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DOI: 10.1103/PhysRevC.104.024308

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(Dated: June 29, 2021)

The β -decay of ^{125m,125}Ag into levels in ¹²⁵Cd was investigated at the Holifield Radioactive Ion Beam Facility (HRIBF). Uranium-238 targets were bombarded with 50-MeV protons with an intensity of 15 μ A, and the induced fission products were mass separated and deposited on a moving tape in the center of the VANDLE array consisting of γ -detectors and plastic scintillators. A partial decay scheme has been assigned for both β -decay of the (9/2⁺) ground state of ¹²⁵Ag and its low-lying (1/2⁻) isomer, with the energy of the low-lying (11/2⁻) isomeric state in ¹²⁵Cd assigned as 188.5 keV. In addition, β -delayed neutron emission probabilities were also determined, as 1.2(2)% for the (9/2⁺) ¹²⁵Ag ground state and 4.6(10)% for the (1/2⁻) isomer, respectively, that are substantially lower than the previously reported value.

PACS numbers: 23.40.-s, 21.10.-k, 21.60. Ev, 23.35.+g

I. INTRODUCTION

In the region of the Segre chart near the doubly magic nucleus ¹³²Sn there is an abundance of long-lived isomeric states. This is largely due to low-lying $\pi g_{9/2}$ and $\nu h_{11/2}$ configurations near several low-j orbitals. One of the best systems to pursue such studies is that of the neutron-rich cadmium nuclei, as they are near the closed proton and neutron shells at Z = 50 and N = 82, respectively. The odd-A isotopes of cadmium, near N = 82 in particular, are all known to exhibit low-energy isomers with halflives similar to those of their ground states. A systematic study of these levels allows one to map the single particle states of these and nearby configurations across the chart. For nuclei that are far from stability, β -decay is the best way to populate these low-energy and low-spin levels. For the cases of odd-mass Ag and Cd isotopes, the study of the β -decay of these nuclei and their isomers gives information on the $\pi 1g_{9/2}$ and $\pi 2p_{1/2}$ states in Ag and the $\nu 1h_{11/2}$ and $\nu 2d_{3/2}$ states in Cd, and the levels built on them.

On the neutron-rich side of stability, the levels of oddmass Ag isotopes from ¹¹¹Ag to ¹²³Ag [1–7] have been extensively measured via the β -decay of Pd isotopes. In all these cases long-lived isomers were observed in the Ag isotopes. The spins and parities of the ground state of ¹¹¹Ag and ¹¹³Ag have been measured via atomic beam as $1/2^{-}$ [8] with the isomer assigned as $7/2^{+}$ based on the β and γ decay patterns. Silver-115 and ¹¹⁷Ag have been assigned as $(1/2^{-})$ for the ground state and $(7/2^{+})$ for the isomeric levels based on the systematics of these lighter Ag isotopes [3, 4]. The energy of the $1/2^{-}$ and $7/2^{+}$ levels were reported to have been reversed in ¹¹⁹Ag based on the β decay pattern of ¹¹⁹Pd [9] with the $(7/2^{+})$ level as the ground state. As such, the ground state of ^{121,123}Ag have also been assigned as $(7/2^{+})$ based on systematics.

For levels in the Cd isotopes, β -decay studies of oddmass Ag isotopes have identified levels in Cd up to ¹²³Cd [6, 10] including a long-lived 1.81(3)s (11/2⁻) isomer whose excitation energy was assigned as 316.4 keV based on decay pattern and systematics. A later study by Kankainen *et al.* involving the JYFLTRAP [11] assigned a value of 144(4) keV to the excitation of the isomer. For the case of ¹²⁵Cd, a long-lived isomer was observed with a half-life of 0.48 (3) s by Huck *et al.* [10] and the excitation energy measured via JYFLTRAP [11] as 186(5) keV and 190(26) keV in the TITAN experiment at TRI-UMF [12]. In addition, long-lived isomers have also been observed in ¹²⁷Cd [13, 14] and ¹²⁹Cd [15]. In neither of these two cases has the relative isomer excitation been established.

TABLE I: List of γ -rays resulting from the beta decay of ground state of ¹²⁵Ag to levels in ¹²⁵Cd. The intensities in columns 2 and 4 are normalized to the 100% 384-keV transition. The I_{tot} intensities in column 4 are the total transition strength (I_{γ} +c.e.). Gamma intensities are based on a) singles, or b) coincidence relationships. The * symbol denotes that the γ was not placed in the level scheme. Contaminants from the β -decay of ¹²⁵Cd are marked with the symbol **, and ¹²⁵In with ***.

$\overline{\mathrm{E}_{\gamma}}$ (keV)	$I_{\gamma}(\%)$	α_{tot}	$I_{tot}(\%)$	E_i	\mathbf{E}_{f}	Mult	coincident γ -rays
$\frac{1}{50.5(1)}$	$17.1(3)^a$	3.30	72.5(12)	239.0	188.5	(M1)	368, 405, 514, 564, 634, 748, 812, 841
	. ,					. ,	882, 926, 1031, 1316, 1654*, 1888
							2106, 2374, 2479
222.4(2)	$16.3(4)^{b}$		17.6(5)	606.6	384.1		384, 514, 693
259.3(2)	$9.0(10)^{b}$	0.044	8.8(10)	643.7	384.1	(M1, E2)	384, 397, 455, 478, 656
292.7(3)	$1.3(4)^{b}$		1.3(4)	1615.4	1322.7		384, 1323
309.1(2)	$1.4(4)^{b}$		1.4(4)	1632.0	1322.7		384, 939, 1323
334.0(6)*	$0.30(15)^{b}$		0.30(15)				191, 574
367.6(2)	$16.4(10)^{b}$		16.2(10)	606.6	239.0		51, 514, 693, 949, 1031
384.1(1)	100^a	0.0041	100	384.1	0.0	(E1)	191, 222, 259, 293, 397, 451, 457
							514, 524, 575, 615, 656, 667, 693, 704
							714, 737, 797, 839, 869, 879, 939, 960
							994, 1042*, 1084*, 1115*, 1231, 1248
							$1404, 1460, 1543^*, 1601 1639, 1738$
							$1819, 1882, 1941^*, 2255, 2278$
							2643, 2843, 2897, 2999, 3074
396.6(3)	$6.4(4)^{b}$		6.3(4)	1040.0	643.7		259, 384, 405, 575, 644, 704, 714
404.7(3)	$4.6(6)^{b}$	0.004	4.5(6)	643.7	239.0	(E1)	51, 397, 656
451.4(2)	$1.2(4)^{b}$		1.2(4)	1632.0	1180.8		384, 797
454.8(3)	$1.9(5)^{a}$		1.9(5)	1754.3	1299.7		259, 384, 607, 644, 656, 693
457.4(3)	$1.3(4)^{a}$		1.3(4)	1638.0	1180.8		384, 797
477.5(2)	$2.5(5)^{b}$		2.5(5)	1121.1	643.7		259, 384, 644
514.4(4)	$10.7(27)^a$		10.7(27)	1121.1	606.6		51, 222, 368. 384, 607
523.5(4)	$0.58(12)^{o}$		0.57(12)	907.6	384.1		384
537.1(4)	$1.3(2)^{b}$		1.3(2)	1180.8	643.7		644
563.7(5)	$0.6(4)^{a}$		0.6(4)	1550.7	986.9		51, 748
574.5(2)	$3.2(9)^{\circ}$		3.1(9)	1180.8	606.6		222, 334*, 384, 607,1558*
575.4(2)	$1.6(3)^{\circ}$		9.1(18)	1615.4	1040.0		384, 397, 656, 1057*
606.6(1)	$32.8(5)^{\circ\circ}$		32.4(5)	606.6	0.0		191^{**} , 455, 514, 575, 693, 948, 1031, 1514
b15.1(4)	$1.9(3)^{\circ}$		1.9(3)	999.2	384.1		384
034.4(3)	$5.9(5)^{-1}$	0.002	5.8(5)	813.5	239.0	(E9)	$\frac{1}{200}$ $\frac{1}{100}$ $\frac{1}{100}$ $\frac{1}{100}$ $\frac{1}{100}$ $\frac{1}{100}$ $\frac{1}{100}$
643.8(2)	$87.9(25)^{\circ}$	0.003	86.8(24)	643.7 1040.0	0.0	(E2)	398, 478, 537, 656, 912, 994, 1163*, 1477
000.8(2)	24.3(8) E $7(c)^{b}$		24.0(8)	1040.0 1200 7	384.1 642.7		384 250 284 405 455 575 644 704 714 060
666.5(2)	0.7(0) $0.66(21)^{b}$		0.65(21)	1299.7	045.7		259, 584, 405, 455, 575, 644, 704, 714, 900
603.1(4)	$4.6(12)^{b}$		0.05(21)	1000.0 1200.7	504.1 606 6		504 51 222 368 384 607
7043(3)	$1.5(3)^{b}$		$\frac{4.0(12)}{1.5(3)}$	1299.7 1744.3	1040.0		384 207 656
704.3(3) 714.3(4)	$1.3(3)^{b}$ 1.2(3) ^b		1.0(3) 1.2(3)	1744.3 1754.3	1040.0		384, 398, 656
$737\ 2(4)$	1.2(0) $1.1(2)^{b}$		1.2(0) 1.1(2)	1121.1	384.1		146** 191** 384 722** 730**
101.2(4)	1.1(2)		1.1(2)	1121.1	504.1		910**, 1173**
747.9(3)	$4.9(5)^{b}$		4.9(5)	986.9	239.0		51
796.5(2)	$8.9(6)^{b}$		8.7(6)	1180.8	384.1		384, 451, 457, 1402, 1460
812.2(3)	$4.2(8)^{b}$		4.2(8)	1051.2	239.0		51
839.3(2)	$1.2(2)^{b}$		1.2(2)	1223.4	384.1		384
868.8(4)	$1.4(3)^{b}$		1.4(3)	1908.8	1040.0		384, 397, 644, 656
878.7(3)	$2.1(2)^{b}$		2.1(2)	1262.8	384.1		384
882.1(2)	$20.0(19)^b$		19.8(19)	1121.1	239.0		51
Continu	ued on next j	page					

