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# Insulating titanium oxynitride for visible light photocatalysis

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We propose insulating titanium oxynitrides  $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$  as promising water-splitting photocatalytic materials in the visible-light range. Using first-principles many-body perturbation theory based on the GW approximation, we show that corundum-type  $\text{Ti}_2\text{N}_2\text{O}$  (an example  $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$  compound with  $n = 2$ ) has a smaller band gap of about 2.5 eV, which is more suitable to absorb visible light, compared to other Ti-based oxides such as  $\text{TiO}_2$  and  $\text{SrTiO}_3$  with a band gap of more than 3 eV. Band gap reduction in  $\text{Ti}_2\text{N}_2\text{O}$  is caused by an upward shift of the valence band (negative shift to the oxidation potential of  $\text{H}_2\text{O}$  to  $\text{O}_2$ ) due to the presence of nitrogen 2p states. The conduction band is dominated by Ti 3d states and the conduction band minimum is nearly unchanged. As a result, the band edge potentials of  $\text{Ti}_2\text{N}_2\text{O}$  are better aligned to the water reduction and oxidation levels. Our theoretical predictions provide useful insights for the discovery of new visible-light-driven photocatalyst for water splitting.

## I. INTRODUCTION

Titanium dioxide ( $\text{TiO}_2$ ) is extensively studied as a functional material with the potential of being used in various technological applications such as photocatalysis, photovoltaics, and oxide electronics.<sup>1–3</sup> In particular, triggered by the demonstration of photocatalytic system for water splitting using rutile  $\text{TiO}_2$  anode and Pt cathode,<sup>4</sup> numerous efforts have been made to discover  $\text{TiO}_2$ -based materials with better photocatalytic performance.<sup>5</sup> Key properties required for water-splitting photocatalytic materials include the magnitude of the band gap and the position of band edges relative to the water reduction and oxidation levels. To ensure the water-splitting reaction without bias voltage, the conduction band minimum (CBM) needs to be negative (higher in energy) than the  $\text{H}_2/\text{H}_2\text{O}$  level of water [usually set to be the normal hydrogen electrode (NHE)], and the valence band maximum (VBM) needs to be positive (lower in energy) than the  $\text{O}_2/\text{H}_2\text{O}$  level (1.23 V versus NHE). A band gap energy in the visible-light range is desirable for utilizing the solar spectrum to induce the water-splitting reaction.

Positions of the band edges of  $\text{TiO}_2$  are suitable for water reduction and oxidation, while the band gap of  $\text{TiO}_2$  (3.0–3.6 eV in rutile<sup>6–9</sup> and 3.2 eV in anatase<sup>10,11</sup>) is too large to absorb visible light, resulting in poor solar

absorption efficiency. To improve photocatalytic activity of  $\text{TiO}_2$  under visible light, there have been continuing efforts to form a band gap suitable for solar-energy absorption as well as the water-splitting reaction. In particular, nitrogen doping into  $\text{TiO}_2$  was shown to narrow its band gap, forming an absorption tail in the visible-light region.<sup>12–14</sup> Band gap narrowing in N-doped  $\text{TiO}_2$

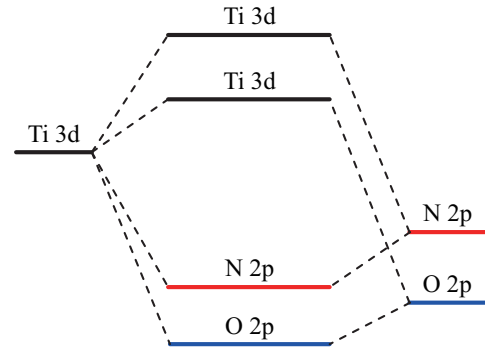


FIG. 1: Schematic band structure of an insulating titanium oxynitride ( $\text{Ti-N-O}$  compound). Valence band consists mainly of hybridized orbitals from N 2p and O 2p states, whereas conduction band is dominated by empty Ti 3d orbitals. Octahedral coordination around Ti atoms splits the 3d band into lower and upper levels.

