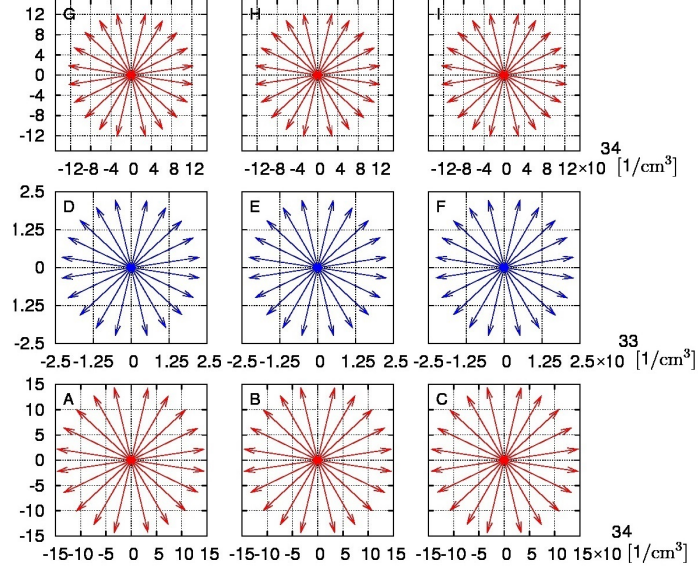


FIG. 11: Same as Fig. 2 but for $t_{pb} = 160$ ms. Note that the scales are different for ν_e (red) and $\bar{\nu}_e$ (blue).

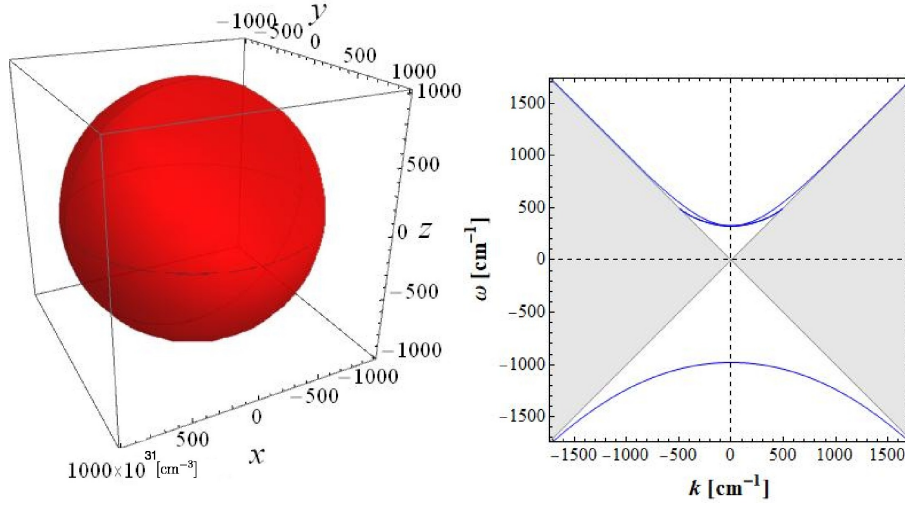


FIG. 12: Same as Fig. 3 but for $t_{pb} = 160$ ms. The red color indicates that the ν_e is dominant over $\bar{\nu}_e$.

that the population of $\bar{\nu}_e$ is comparable to or even larger than that of ν_e in those regions and that the ratio Γ of the former to the latter tends to increase in time. Since the angular distributions of ν_e and $\bar{\nu}_e$ are both nearly isotropic, the ELN crossing is possible only where Γ is very close to 1, which is really the case.

We have found that the chemical potential of ν_e becomes almost vanishing although the electrons are still strongly degenerate and demonstrated that the large chemical potential of electron is almost canceled by the chemical potential difference of neutron and proton. This became possible by the interplay of nuclear physics and hydrodynamics, i.e., the appearance of light elements,

which take in most of the free protons, and lower values Y_e produced by the convective motions of matter. This phenomenon seems to depend on EOS rather weakly and is may occur commonly.

It should be mentioned finally that the fact that we have found the ELN crossings and hence the fast flavor conversions in some regions inside the neutrino sphere is one thing but whether they have really some implications for supernova explosions is quite another. For one thing, the solid angle that these sporadically regions subtend is not very large although it tends to widen in time. We had better look into other simulation results for different progenitors and/or with different EOS's. Axisymmetry we

imposed in our simulations should be removed. All three flavors should be taken into account. The nonlinear evolution of the fast flavor conversion should be investigated [39, 40]. And eventually we need to perform CCSN simulations with these fast flavor conversions begin somehow implemented [41].

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Appendix

We present the results of our analysis for $t_{pb} = 160$ ms as a representative of earlier post-bounce times, at which we find no ELN crossing for the ν_e and $\bar{\nu}_e$ angular distributions. As demonstrated in our previous paper [21] in the early post-bounce phase ν_e are substantially

more abundant than $\bar{\nu}_e$. We pick up the same point ($r = 16.5$ km and $\theta = 2.1$ rad) where the ELN crossing is detected at $t_{pb} = 190$ ms and we gave detailed analyses in the main body of the paper. In Fig. 11 we show the angular distributions of ν_e and $\bar{\nu}_e$ (cf. Fig. 2). It is clear that the scales are different between ν_e and $\bar{\nu}_e$ and ν_e is highly dominant over $\bar{\nu}_e$ in all directions.

For comparison, we give some relevant numbers: the density, temperature and electron fraction are $\rho = 2.36 \times 10^{13} \text{g/cm}^3$, $T = 18.4$ MeV and $Y_e = 0.1$, respectively; the mass fractions are $X_n = 0.827$, $X_p = 0.051$, $X_h = 4.58 \times 10^{-3}$, and $X_l = 0.116$. The chemical potentials of electron, proton and neutron are $\mu_e = 66.5$, $\mu_p = 861$ and $\mu_n = 907$ MeV including rest masses, respectively. Since we obtain $\eta_{\nu_e} \gtrsim 1$, it is natural that $\bar{\nu}_e$ is suppressed. Fig. 12 shows the angular distributions difference between ν_e and $\bar{\nu}_e$ (left panel) for \mathbf{k} oriented in the radial direction. Note that the red color means that the ν_e is dominant over $\bar{\nu}_e$ (cf. Fig. 3) and the corresponding dispersion relation DR (right panel). The DR has a pattern typically observed when there is no ELN crossing [21] where branches that have a gap in the frequency of perturbation ω .

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