TABLE I: List of γ -rays resulting from the beta decay of ground state of ¹²⁵Ag to levels in ¹²⁵Cd. The intensities in columns 2 and 4 are normalized to the 100% 384-keV transition. The I_{tot} intensities in column 4 are the total transition strength (I_{γ} +c.e.). Gamma intensities are based on a) singles, or b) coincidence relationships. The * symbol denotes that the γ was not placed in the level scheme. Contaminants from the β -decay of ¹²⁵Cd are marked with the symbol **, and ¹²⁵In with ***.

$\overline{\mathrm{E}_{\gamma}}$ (keV)	$I_{\gamma}(\%)$	α_{tot}	$I_{tot}(\%)$	E_i	\mathbf{E}_{f}	Mult	coincident γ -rays
$\frac{1}{912.0(2)}$	$2.5(5)^{b}$		2.5(5)	1555.5	643.7		644
925.5(3)	$2.7(7)^{b}$		2.7(7)	1164.3	239.0		51
932.5(2)	$18.2(4)^{a}$		17.9(4)	1121.1	188.5		1888
939.0(2)	$1.8(3)^{b}$		1.8(3)	1322.7	384.1		384
948.9(3)	$3.9(10)^{b}$		3.9(10)	1555.5	606.6		368, 607
960.2(3)	$0.47(17)^{b}$		0.53(19)	2000.2	1040.0		384,656
975.5(5)	$1.4(2)^{a}$		1.4(2)	1164.3	188.5		
994.2(2)	$4.3(5)^{b}$		4.3(5)	1638.0	643.7		384, 644
1031.4(2)	$3.6(10)^{b}$		3.5(10)	1638.0	606.6		51, 222, 368, 384, 607
$1057.2(5)^*$	$0.4(2)^{a}$		0.4(2)				384, 575
1084.1(5)*	$0.8(2)^{b}$		0.8(2)				384
1114.5(5)*	$0.8(2)^{b}$		0.8(2)				384
1163.0(5)*	$2.6(7)^{b}$		2.6(7)				384
1231.3(2)	$1.5(3)^{b}$		1.7(3)	1615.4	384.1		384
1247.8(2)	$1.9(3)^{b}$		1.9(3)	1632.0	384.1		384
1316.3(2)	$5.9(11)^{b}$		5.8(11)	1555.5	239.0		51
1322.8(2)	$4.4(3)^{a}$		4.4(3)	1322.9	0.0		293, 309
1402.2(3)	$3.3(6)^{b}$		3.2(6)	2583.0	1180.8		384, 797
1459.8(3)	$2.5(5)^{b}$		2.5(5)	2640.5	1180.8		384, 797
1476.5(3)	$2.0(4)^{b}$		2.0(4)	2120.4	643.7		644
1514.0(5)	$1.1(4)^{b}$		1.1(4)	2120.4	606.6		607
1543.3(5)*	$1.1(3)^{b}$		1.1(4)				384
1558.1(4)*	$0.7(4)^{b}$		0.7(4)				$384, 574, 618^{***}$
1601.0(5)	$1.6(3)^{b}$		1.6(3)	1985.1	384.1		384
1638.8(4)	$1.6(3)^{b}$		1.6(3)	2022.9	384.1		384
1654.2(4)*	$1.4(5)^{b}$		1.6(3)				50
1737.8(5)	$1.7(3)^{b}$		1.7(3)	2121.9	384.1		384
1818.5(5)	$0.57(16)^{b}$		0.56(16)	2202.6	384.1		384
1882.2(5)	$1.1(3)^{b}$		1.1(3)	2266.3	384.1		384
1887.8(5)	$3.4(11)^{b}$		3.4(11)	3008.9	1121.1		51, 882
1940.6(5)	$0.9(3)^{b}$		0.9(3)				384
2106.4(5)	$9.4(5)^{b}$		9.3(5)	3227.5	1121.1		51,607,882,933
2254.5(3)	$6.4(4)^{b}$		6.4(4)	2638.6	384.1		384
2277.6(3)	$2.9(3)^{b}$		2.9(3)	2661.7	384.1		384
2336.8(5)	$5.6(7)^{a}$		5.6(7)	2980.6	643.7		644
2374.0(5)	$7.0(18)^{b}$		6.9(18)	2980.6	606.6		51, 222, 368, 384, 607
2478.7(5)	$3.1(10)^{b}$		3.1(10)	3085.3	606.6		51, 222, 368, 384, 607
2534.0(5)	$3.2(11)^{b}$		3.1(11)	3140.6	606.6		222, 368, 384, 607
2583.7(5)	$8.9(9)^{b}$		8.8(9)	3227.5	643.7		644
2642.7(5)	$5.2(15)^{b}$		5.1(15)	3249.3	606.6		51, 222, 368, 384, 607
2843.0(4)	$1.6(3)^{b}$		1.6(3)	3227.5	384.1		384
2897.2(4)	$2.0(4)^{b}$		2.0(4)	3281.3	384.1		384
2999.2(4)	$1.1(3)^{b}$		1.1(3)	3383.3	384.1		384
3074.0(4)	$1.1(3)^{b}$		1.1(3)	3458.1	384.1		384

Other information on levels in 125 Cd comes from thermal neutron-induced fission studies of 239 Pu [16, 17] and fragmentation of a 136 Xe and passed through a separator

[13]. In these studies, high spin ($\geq 15/2$) states are observed to feed the $11/2^-$ isomer. These states would not be expected to be populated in β -decay due to the large



FIG. 1: Total mass 125 γ spectrum gated on the β detectors. In this figure, known γ -rays from ^{124,125}Cd, ¹²⁵In are labeled along with γ 's we assign to the β -decay of ¹²⁵Ag. For the cases where peaks have unresolved contribution from more than one isotope, the major isotope is labeled.

change in angular momentum. Fragmentation studies of ²³⁸U by Rejmund *et al.* [18] observed the previously known 743 γ feeding the 11/2⁻ isomer of ¹²⁵Cd, and 233, 250, 570, 646, 663, 756 and 844 keV γ 's which may be positive-parity yrast transitions feeding the 3/2⁺ ground state of ¹²⁵Cd, of which the lower spin states could be populated in the β -decay of ¹²⁵Ag.