can be explained by the schematic band structure shown in Fig. 1. Valence band is composed of hybridized N  $2p$  and O  $2p$  orbitals. Since the N  $2p$  states are higher in energy than O  $2p$  states, VBM is expected to be shifted higher compared to pristine  $\text{TiO}_2$ . Conduction band is dominated by empty Ti  $3d$  orbitals, being less affected by the dopant. Nitrogen doping is a useful clue to the band-gap engineering of Ti oxides. There is, however, a doping limit in  $\text{TiO}_2$ , making it difficult to optimize the photocatalytic performance under visible light. Several studies suggest that the amount of oxygen vacancy in Ti oxide plays a crucial role in controlling nitrogen concentration.<sup>15,16</sup>

Here, we study an alternative way to achieve much higher nitrogen concentrations in Ti oxides. We focus on Ti oxides with a composition of  $\text{Ti}_n\text{O}_{2n-1}$ , which can be viewed as oxygen-deficient variants of  $\text{TiO}_2$ . Substituting oxygens with nitrogens can yield a series of Ti oxynitrides ( $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$ ). By performing first-principles calculations based on the density-functional theory (DFT)<sup>17,18</sup> and the GW approximation,<sup>19</sup> we show that corundum-type  $\text{Ti}_2\text{N}_2\text{O}$  (an example  $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$  compound with  $n = 2$ ) has a smaller band gap compared to pristine  $\text{TiO}_2$  phases. We also show that the band edge potentials of  $\text{Ti}_2\text{N}_2\text{O}$  are much better aligned to the water splitting levels.

## II. COMPUTATIONAL METHODS

The ground-state total energies, the Kohn-Sham eigenfunctions, and corresponding eigenenergies are calculated in the framework of DFT<sup>17,18</sup> as implemented in the QUANTUM ESPRESSO package.<sup>20</sup> We use the generalized gradient approximation (GGA) with the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional.<sup>21</sup> We employ norm-conserving pseudopotentials.<sup>22</sup> For Ti pseudopotential, we include semicore  $3s$  and  $3p$  states as valence states. In generating pseudopotentials, core radii (in a.u.) of  $2s$ ,  $2p$ ,  $3s$ ,  $3p$ , and  $3d$  states are set to be 1.05, 1.05, 0.90, 0.90, and 1.00, respectively. Wavefunctions are expanded by a plane wave basis set with a cutoff energy of 160 Ry in order to get a DFT band gap converged to better than 10 meV. Brillouin-zone integration is performed using the Monkhorst-Pack method<sup>23</sup> with  $5 \times 5 \times 8$  and  $5 \times 5 \times 5$  grids for rutile  $\text{TiO}_2$  and anatase  $\text{TiO}_2$ , respectively. For metallic corundum  $\text{Ti}_2\text{O}_3$ , we use a  $5 \times 5 \times 5$   $k$ -point grid and a gaussian smearing with a 0.01 Ry spreading. For corundum-type  $\text{Ti}_2\text{N}_2\text{O}$ , a  $5 \times 5 \times 5$  ( $5 \times 5 \times 2$ )  $k$ -point grid is used for calculations with primitive unit cell (hexagonal conventional cell). Atomic positions are optimized by using the BFGS algorithm.<sup>24-26</sup>

Quasiparticle band structures are calculated by the “one-shot” GW approach combined with the generalized plasmon-pole model<sup>19,27</sup> as implemented in the BERKELEYGW code.<sup>28</sup> Our GW calculations are performed on top of the eigenvalues and eigenfunctions obtained from

the above-mentioned DFT calculations. We employ a fine grid for the dielectric screening to achieve convergence of a band gap energy within 0.01 eV, a  $9 \times 9 \times 15$  grid for rutile  $\text{TiO}_2$ , an  $8 \times 8 \times 8$  grid for anatase  $\text{TiO}_2$ , and a  $4 \times 4 \times 4$  grid for corundum-type  $\text{Ti}_2\text{N}_2\text{O}$ . We adopt simple approximate physical orbitals (SAPO)<sup>29</sup> to get a large number of empty states that are required to compute the polarizability and the self-energy. For rutile  $\text{TiO}_2$  and anatase  $\text{TiO}_2$ , the polarizability matrices are computed using a total of 2000 bands and a screened Coulomb cutoff energy of 20 Ry, and the self-energy corrections are calculated using 6000 bands. For corundum-type  $\text{Ti}_2\text{N}_2\text{O}$ , we employ 2000 bands for sums in the polarizability matrices and 4000 bands for sums in the self-energy operators. Summations over empty states in the self-energy corrections are done with a static remainder approach.<sup>30</sup>