Information on the β decay of ¹²⁵Ag and the resulting populated levels in ¹²⁵Cd is very limited. The half-life of ¹²⁵Ag has been measured via the delayed neutron curve as 166(7) ms [19], 163_{-6}^{+9} ms [20], 146(11) ms [21], also via an ion-implanted decay curve as 150(8) ms [22]. A β -decaying isomer of ¹²⁵Ag was observed at 97.1 keV in the β -decay of ¹²⁵Pd [23]. The J^{π} of the ground state was assigned as (9/2⁺) and the isomer as (1/2⁻) based on the decay pattern of ¹²⁵Pd and shell-model calculations performed using the KSHELL code with a monopole-based universal interaction plus a spin-orbit force [23]. This work also reported γ -rays of 383.6 and 643.2 keV assigned as γ 's that feed the ground state decay of ¹²⁵Cd, along with γ -rays of 352.7, 383.6 and 643.2 keV assigned as γ 's that feed the isomer of ¹²⁵Cd. These γ -rays were not assigned to a decay scheme and no uncertainties were assigned to the energies in this work. The mass of ¹²⁵Ag has been measured the FRagment Separator (FRS) at GSI [11] resulting in Q_{β}=8830(430) and Q_{β n}=4110(430) [24]. The currently known levels in ¹²⁵Cd are limited to an 11/2⁻ isomer and a high spin band built on the 11/2⁻ isomer [13, 25], whose levels would not be expected to be populated in the β -decay of ¹²⁵Ag.

Measurement of the β -delayed neutron branch also provides insight into nuclear structure above the neutron separation energy of nuclei near doubly-closed ¹³²Sn and the A = 130 r-process abundance peak, and is vitally



FIG. 2: Beta-gated γ -ray spectrum coincident with the 384.1-keV transition. The energies of prominent γ -peaks are labeled. Those transitions that are not included in the proposed decay scheme are labeled with a *.

important to the understanding of the astrophysical rprocess. Significant β -delayed neutron emission can potentially affect the final abundances of nuclei produced in this process. Although ¹²⁵Ag is not directly on the rprocess path, a knowledge of its β -n branching ratio can provide information with which to evaluate astrophysical r-process network calculations. The β -n branching ratio for ¹²⁵Ag has been reported as 11.8(10)% in the unpublished thesis [20] and more recently in [21] as 2.2(11)%, without distinguishing between the ground state and isomeric decays in either work.

II. EXPERIMENTAL METHOD

Fission products were produced via the proton-induced fission of ²³⁸U at the Holifield Radioactive Ion Beam Facility (HRIBF). Fifty MeV protons with an intensity of $\approx 15 \ \mu A$ from the ORIC cyclotron were used to bombard a 238 UC_x target [26] in a plasma ion source installed at the IRIS-2 [27]. The proton induced fission products were then mass-separated via a high resolution ($\Delta m/m$ $\approx 10,000$) magnet and delivered to the VANDLE (the Versatile Array of Neutron Detectors at Low Energy) detector array, where they were embedded on a tape located in the center of the array [28]. The array consisted of two high-purity Ge (HPGe) clover detectors which measured γ -transitions, two plastic scintillators surrounding the collection spot for detection of β particles, and 48 plastic scintillators for detection of neutron energies via their time of flight between the two scintillators surrounding the beam spot and the scintillators at a distance of 50 cm. The clover detector's photopeak efficiency was $\approx 4\%$ at 1.33 MeV. The results contained herein were obtained from γ -ray information, while information obtained from the neutron array will be published in the future.

The beam of mass 125 was pulsed via electrostatic plates with the beam on for the first 2 s, and then turned off for 0.5 s for a total of 2.5 s to allow the half-life measurements. This resulted in a 2 s grow-in, followed by 0.5 s decay cycle, during which the ion beam was deflected away by the electrostatic plates. Afterwards the tape moved the collected radioactivity to a Pb shielded area 50 cm away from the collection spot. Coincidences between the HPGe detectors and the scintillators (which were shaped to surround the beam-line) allowed us to remove the vast majority of γ -rays from either the room or longer-lived daughter products on the used tape. The tape cycle time was chosen to enhance the observation of the short-lived ¹²⁵Ag nucleus. Cadmium-125 and ¹²⁵In are present in this data set both from being deposited on the tape directly and as the β -daughter products. The direct production rate of ¹²⁵Cd and ¹²⁵In from ²³⁸U fission are orders of magnitude higher than Ag. With no isobaric separation or tape movement, the ratio of Ag/Cd/In was expected to be [29] $\approx 1/70/650$ for mass 125.

The short tape cycle compared to the half-lives of the Cd and In isomers greatly favors the decays of Ag, however, this will only improve the ratio to $\approx 1/23/63$. After tuning the high resolution magnet to sufficiently enhance the γ -rays resulting from the decay of Ag compared to Cd and In, the resulting efficiency corrected measured ratios



FIG. 3: Beta-gated γ -ray spectrum coincident with the 50.5-keV transition. Prominent γ -peaks that are included in our proposed decay scheme of ¹²⁵Ag are labeled.

(running at the short tape cycle) for $^{125}\text{Ag}/^{125}\text{Cd}/^{125}\text{Im}$ was $\approx 1/1.1/1.7$. The other members of the isobaric chain (Pd, Rh) are refractory metals, which have a very low efficiency for release from the ion source, and are not observed. The only other nuclei observed from any other mass results from the β -delayed neutron emission of ^{125}Ag .

The data acquisition system made use of a digital spectroscopy system based on XIA Pixie16 Rev. F modules (produced by X-ray Instrumentation Associates) [30]. These modules incorporate 250 MHz flash ADC's, and serve as a replacement for amplifiers, discriminators, conventional ADC's and TDC's. All signals from the preamps are connected directly into the PIXIE 16 modules then analyzed via the on-board processors to determine their amplitude by fitting the waveform, and time-stamped by a continuously running clock.

A 16 ns (4 channel) coincidence time gate was used to produce prompt $\gamma - \gamma$ matrices. A time-delayed matrix gated with the same time width was used to subtract the random background from the prompt $\gamma - \gamma$ matrices. To calibrate the efficiency of the Ge detectors, standard sources of ¹³³Ba, ⁶⁰Co and ¹⁵²Eu were used. The relative error in efficiency in the energy range of these isotopes was determined to be 6%. For γ -rays above 1.4 MeV, the extrapolated efficiencies were assigned the following errors: 6% for E \leq 1.8 MeV, 10% for 1.8 MeV \langle E \leq 2.5 MeV, 15% for 2.5 MeV \langle E \leq 3.0 MeV, and 20% for E >3.0 MeV.

III. EXPERIMENTAL RESULTS

Transitions have been placed in the decay schemes from information obtained via $\gamma - \gamma$ and $\beta - \gamma - \gamma$ coincidences. All possible coincidence spectra were analyzed to determine this. The relative intensities of the γ -transitions de-exciting a given level were determined wherever possible both by coincidences with γ -transitions feeding the given level and by gating from below. Gammas that arise from decaying nuclei on the tape from previous cycles and background events will not cause coincidences between the β and γ detectors and can therefore be determined not to arise from ¹²⁵Ag.

Fig. 1 shows the total γ spectrum gated on the β detectors. In this figure, known γ -rays from ¹²⁵Cd, ¹²⁵In, along with transitions that we assign to γ transitions resulting from the β decay and β -delayed neutron emission of ¹²⁵Ag. The previously reported 352.7, 383.6 and 643.2 keV γ 's [23] from ¹²⁵Ag are all observed in this work.