## III. CRYSTAL STRUCTURE

We focus on  $\text{Ti}_n\text{O}_{2n-1}$  compounds as starting materials for insulating Ti oxynitrides without introducing a localized state in the energy gap. Figure 2 illustrates the crystal structure of corundum  $\text{Ti}_2\text{O}_3$ ,<sup>31</sup> a  $\text{Ti}_n\text{O}_{2n-1}$  compound with  $n = 2$ . Ti ions occupy two thirds of the octahedral interstitial sites, whereas oxygen ions nearly form a hexagonal close-packed structure. In this Ti sesquioxide, each Ti atom is coordinated octahedrally by neighboring six O atoms.  $\text{Ti}_n\text{O}_{2n-1}$  compounds with  $n \geq 3$  are known as Magnéli phases,<sup>32,33</sup> which are characterized by the presence of a “shear plane,” an ordered plane

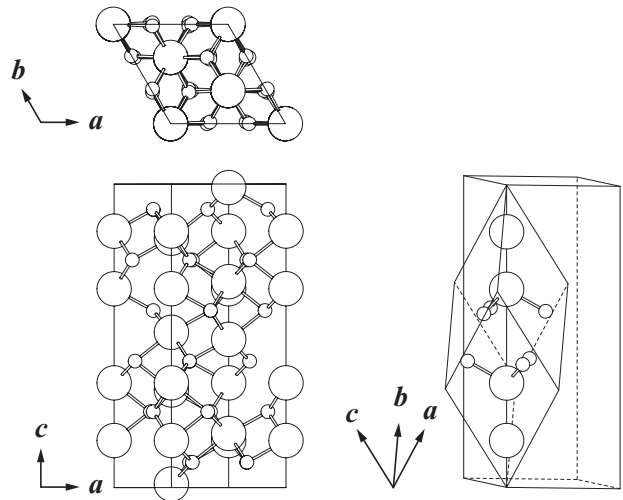


FIG. 2: Crystal structure of corundum  $\text{Ti}_2\text{O}_3$ . Large and small spheres represent Ti and O atoms, respectively. Hexagonal conventional cell (left) contains six formula units, whereas rhombohedral primitive cell (right) has two formula units in it. Each Ti atom is surrounded by six O atoms, forming  $\text{TiO}_6$  octahedron.

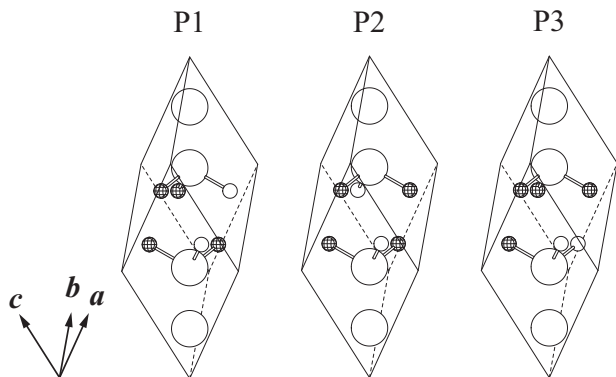


FIG. 3: Crystal structures of corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  with different nitrogen-substitution patterns. Large spheres represent Ti atoms. Small spheres with and without a pattern represent N and O atoms, respectively. The structure labeled “P2” has inversion symmetry and the other two structures do not.

of oxygen vacancies. Most of the  $\text{Ti}_n\text{O}_{2n-1}$  compounds are metallic due to partially occupied Ti 3d states, while some of the low-temperature phases are found to be semi-conducting with a relatively small band gap ( $\sim 0.1$  eV).<sup>34</sup>

Nitrogen doping into  $\text{Ti}_n\text{O}_{2n-1}$  phases can be an alternative way to design new Ti–N–O compounds. For example, by replacing two out of  $2n - 1$  oxygen atoms of  $\text{Ti}_n\text{O}_{2n-1}$  compounds with nitrogen, we can construct a series of Ti oxynitrides  $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$  that are expected to have a large insulating gap. With this composition, the resultant materials can keep the stable +IV valency of titanium.

We carry out a case study on compounds with  $n = 2$  to demonstrate the feasibility of our concept for new Ti oxynitrides. The primitive unit cell of corundum  $\text{Ti}_2\text{O}_3$  has two formula units (4 Ti atoms and 6 O atoms) in it. Suppose that four out of six oxygens are substituted with nitrogen to construct a  $\text{Ti}_2\text{N}_2\text{O}$  compound (two  $\text{Ti}_2\text{N}_2\text{O}$  formula units in an unit cell). There are 15 [=  $6!/(4!2!)$ ] possibilities for nitrogen substitution patterns. We actually have only three kinds due to the symmetries of the parent material as shown in Fig. 3. We call them P1, P2, and P3 (P for primitive). We also construct additional three  $\text{Ti}_2\text{N}_2\text{O}$  structures using the hexagonal conventional cell that contains six formula units. We call them S1, S2, and S3 (S for supercell). The purpose of constructing S structures is to examine the effect of the disorder (different arrangements of N atoms) on the structural stability of corundum-based  $\text{Ti}_2\text{N}_2\text{O}$ .

We evaluate the formation energy per formula unit of a  $\text{Ti}_2\text{N}_2\text{O}$  structure, defined as

$$E = E(\text{Ti}_2\text{N}_2\text{O}) - E(\text{Ti}_2\text{O}_3) - E(\text{N}_2) + E(\text{O}_2). \quad (1)$$

Here,  $E(\text{Ti}_2\text{N}_2\text{O})$  and  $E(\text{Ti}_2\text{O}_3)$  are the total energies per formula unit of  $\text{Ti}_2\text{N}_2\text{O}$  and  $\text{Ti}_2\text{O}_3$  phases, and  $E(\text{N}_2)$  and  $E(\text{O}_2)$  are the total energies of  $\text{N}_2$  and  $\text{O}_2$  molecules, respectively. As shown in Fig. 4, the formation energies of

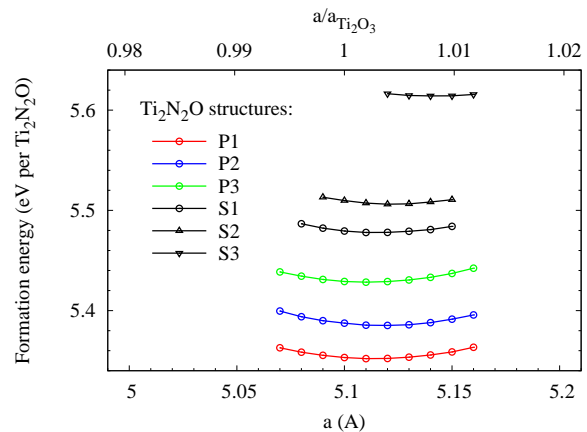


FIG. 4: Formation energies [Eq. (1)] of corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  structures (P1–P3 and S1–S3) as a function of lattice constant  $a$ . Lattice constant  $c$  is optimized at a given  $a$ . Atomic positions are relaxed as well.

P structures are less than those of S structures. We find that this trend is well correlated with the coordination of titanium atoms with neighboring nitrogen and oxygen atoms. In the P structures, each Ti atom is surrounded by four N atoms and two O atoms, forming  $\text{TiN}_4\text{O}_2$  octahedron network. The three S structures have two types of octahedral coordination around Ti atoms, containing  $\text{TiN}_3\text{O}_3$  and  $\text{TiN}_6$  octahedra. Our total-energy calculations on bulk  $\text{Ti}_2\text{N}_2\text{O}$  phases, together with an experimental report<sup>38</sup> on the existence of molecules with a composition  $\text{Ti}_2\text{N}_2\text{O}$ , provide helpful clues to the synthesis and the stabilization of  $\text{Ti}_2\text{N}_2\text{O}$  structures as well as higher  $\text{Ti}_n\text{N}_2\text{O}_{2n-3}$  phases. Table I summarizes structural parameters of corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  structures determined by the present DFT-GGA calculations. In Table I, we also list experimental lattice constant of pristine Ti oxides including corundum  $\text{Ti}_2\text{O}_3$ .