Beta decay from the $(1/2^{-})$ isomer and $(9/2^{+})$ ground state will feed very different states due to the large difference in spin and opposite parities. One would expect that any levels observed in a given state would only be populated by one of the two β -decaying states. The observed gammas can be placed in two groups via their γ - γ coincidences. Group 1 are those γ 's in coincidence with either the 384.1 (see Fig. 2) or 50.5 keV γ transitions (see Fig. 3), and group 2 being those in coincidence with the 352.6-keV γ transition (see Fig. 4). The 50.5 and 384.1 keV transitions are connected via feeding from



FIG. 4: Beta-gated γ -ray coincident with the 352.6-keV transition. Prominent γ -peaks that are included in our proposed decay scheme of 125m Ag are labeled.



FIG. 5: Beta-gated $\gamma\text{-ray}$ spectrum in coincidence with the 514.4-keV transition.

above from several γ -rays, the strongest of which is the 514.4 keV transition (see Fig. 5). We propose that these represent β -decays from two isomers in ¹²⁵Ag: a high-spin (9/2⁺) ground state and a low-spin (1/2⁻) isomeric state, with the assignment based on their respective decay to the low-spin $3/2^+$ and high-spin $11/2^-$ isomers in ¹²⁵Cd. Overall, we have observed a total of 73 γ -rays

from 47 levels that are assigned to the β -decay of the ground state of 125 Ag and 17 γ -rays from 14 levels we assign to the β -decay of 125m Ag.



FIG. 6: Proposed partial decay scheme for the β -decay of ¹²⁵Ag and ^{125m}Ag (part 1 of 2). For each transition, the energy of the transition is listed in keV, followed by the total intensity (γ + c. e.) normalized to the 384.1 keV transition for γ s arising from the β -decay of the ground state and normalized to the 352 keV transition for γ from the beta decay of the 1/2⁻ isomer (shown in italics). The ratios of γ -transitions de-exciting a given state were determined by γ - γ coincidences whenever possible.

A. Decay of the $(9/2^+)$ ground state of ¹²⁵Ag

159(21) ms

97.1, (1/2-)

The observed γ -ray transitions resulting from the β decay of the $(9/2^+)$ ground state of ¹²⁵Ag are listed in Table 1. For γ -transitions with assigned multipolarities the K-conversion coefficients (σ_k) are calculated using the conversion coefficient calculator BRICC [31], along with the assigned J^{π} and energies of their initial and final states. In all cases except for the 50.5 keV γ , these values are relatively small and won't make large changes in the total intensities. The assignment of the multipolarity of the 50.5 keV transition is detailed in the next section. The last column in Table 1 list the γ -rays in coincidence with the given transition. Transitions and levels that require further explanation are listed in the following sections by level.

Table 2 lists the levels in ¹²⁵Cd populated in the β decay of the (9/2⁺) ground state of ¹²⁵Ag. The corresponding assignment of J^{π} for the low-energy levels is based on the decay pattern of the γ -rays de-exciting these levels. The calculated log ft values are based on the apparent β feeding to the given state and should be considered as lower limits (especially in the weaker transitions) in this table and Table 5 due to the potential "pandemonium" effect [32].

Any β -decay from directly to the ground state of ¹²⁵Cd

or its β -decaying 188.5 keV isomer would not be observed in this work. We therefore have attempted to assign a reasonable upper limit based on the J^{π} values of the initial and final levels. The direct β -decay of the (9/2⁺) ground state to the 11/2⁻ isomer in ¹²⁵Cd would require a first forbidden β -decay with a corresponding log ft>6, and partial half-life of > 20 sec resulting in a branching ratio of < 4%. A direct β -decay to the (3/2⁺) ground state would be a third forbidden transition and therefore have a branching ration of ≈ 0 . The resulting decay scheme for the ground state and isomer of ¹²⁵Ag is shown in Figs 6 and 7.

1. 188.5, 239.1, and 606.6 keV levels

The 606.6 keV level is fed from above by several transitions with the 514.4 keV being the largest. Fig. 5 shows the γ -spectrum in coincidence with the 514-keV γ transition. The 606.6 keV state decays via γ -rays of energy 606.6 (to the ground state), 367.6 and 222.4 keV. The 222.4 keV γ decays to the 384.1 keV (5/2⁺) level which then decays by a 384.1 keV γ . The 367.6 keV γ decays to a state at 239.1 keV, and is in coincidence with the 50.5 keV γ , which then decays to a state at 188.5 keV. No evidence is observed for any γ 's of energy 188.5 keV in this



FIG. 7: Proposed partial decay scheme for the β -decay of ¹²⁵Ag and ^{125m}Ag (part 2 of 2). For each transition, the energy of the transition is listed in keV, followed by the total intensity (γ + c. e.) normalized to the 384.1 keV transition for γ s arising from the β -decay of the ground state and normalized to the 352 keV transition for γ from the beta decay of the 1/2⁻ isomer (shown in italics). The ratios of γ -transitions de-exciting a given state were determined by γ - γ coincidences whenever possible.

work. The lack of observed transitions de-exciting this state is strongly suggestive of a relatively long-lived state that decays via β -emission to ¹²⁵In. As noted previously, Ref. [11] reported an 11/2⁻ isomeric state in ¹²⁵Cd at 186(4) keV. We therefore assign the 188.5 keV level as the expected 11/2⁻ isomer. A γ emitted from this state to the 3/2⁺ ground state would have a multipolarity of M4 and would be much slower than the measured T_{1/2} of 480(30) ms [10], and have a vanishingly small branching ratio.

The 50.5-keV transition feeds a $11/2^{-}$ state, the assignment of the multiplicity of this decay will constrain the J^{π} of the 239 keV state and hence the type of β de-

cay to that state. The 50.5 keV γ ray is expected to be highly converted, with a large range of σ_{tot} values and consequently β feeding, depending on the multiplicity of the γ transition. Table 3 details the effect on the total intensity of the transition (with the sum of all observed transitions between levels equal to 100%), β -branching ratio to the 239 keV level and the log ft of the multipolarity assignment of the 50.5 keV γ as E1, M1, E2, or M2.

As detailed in the table, if the 50.5 keV transition is E1, the feeding to this level from above would exceed the amount de-exciting this level despite the fact that a $9/2^+$ assignment would make this an allowed decay,

TABLE II: Levels in $^{125}\mathrm{Cd}$ populated by the $\beta\text{-decay}$ of the $(9/2^+)$ ground state of $^{125}\mathrm{Ag}.$

Level (keV)	$\beta~\%$	log ft
0.0	0	
188.5(2)	<4	>6
239.1(1)	4.2(9)	6.0(1)
384.1(1)	0	
606.6(1)	6.2(15)	5.7(2)
643.7(2)	19.6(10)	5.2(1)
873.5(3)	1.9(2)	6.2(1)
907.6(4)	0.16(3)	7.2(2)
986.9(3)	1.4(2)	6.2(1)
999.2(4)	0.55(7)	6.6(1)
1040.0(2)	7.0(3)	5.5(1)
1050.6(3)	0.19(6)	7.1(2)
1051.2(3)	1.4(3)	6.2(2)
1121.1(3)	12.7(11)	5.3(1)
1164.6(3)	1.4(2)	6.1(2)
1180.8(2)	1.3(4)	6.1(3)
1223.4(2)	0.35(6)	6.8(2)
1262.8(3)	0.60(7)	6.5(1)
1299.7(3)	2.7(11)	5.9(3)
1322.9(2)	1.1(2)	6.3(2)
1550.7(5)	0.19(12)	7.0(3)
1555.5(2)	4.0(3)	5.6(1)
1615.4(2)	1.4(2)	5.9(1)
1632.0(2)	1.4(2)	5.9(2)
1638.0(2)	2.9(3)	5.7(1)
1744.3(3)	0.5(1)	6.5(2)
1754.3(4)	1.0(9)	6.2(4)
1908.8(4)	0.45(10)	6.5(2)
1985.1(5)	0.5(1)	6.4(2)
2000.2(3)	0.15(5)	6.9(2)
2022.9(4)	0.2(1)	6.4(2)
2120.4(3)	0.98(17)	6.3(2)
2121.9(5)	0.5(1)	6.4(2)
2202.6(5)	0.17(5)	6.9(2)
2266.3(5)	0.36(10)	6.5(2)
2583.0(3)	1.0(2)	5.9(2)
2638.6(3)	2.0(1)	5.6(2)
2640.5(3)	0.80(16)	5.7(2)
2661.7(3)	0.91(9)	6.0(2)
2980.6(4)	4.0(6)	5.2(2)
3008.9(5)	1.1(3)	5.5(2)
3085.3(5)	1.0(3)	5.8(2)
3140.6(5)	1.0(4)	5.8(2)
3227.5(5)	6.4(4)	5.0(2)
3249.3(5)	1.7(5)	5.5(2)
3281.3(4)	0.6(1)	5.9(2)
3383.3(4)	0.33(8)	6.2(2)
3458.1(4)	0.35(8)	6.1(2)
>4718	1.2(2)	