It has come to our attention that nitrogen doping into  $\text{Ti}_n\text{O}_{2n-1}$  compounds with  $n \geq 3$  has been studied theoretically<sup>39</sup> and experimentally.<sup>40–42</sup> In particular, Wu and his co-workers used a high throughput screening method within DFT.<sup>39</sup> They find that  $\text{Ti}_3\text{N}_2\text{O}_3$  (the  $n = 3$  case) adopts the crystal structure of  $\text{Ta}_3\text{N}_5$ , which is almost isomorphic to the  $\alpha$ - $\text{Ti}_3\text{O}_5$  structure. The band gap value of  $\text{Ti}_3\text{N}_2\text{O}_3$  is estimated to be 2.37 eV within their  $\Delta$ -sol methods.<sup>43</sup> The material is also predicted to have a favorable band-edge position for water splitting.

#### IV. ELECTRONIC STRUCTURE

Electronic band structures of corundum  $\text{Ti}_2\text{O}_3$  and corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  phases, calculated by the DFT-GGA and GW methods, are shown in Fig. 5. Since the energies of P structures are less than those of S structures (Fig. 4), GW calculations were performed only on P structures. Corundum  $\text{Ti}_2\text{O}_3$ , a parent phase of  $\text{Ti}_2\text{N}_2\text{O}$

TABLE I: Structural parameters of rutile  $\text{TiO}_2$  [space group  $P4_2/mmm$  (No. 136), tetragonal], anatase  $\text{TiO}_2$  [ $I4_1/amd$  (No. 141), tetragonal], corundum  $\text{Ti}_2\text{O}_3$  [ $R\bar{3}c$  (No. 167), trigonal], and corundum-based  $\text{Ti}_2\text{N}_2\text{O}$ . Atomic positions of  $\text{Ti}_2\text{N}_2\text{O}$  structures are given in the Supplemental Material.<sup>35</sup>

Material	Method	Lattice constant ( $\text{\AA}$ )	Atomic positions	
rutile $\text{TiO}_2$	X-ray [36]	$a = 4.594, c = 2.959$	Ti	$2a$ (0, 0, 0)
			O	$4f$ ( $v, v, 0$ ) <span style="float: right;"><math>v = 0.305</math></span>
anatase $\text{TiO}_2$	X-ray [37]	$a = 3.785, c = 9.514$	Ti	$4a$ (0, 0, 0)
			O	$8e$ (0, 0, $v$ ) <span style="float: right;"><math>v = 0.208</math></span>
corundum $\text{Ti}_2\text{O}_3$	X-ray [31]	$a = 5.149, c = 13.642$	Ti	$12c$ (0, 0, $\pm u$ ), (0, 0, $1/2 \pm u$ ) <span style="float: right;"><math>u = 0.345</math></span>
			O	$18e$ ( $\pm v, 0, \pm 1/4$ ), ( $\pm v, \pm v, \pm 3/4$ ) <span style="float: right;"><math>v = 0.317</math></span>
corundum $\text{Ti}_2\text{O}_3$	GGA (This work)	$a = 5.10, c = 14.0$	$u = 0.345, v = 0.312$	
$\text{Ti}_2\text{N}_2\text{O}$ , P1	GGA (This work)	$a = 5.11, c = 14.1$		
$\text{Ti}_2\text{N}_2\text{O}$ , P2	GGA (This work)	$a = 5.12, c = 13.9$		
$\text{Ti}_2\text{N}_2\text{O}$ , P3	GGA (This work)	$a = 5.11, c = 14.0$		
$\text{Ti}_2\text{N}_2\text{O}$ , S1	GGA (This work)	$a = 5.11, c = 14.0$		
$\text{Ti}_2\text{N}_2\text{O}$ , S2	GGA (This work)	$a = 5.12, c = 14.0$		
$\text{Ti}_2\text{N}_2\text{O}$ , S3	GGA (This work)	$a = 5.14, c = 14.1$		

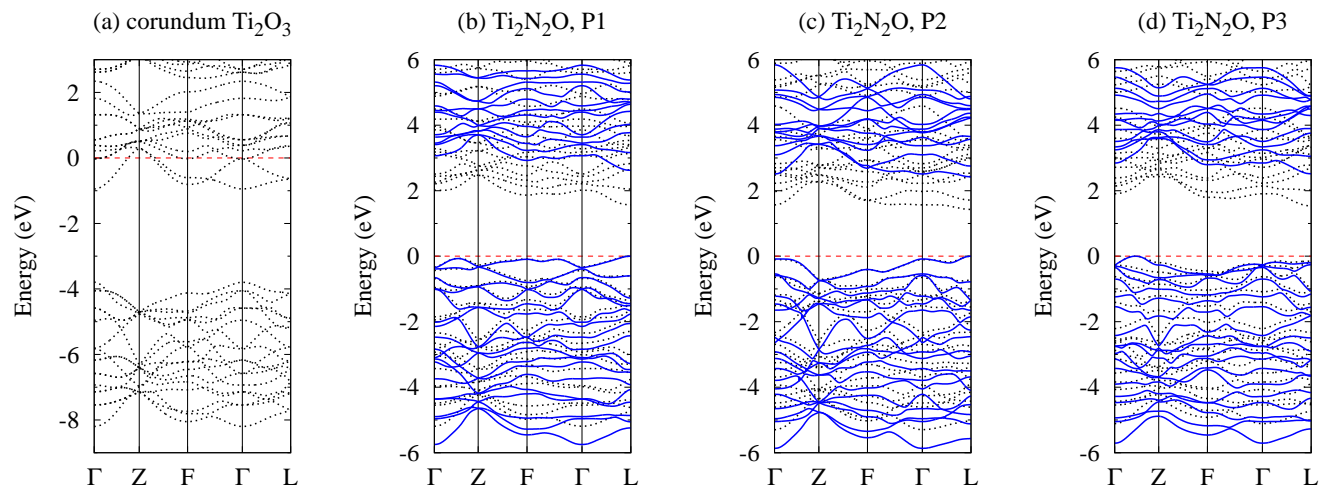


FIG. 5: Band structures of (a) corundum  $\text{Ti}_2\text{O}_3$  and (b)–(d) corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  phases, calculated by GGA (black, dotted) and GW methods (blue, solid). Red dashed line indicates the Fermi level in (a) and the top of the valence band in (b)–(d). P1 and P2 structures are direct-gap semiconductors, whereas P3 structure is an indirect-gap semiconductor. In (b)–(d), the GGA and GW results are displayed with the top of the valence band aligned.

structures, is metallic, being consistent with experimental findings. As shown in Fig. 5(a), the O  $2p$  band (fully occupied) and lower Ti  $3d$  band (partially occupied) in  $\text{Ti}_2\text{O}_3$  are separated by an energy gap of 2.85 eV within the DFT. Nitrogen substitution into  $\text{Ti}_2\text{O}_3$  opens a substantial band gap in the resultant  $\text{Ti}_2\text{N}_2\text{O}$  phases [Figs. 5(b)–(d)] as expected from the electron-counting rule. We find that P1 and P2 structures are direct-gap semiconductors, while P3 structure is an indirect-gap semiconductor. Three P structures have no localized state in the energy gap. Through the projected density

of states calculations, we have confirmed that the valence band of  $\text{Ti}_2\text{N}_2\text{O}$  is composed mainly of N  $2p$  and O  $2p$  states. Major contributions to the conduction band of  $\text{Ti}_2\text{N}_2\text{O}$  are Ti  $3d$  states. These results support the realization of our band-engineering concept shown schematically in Fig. 1.

Band gap values of pristine  $\text{TiO}_2$  phases (rutile and anatase) and corundum-based  $\text{Ti}_2\text{N}_2\text{O}$  structures are summarized in Table II. For calculations on rutile and anatase, we used the experimental lattice constant listed in Table I. The present GW methods give band gap val-

TABLE II: Band gap values in eV, calculated by DFT-GGA and GW methods. Experimental values, determined by photoemission spectroscopy (PES) and optical methods (Opt.), are also listed for comparison.