which rules out E1. An assignment of E2 (with $J^{\pi} = 7/2^{-}$) would lead to very large apparent beta feeding to the 239 keV level of 35% and an apparent logft value of 5.0, which is far too low for a 1st forbidden transition. A M1 assignment is the only one that gives a reasonable value for β -feeding of 2.7(8)% and logft equal to 6.1 corresponding to a 1st forbidden transitiont. We therefore assign a $J^{\pi} = (9/2^{-})$ to this state.

Further evidence for the assignment of the 239-keV level as $9/2^-$ comes from the level systematics of the lighter odd-mass Cd isotopes. In the cases of 117,119,121,123 Cd, the first excited state above the $11/2^-$ isomer that decays to the isomer is a M1 transition [4–7].

2. 384.1 and 643.7 keV levels

Within error bars the intensity of the γ 's feeding the 384.1 keV level are equal to the intensity of the 384.1 keV γ making the observed β feeding to this state equal to ≈ 0 . The level is fed by the $(7/2^+)$ 644-keV level and de-excites by a 384 γ -keV which feeds the $3/2^+$ ground state ¹²⁵Cd. This is consistent with our proposed J^{π} of $(5/2^+)$, which would be a second forbidden β -decay form the $(9/2^+)$ ground state.

The 643.7-keV level is strongly fed in the β decay of the $9/2^+$ ground state of ^{125}Ag , with an apparent feeding of 18.9(9)% and corresponding log ft = 5.2(1). This strongly suggests an allowed β decay to this state. It decays strongly to the $3/2^+$ ground state, and weakly to 239-keV $(9/2^{-})$ state and the 384-keV $(5/2^{+})$ state. The strong decay to the $3/2^+$ state rules out $11/2^+$ and $9/2^+$, leaving $7/2^+$ as the likely J^{π} of this state. Further evidence for this assignment comes from studies of the neutron-rich ^{121,123,125}Cd nuclei by Rejmund *et al.* [18]. In this study, they observed a 646 keV γ in the prompt (A and Z gated) γ spectrum. It's likely that this γ is the 643.8(2) keV γ observed in this work. This would establish the 643.8 keV transition as the lowest positive parity $7/2^+$ to the $3/2^+$ yrast transition. The other γ 's observed in [18] were not observed in this work.

Figure 8 shows the γ - γ coincidence spectrum generated by gating on the 643.8 keV γ transition. This shows that the (7/2⁺) 643.7 keV level is fed by γ 's from the 1040.0, 1121.1, 1299.7, 1555.5, 1638.0, 2120.4 keV levels (see inset of figure). The decays from these levels can be divided into two groups: 1) levels which feed the (5/2⁺) 384.1 and (7/2⁺) 643.7-keV levels and 2) levels which feed the 606.6 and (7/2⁺) 643.7-keV levels. In the first group are the 1040.0 and 1121.1 keV levels. In the second group are the 1299.7, 1555.5, 1638.0 and 2120.4 keV levels. Of this second group, the levels at 1299.7, 1638.0 are only observed to feed the 643.7 and 606.6-keV levels and have an apparent logft < 6 (consistent with an allowed β decay for the (9/2⁺) ground state), making them possible candidates for the (11/2⁺) yrast state.



FIG. 8: Beta-gated γ -ray spectrum coincident with the 643.8-keV transition. Prominent γ -peaks that are included in our proposed decay scheme of ¹²⁵Ag are labeled.



FIG. 9: Beta-gated γ -ray spectrum coincident with the 575-keV multiplet. Prominent γ -peaks are labeled. The inset shows the energy of the peak in the singles spectrum (top), in coincidence 606.6-keV γ (middle), and the 655.9-keV γ (bottom), with the fitted peak energy displayed.

3. 1121.1 keV level

Like the above 643 keV level, the 1121.1 keV level is strongly fed in the β decay of the (9/2⁺) ground state in ¹²⁵Ag, with an apparent β feeding of 12.1(11) % with a logft of 5.3(1), suggesting an allowed decay to this state with possible J^{π} of 7/2⁺, 9/2⁺ and 11/2⁺. The state strongly decays to the (9/2⁻) 239 and 11/2⁻ 189-keV states and has weaker decays to the (5/2⁺) 384 and (7⁺) 644-keV states. The spin of 7/2⁺ is excluded as the level doesn't have a branch to the 3/2⁺ ground state (which would be E2). Spin equal to 11/2⁺ is also excluded as the transition to the (5/2⁺) 384 keV state would be M3. We therefore assign this level as (9/2⁺).

4. 873.5, 986.9, 1051.2, and 1164.6 keV levels

The four levels at energies of 873.5, 986.9, 1051.2, and 1164.6 keV have been observed to decay to the $(9/2^{-})$ 239.1-keV level, with no observed decay into the positive parity states at 643.7 $(7/2^+)$, 384.1 $(5/2^+)$ or 0.0 $(3/2^+)$ keV. The apparent logfts to these levels are all > 6, which is consistent with a first forbidden decay from the $(9/2^+)$ ground state of ¹²⁵Ag to negative parity states. Of these states, only the 1164.6 level has been observed to decay to the $11/2^{-}$ 188.5 keV isomer. The transition from the 986.9-keV level to the 188.5 keV level is not observed with a lower limit of < 0.3% (normalized to the 384.1) transition). Low intensity transitions from the 873.5 and 1051.2-keV states of 685.0 and 862.8- keV respectively, to the $11/2^{-}$ isomer would not be observed in the singles spectrum because of contamination from relatively large peaks resulting from the decay of ¹²⁵Cd at 683.6 and 859.7 keV [10]. There is a $(15/2^{-})$ level reported in [13] at x + 719.7(2) keV, where x is the energy of the $11/2^{-}$ isomer. With our value of 188.5(2), this puts the $(15/2^{-})$ level at 931.8(3) keV. Based on the closeness of energies to the $(15/2^{-})$ level, the 873.5- and 986.9-keV levels are possible candidates for $7/2^-$ and $11/2^-$ core coupled levels.

B. Decay of the $(1/2^{-})$ isomer of ¹²⁵Ag

The 352.6-keV γ transition is not connected via $\gamma - \gamma$ coincidence to any transitions we have assigned to the decay of the (9/2⁺) ground state of ¹²⁵Ag (although the decay schemes of the two isomers are connected by the 190.5 keV transition - see 574.8 keV level subsection below). The energy of this γ matches well with the reported 352.7-keV γ assigned to the decay of ^{125m}Ag from Ref. [23]. We therefore assign this transition and those in coincidence with it to the β -decay of the (1/2⁻) isomer. Table 4 lists the transitions assigned to this decay along with their intensities, initial and final states, assigned mutipolarities and coincidence relationships. Transitions

and levels that require further explanation are listed in the following sections by level.

Table 5 lists the levels in ¹²⁵Cd populated in this decay. As noted previously, the calculated log ft values are based on the apparent β feeding and should be considered as lower limits and direct decays to the ground state and isomer are assigned an upper limit based on the J^{π} values of the initial and final levels. The direct β -decay of the (1/2⁻) ground state to the 3/2⁺ ground state of ¹²⁵Cd would be a first forbidden β -decay with a corresponding log ft>6, resulting in a branching ratio of < 4%. A direct β -decay to the 11/2⁻ isomer in ¹²⁵Cd would require a fourth forbidden β -decay and therefore have a branching ratio of ≈ 0 . The resulting decay scheme is shown in Figs 6 and 7.

1. 574.8 keV level

The 574.8 keV level is strongly fed by the β -decay of the $(1/2^-)$ isomer of ¹²⁵Ag with an apparent logft of 5.1, which is consistent with an allowed β -decay. It then decays to the $3/2^+$ ground state, the $(1/2^+)$ 352.6-keV, and $(5/2^+)$ 384.1-keV states. This is suggestive of a $(3/2^-)$ assignment for this level, however, a mechanism for lowering a $(3/2^-)$ to that energy is unknown to the best of our knowledge. However, systematics would suggest an assignment of $(3/2^+)$ to this state resulting from a 1st forbidden decay. It is likely that the actual logft to this level is much higher that the apparent value of 5.1 due to unobserved transitions to the 574.8 level from higher lying states. Note that the statistics from the decay of the $(1/2^-)$ beta decay are much lower than those from decay of the $(9/2^-)$ ground state.

We assign three γ 's in this work that have energies close to 575 keV: 1) 574.5 keV de-exciting the 1180.8level to the 606.6-keV level, 2) 575.4 keV de-exciting the 1615.4 level to the 1040.0-keV level, and 3) 574.8 keV deexciting the 574.8-keV level to the ground state. For the first two transitions, the energy and intensity were determined via coincident γ gates on the 606.6 and 655.8-keV γ 's respectively. A comparison of these two intensity values to the 575 peak in the singles spectra shows that there is an excess of counts which indicates a third peak. The energy of this peak at 574.8 keV and it's intensity were determined from the β -gated γ spectrum after subtracting out the contributions of the first two transitions from the multiplet. The spectrum obtained by gating on the 575-keV multiplet is shown in Fig. 9, with a closeup of the three 574.8, 574.5 and 575.4-keV γ 's that make up the multiplet in the singles spectrum and gated on the 606.6 and 655.9-keV γ 's respectively.

The 190.5 keV transition decays from the 754.8 keV level to the 384.1 keV level and connects the levels fed by the $(1/2^{-})$ isomer with those fed by the $(9/2^{-})$ ground state. In our data there is a significant amount of γ 's from the decay of ¹²⁵Cd including the known 191.88(15)-keV transition [10]. The $\gamma - \gamma$ coincident gate on the 191-keV

TABLE III: Effects of the assignment of possible multipolarites for the 50.5-keV transition in 125 Cd. Column three reflects the total intensity of the transition (with the sum of all observed transitions between levels equal to 100%), *i.e.* it is not normalized to the 384 keV transition.

Mult.	σ_{tot}	I_{tot}	$\%~\beta$ feeding	log ft	J^{π}	β transition
E1	1.14	5.8(1)%	-9.8(9)%		$9/2^{+}$	allowed
M1	3.3	11.1(2)%	2.7(8)%	6.1(2)	$9/2^{-}$	1st forbidden
E2	17.1	34.6(6)%	43.6(5)%	5.0(1)	$7/2^{-}$	1st forbidden
M2	64.2	65.5(11)%	76.8(2)%	4.7(1)	$7/2^{+}$	allowed $(l$ -forbidden)

TABLE IV: List of γ -rays resulting from the beta decay of 125m Ag to levels in 125 Cd. The intensities in columns 2 and 4 are normalized to the 100% 352-keV transition. The I_{tot} intensities in column 4 are the total transition strength (I_{γ} +c.e.). Gamma intensities are based on a) singles, or b) coincidence relationships. The * symbol denotes that the γ was not placed in the level scheme. Contaminants from the β -decay of 125 Cd are marked with the symbol **.

$E_{\gamma} (keV)$	$I_{\gamma}(\%)$	α_{tot}	$I_{tot}(\%)$	E_i	\mathbf{E}_{f}	Mult	coincident γ -rays
190.5(4)	$10.5(9)^{b}$	0.152	11.7(7)	574.8	384.1	(E2)	193*, 334*, 384, 606**, 737**, 1173**
222.0(3)	$7.1(11)^{b}$		7.4(11)	574.7	352.6		353
352.6(1)	100^a	0.018	100	352.6	0.0	(M1, E2)	222, 1023, 1079, 1167, 1725, 1781, 1785
							1888,1940,2047,2805,2825,2969
384.1	$11.7(7)^{b}$	0.0041	11.7(7)	384.1	0.0	[E1]	191
574.8(2)	$19.7(9)^{a}$		19.0(9)	574.7	0.0		
1022.7(3)	$2.8(9)^{b}$		2.7(9)	1375.3	352.6		353
1079.2(4)	$3.2(8)^{b}$		3.1(8)	1431.7	352.6		353
1166.7(4)	$6.1(11)^{b}$		6.0(10)	1519.3	352.6		353
1725.2(4)	$5.6(12)^{b}$		5.5(12)	2077.8	352.6		353
1780.9(4)	$13.9(16)^{b}$		13.7(16)	2133.7	352.6		353
1795.4(3)	$2.9(10)^{b}$		2.8(10)	2148.0	352.6		353
1888.0(3)	$5.3(12)^{b}$		5.2(12)	2240.6	352.6		353
1940.4(3)	$7.4(15)^{b}$		7.2(15)	2293.0	352.6		353
2046.6(5)	$7.2(15)^{b}$		7.1(15)	2399.2	352.6		353
2133.8(3)	$25.7(18)^a$		25.3(18)	2133.7	0.0		
2804.7(4)	$18(2)^{b}$		18(2)	3157.3	352.6		353
2825.0(3)	$4.2(10)^{b}$		4.1(10)	3177.6	352.6		353
2968.7(5)	$6.2(12)^{b}$		6.1(11)	3321.3	352.6		353
3157.5(4)	$3.8(5)^{a}$		3.7(5)	3157.3	0.0		

multiplet is shown in Fig. 10. The prominent peaks in coincidence are the 384.1-keV (from ¹²⁵Ag) and the 736.7 and 1173.2-keV transitions (from ¹²⁵Cd). Gating on the 736.7-keV γ gives a peak at 191.8 keV which agrees well with the literature value [10], and gating on the 384.1-keV peak yields a peak at 190.5 keV (see inset of Fig. 10).

C. Half-lives

The measurements of the half-lives of the ground state and isomer in 125 Ag were done by measuring the apparent half-life of each γ -ray with time equal zero defined as when the tape moved the sample to the front of the detectors. As detailed above, the A= 125 beam was on for two seconds, then turned off for 500 ms to allow measurements of the half-lives of the ground state and isomer in ¹²⁵Ag (and ¹²⁵Cd). This was done by measuring the apparent half-life of the most intense γ -rays of each starting when the beam was tuned off. The measured half-life $(T_{1/2}^{tot})$ is a combination of the β -decay $T_{1/2}^{\beta}$ and the γ -ray decay $T_{1/2}^{\gamma}$, where it was assumed $T_{1/2}^{\gamma} << T_{1/2}^{\beta}$. This is shown in Fig. 11 for the 384.1 keV γ -transition from the β -decay of ¹²⁵Ag and the 436.3 keV γ -transition from the β -decay of ¹²⁵Cd. In this figure, the grow-in (beam on) and decay (beam off) of both nuclei is observed. In this figure, time equal zero corresponds to the tape movement. Note that ¹²⁵Cd is produced both directly and as



FIG. 10: Beta-gated γ -ray spectrum coincident with the 190.5-keV transition. Prominent γ -peaks are labeled. There is significant contamination in this peak from the 191.88(15)-keV γ resulting from the β -decay of ^{125m}Cd. The inset shows the energy of the peak in coincidence with the ¹²⁵Ag 384.1-keV γ (top), and the 736.7-keV γ (bottom) from ¹²⁵Cd. A dashed line was added to the inset to guide the eye.

the β -daughter, making the time spectrum more complicated than for the ¹²⁵Ag activity, which is only produced directly.

For the ground state β -decay of ¹²⁵Ag, the half-life was obtained from the weighted average of the largest gamma peaks in its decay. The γ 's used were the 50.5, 384.1, 643.8, and 606.6-keV transitions with measured $T_{1/2}$ values of 176.1(49), 176.4(50), 173.5(80), and 174(15) respectively, resulting in a weighted average of 176(3) ms. For the isomer of ¹²⁵Ag, the half-life was obtained from the 352.6-keV transition to be 159(21) ms.

D. β -delayed neutron emission

Known γ -rays from levels in ¹²⁴Cd are present in the data with the assignment made by $\gamma - \gamma$ coincidences. The available energy for β -decay into neutron unbound states ($Q_{\beta n}$) is 4110(430) keV [24] out of a total Q_{β} of 8830(430) keV. The lowest energy levels in ¹²⁴Cd above the 0⁺ ground state are the 2⁺ 613.0 keV, 4⁺ 1385.2 keV, 2⁺ 1427.2, 0⁺ 1573.1, and the (5⁻) 1846.5 keV states [33]. Gammas from the 613.0, 1385.2 and 1846.5 keV states have been observed with the expected coincidence relationships. Fig. 10 shows the coincidence spectra gated on the a) 461.3 keV, b) 772.2 keV, and c) 613.0 keV γ -rays. In all three spectra, the expected γ -rays are all observed. In addition, the gate on the 772-keV γ has a peak at

436.5 keV, which results from coincidence with the tail of the unresolved 774.5 keV γ arising from the β -decay of ¹²⁵Cd [10]. The 384 keV peak in the 613.0 keV gated spectrum arises from the tail of the unresolved 615.1 keV γ from ¹²⁵Ag.

We also observed γ 's with energies of 62.2(3), 142.9(3), 179.4(3) and 241.7(4) keV which match well with the known 62.2(1), 143.0(2), 179.6(1) and 242.0(3) keV γ 's [34] arising from the β -decay of ¹²⁴Cd. These are listed in Table 6. Fig. 11 shows the γ coincidence spectrum gated on the 62-keV γ . The expected 179 and 143-keV γ s are observed in the spectrum. The 36.5 keV line is absent as it is below the energy threshold. The inset of the figure shows the decay scheme from Ref. [34].

Neutron emission to states in ¹²⁴Cd are expected to occur predominately from excited levels in ¹²⁵Cd with similar J^{π} values. These neutron-emitting states are populated primarily by allowed (and to a lesser extent first forbidden) β -decays of ¹²⁵Ag. These neutron unbound states will then decay by neutron emission to states in ¹²⁴Cd requiring the smallest change in angular momentum - primarily $\Delta \ell$ equal 0. Therefore, we assign those β -delayed neutrons to the (4⁺) 1385.2 and (5⁻) 1846.5 keV levels in ¹²⁴Cd as arising from the β -delayed neutron emission of the (9/2⁺) ground state and the delayedneutrons populating the 2⁺ 613.0 keV and 0⁺ ¹²⁴Cd ground state as arising from the (1/2⁻) isomer of ¹²⁵Ag. This is detailed in Table 7, showing ℓ =0 delayed neutron



FIG. 11: Timing spectrum of 384.1-keV transition arising from the β -decay of ¹²⁵Ag compared to the 436.3-keV transition from the β -decay of ¹²⁵Cd.

TABLE V: Levels in $^{125}\mathrm{Cd}$ populated by the $\beta\text{-decay}$ of $1/2^{-125m}\mathrm{Ag}.$

Level (keV)	β %	log ft	
0.0	<4	>6	
188.5(2)	0		
352.6(1)	7.6(28)	5.6(3)	
574.8(3)	22.7(9)	5.1(2)	
1375.3(3)	1.6(5)	6.1(2)	
1431.7(4)	1.9(5)	6.0(2)	
1519.3(4)	3.5(6)	5.7(2)	
2077.8(4)	3.3(7)	5.6(2)	
2133.7(4)	23.0(9)	4.7(2)	
2148.0(3)	1.7(6)	5.8(2)	
2240.6(3)	3.1(7)	5.5(2)	
2293.0(3)	4.3(8)	5.4(2)	
2399.2(5)	4.2(9)	5.4(2)	
3157.3(4)	12.7(11)	4.6(2)	
3177.6(3)	2.4(6)	5.4(2)	
3321.3(5)	3.6(7)	5.1(2)	
>4621	4.6(12)		

emission to the 5⁻ and 4⁺ ¹²⁴Cd states from the (9/2⁺) ground state and emission to the 2⁺ and 0⁺ ¹²⁴Cd states from the (1/2⁻) isomeric state. Determining the delayed neutron branching ratio to the excited states in ¹²⁴Cd is a straightforward ratio compared to other observed γ 's

from ¹²⁵Ag. The branching ratio to the ground state of ¹²⁴Cd was determined by subtracting the relative amount of ¹²⁴Cd beta decay determined by the 179 keV γ 's from β -decay of ¹²⁴Cd to the observed 613 keV γ in ¹²⁴Cd from β -delayed neutron decay of ¹²⁵Ag.

This results in a branching ratio for the $(9/2^+)$ ¹²⁵Ag ground state of 1.2(2)% (0.56(5)% to the 1846 keV level and 0.59(10)% to the 1385 keV level), and 4.6(12)% (1.5(2)% to the 613 keV level and 3.1(10)% to the ground state) for the $(1/2^-)$ isomer. The assigned value of the β -n branching to the ¹²⁴Cd ground state should be considered as a lower limit as the beta branching ratios to the 241.8-keV (< 67.7%) and 205.2-keV level (< 15.8%) are considered to be upper limits [34] due to the pandemonium effect [32]. The resulting decay scheme for β -delayed neutron emission from ^{125,125m}Ag is shown in Fig. 13.

A comparison of the β -n branching ratios calculated by Moller *et al.* [35] and Marketin *et al.* [36] versus experimental literature values for the neutron-rich odd-mass is shown in table 7. Our experimental values of 1.2(2)% for the $9/2^{+}$ ¹²⁵Ag and 4.6(10)% for the $1/2^{-}$ ^{125m}Ag are significantly lower than the previously reported value of 11.8(10)% [20] and matches somewhat with the value of 2.2(11)% [21] which is a combination of decays from the ground state and isomer.

In conclusion, we have investigated the β -decay of ¹²⁵Ag. We have found evidence for β decay and β -delayed neutron emission from both the (9/2⁺) ground state and the (1/2⁻) isomer of ¹²⁵Ag. The resulting partial decay scheme of ^{125m}Ag consists of 16 γ 's from 14 levels in the β daughter, while the partial decay scheme of the ground

$E_{\gamma} (keV)$	I_{γ}	α_{tot}	I_{tot}	E_i	E_{f}	Mult	coin. γ -rays
62.2(3)	28(17)	0.664	47(29)	62.3	0.0	${ m E1}$	142.9, 179.3
142.9(3)	10(4)	0.063	11(4)	205.2	62.3	${ m E1}$	62.3
179.4(3)	42(12)	0.033	43(12)	241.6	62.3	${ m E1}$	62.3
241.7(4)	12(3)	0.07	13(3)	241.6	0.0	E2	

TABLE VI: Observed $\gamma\text{-rays}$ in this work from the β decay of $^{124}\text{Cd.}$

TABLE VII: Beta-delayed neutron emission results from 125 Ag and 125m Ag.

nuclide	$J\pi$	allowed β decays	final state (^{124}Cd)	$\Delta \ell$	β -n branching
125 Ag	$(9/2^+)$	$7/2^+, 9/2^+, 11/2^+$	$5^-, 1846 \text{ keV}$	0, 1, 2	0.56(5)%
			$4^+, 1385.2 \text{ keV}$	0, 1	0.59(10)%
			$2^+,613.0~{\rm keV}$	1, 2, 3	0
			0^+ , ground state	3, 4, 5	0
$^{125m}\mathrm{Ag}$	$(1/2^{-})$	$1/2^{-}, 3/2^{-}$	5-, 1846 keV	3, 4	0
			$4^+, 1385.2 \text{ keV}$	2, 3	0
			$2^+,613.0~{\rm keV}$	0, 1	1.5(2)%
			0^+ , ground state	0, 1	3.1(10)%

state of $^{125}\mathrm{Ag}$ consists of 72 γ 's from 47 levels. In both the isomer and ground state, evidence for β -delayed neutron emission was observed, with the resulting branching ratios of 4.6(12)% for the isomer, and 1.2(2)% for the ground state.

Acknowledgments

This work has been supported by the U. S. Department of Energy, Office of Nuclear Physics under contracts DE-AC02-05CH11231, DOE-AC05-00OR22725, DE-FG05-88ER40407, DE-FG02-96ER40983 and the National Nuclear Security Administration (NNSA) under the Stewardship Science Academic Alliances program through U.S. Department of Energy (DOE) Cooperative Agreements No. DE- FG52-08NA28552 and DE-NA0002132.

TABLE VIII: Comparison of experimental and theoretical β -delayed neutron probabilities for the odd-mass neutron-rich Cd isotopes.

Nuclide	theory [35]	theory [36]	Lit. value (exp)	this work
¹²⁰ Ag	0	0.5%	<0.003% [37]	
^{121}Ag	0	0.7%		
^{122}Ag	0	0.6%	0.186(10)% [37]	
^{123}Ag	7%	0.9%	0.55(9)% [37]	
^{124}Ag	10%	0.7%	1.3(9)% [38]	
$^{125}\mathrm{Ag}$	4%	1.1%	11.8(10)% [20], $2.2(11)%$ [21]	$4.6(12)\% (^{125m} \text{Ag})$
				1.2(2)% (¹²⁵ Ag)
^{126}Ag	4%	1.0%	13.7(11)% [20]	
$^{127}\mathrm{Ag}$	7%	1.9%	14.6(15)% [20]	
$^{128}\mathrm{Ag}$	9%	1.8%	20(5)% [20]	
$^{129}\mathrm{Ag}$	10%	13.3%	$<\!\!20\%$ [20]	



FIG. 12: Gamma spectra gated on the known γ -rays in ¹²⁴Cd arising from the β -delayed neutron emmission of ¹²⁵Ag.



FIG. 13: Beta-gated γ -ray spectrum coincident with the 62.2-keV γ -ray arising from the β decay of ¹²⁴Cd. The inset of the figure shows the decay scheme from Ref. [34].



FIG. 14: Proposed partial decay scheme for β -delayed neutron emission from 125,125m Ag.

- T. Kracikova, I. Prochazka, Z. Hons, M. Fiser, A. Kuklik, Czech. J. Phys. B27, 1099 (1977).
- [2] B. Fogelberg, Z. Ye, B. Ekstrom, E. Lund, K. Aleklett, L. Sihver, Z. Phys. A337, 251 (1990).
- [3] W. Bruchle, G. Herrmann, Radiochim. Acta 30, 1 (1982).
- [4] B. Fogelberg, Y. Kawase, J. McDonald, A. Backlin, Nucl. Phys. A267, 317 (1976).
- [5] B. Fogelberg, P. Hoff, Nucl. Phys. A**391**, 445 (1982).
- [6] H. Huck, A. Jech, G. Marti, M. L. Perez, J. J. Rossi, H. M. Sofia, Phys.Rev. C40, 1384 (1989).
- [7] Y. Kawase, B. Fogelberg, J. McDonald, A. Bäcklin, Nucl. Phys. A 241, 237 (1975).
- [8] G. H. Fuller, J. Phys. Chem. Ref. Data 5, 835 (1976).
- [9] H. Penttila, J. Aysto, K. Eskola, Z. Janas, P. P. Jauho, A. Jokinen, M. E. Leino, J. M. Parmonen, P. Taskinen, Z. Phys. A338, 291 (1991).
- [10] H. Huck, A. Jech, G. Marti, M. L. Perez, J. J. Rossi, H. M. Sofia, Phys.Rev. C39, 997 (1989).
- [11] A. Kankainen, et al., Phys.Rev. C 87, 024307 (2013).
- [12] D. Lascar, et al., Phys. Rev. C 96, 044323 (2017).
- [13] N. Hoteling, et al., Phys. Rev. C 76, 044324 (2007).
- [14] Ch. Lorenz, et al., Phys.Rev. C 99, 044310 (2019).
- [15] J. Taprogge et al., Phys.Lett. B 738, 223 (2014).
- [16] A. Scherillo, J. Genevey, J. A. Pinston, A. Covello, H. Faust, A. Gargano, R. Orlandi, G. S. Simpson, I. Tsekhanovich, N. Warr, Phys. Rev. C 70, 054318 (2004).
- [17] G. S. Simpson, A. Scherillo, J. Genevey, R. Orlandi, J. A. Pinston, I. S.Tsekhanovich, N. Warr, A. Covello, A.Gargano, J. Phys. Conf. Ser. 267, 012031 (2011).
- [18] M. Rejmund, A. Navin, S. Bhattacharyya, M. Caamano, E. Clement, O. Delaune, F. Farget, G. de France, B. Jacquot, A. Lemasson, Phys. Rev. C 93, 024312 (2016).
- [19] V. N. Fedoseyev, et al., Z. Phys. A**353**, 9 (1995).
- [20] K. I. Smith Thesis, Univ. of Notre Dame (2014) (unpub-

lished).

- [21] O. Hall, et al., Phys. Lett. B 816, 136266 (2021).
- [22] G. Lorusso, et al., Phys. Rev. Lett. 114, 192501 (2015).
- [23] Z. Q. Chen, et al., Phys. Rev. Lett. 122, 212502 (2019).
- [24] W. J. Huang, M. Wang, F. G. Kondev, G. Audi, S. Naimi, Chin. Phys. C 45, 030002 (2021).
- [25] B. E. Tomlin, P. F. Mantica, W. B. Walters, Eur. Phys. J. Spec. Top. **150**, 183 (2007).
- [26] D. W. Stracener, in Proceedings of the Sixteenth International Conference on the Application of Accelerators in Research and Industry, edited by J. L. Duggan and I. L. Morgan, CP576, AIP Press, New York (2000) pp. 257-260.
- [27] J. R. Beene, et al., J. Phys. G 38, 024002 (2011).
- [28] W. A. Peters, et al., Nucl. Instr. and Meth. A 836, 122 (2016).
- [29] Y. Zhao, unpublished, Ph. D. Thesis, Tokyo Metropolitan University, (1996).
- [30] http://www.xia.com/
- [31] T. Kibédi, T. W. Burrows, M. B. Trzhaskovskaya, P. M. Davidson, C. W. Nestor, Jr., Nucl. Instr. and Meth. A 589, 202 (2008).
- [32] J. C. Hardy L. C. Carraz, B. Jonson, P. G. Hansen, Phys. Lett. B 71, 307 (1977).
- [33] J. C. Batchelder, et al., Phys. Rev. C 89, 054321 (2014).
- [34] J. C. Batchelder, et al., Phys. Rev. C 94, 024317 (2016).
- [35] P. Möller, M. R. Mumpower, T. Kawano, W. D. Myers, At. Data Nucl. Data Tables 125, 1 (2019).
- [36] T. Marketin *et al.*, Phys. Rev. C **93**, 025805 (2016).
- [37] P. L. Reeder, R. A. Warner, R. L. Gill, Phys. Rev. C 27, 3002 (1983).
- [38] F. Montes, et al., Phys. Rev. C 73, 035801 (2006).