	TiO <sub>2</sub>		Ti <sub>2</sub> N <sub>2</sub> O		
	rutile	anatase	P1	P2	P3
GGA	1.90	2.20 (2.62 <sup>†</sup> )	1.57	1.43	1.54
GW	3.38	3.86 (4.32 <sup>†</sup> )	2.62	2.42	2.52
GW [44]	3.59	3.83 (4.29 <sup>†</sup> )			
GW [45]	3.34	3.56			
GW [46]	3.46	3.73			
GW [47]	3.40	3.70			
GW [48]	3.23				
GW [49]	3.13				
PES [9]	3.6 ± 0.2				
Opt. [6,11]	3.05	3.2			

<sup>†</sup> Direct gap at the  $\Gamma$  point.

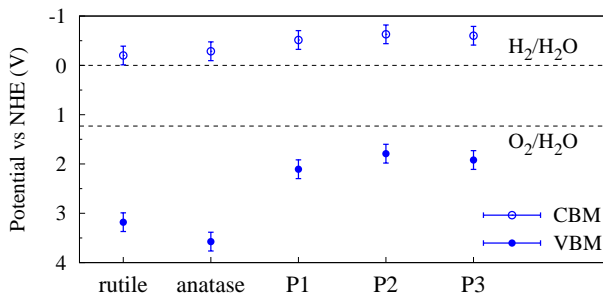


FIG. 6: Estimated GW CBM and VBM positions relative to the water reduction and oxidation levels for pristine TiO<sub>2</sub> phases (rutile and anatase) and corundum-based Ti<sub>2</sub>N<sub>2</sub>O structures (P1–P3).

ues of 3.38 eV for rutile TiO<sub>2</sub> and 3.86 eV for anatase TiO<sub>2</sub>, respectively. Our GW band-gap values agree well with previous GW calculations<sup>44–49</sup> and are consistent with experiments.<sup>6,9</sup> This indicates that our GW methods can give a band gap with high accuracy. The calculated band gap values of corundum-based Ti<sub>2</sub>N<sub>2</sub>O structures, ranging from 2.42 to 2.62 eV within the GW approximation, are significantly smaller than those of pristine TiO<sub>2</sub> phases.

We estimate the positions of CBM and VBM from the GW calculations relative to the water reduction and oxidation levels using the method proposed in Ref. 50. In this method, the average Hartree potential is used as an energy reference for band alignment. For the sake of simplification in calculations for semiconductor-water interfacial systems, this method typically gives the band edge positions with a mean absolute error of 0.19 eV<sup>50</sup>

compared to the experimental data (indicated by an error bar in Fig. 6). Figure 6 shows the band edge positions estimated for pristine TiO<sub>2</sub> phases (rutile and anatase) and Ti<sub>2</sub>N<sub>2</sub>O structures. In rutile and anatase, the CBM locates in a nearly optimal way, while the position of the VBM is rather “deep” in comparison to the O<sub>2</sub>/H<sub>2</sub>O level. In Ti<sub>2</sub>N<sub>2</sub>O structures, the presence of occupied N 2*p* states brings an upward shift of the VBM, making the VBM closer to the O<sub>2</sub>/H<sub>2</sub>O level. The positions of the CBM of Ti<sub>2</sub>N<sub>2</sub>O structures are still within the range of –1 V to the H<sub>2</sub>/H<sub>2</sub>O level. The estimated band edge positions of Ti<sub>2</sub>N<sub>2</sub>O are similar to those of Ti<sub>3</sub>N<sub>2</sub>O<sub>3</sub> (the  $n = 3$  case).<sup>39</sup> The CBM and VBM of Ti<sub>*n*</sub>N<sub>2</sub>O<sub>2*n*–3</sub> compounds are predicted to match better with the water splitting levels.

## V. SUMMARY

We have studied Ti oxynitrides with a composition of Ti<sub>*n*</sub>N<sub>2</sub>O<sub>2*n*–3</sub> that have much higher nitrogen concentrations compared to conventional N-doped Ti oxides. We demonstrate, through first-principles calculations, that Ti<sub>*n*</sub>N<sub>2</sub>O<sub>2*n*–3</sub> compounds can be promising materials to achieve better photocatalytic performance for splitting water. In particular, corundum-type Ti<sub>2</sub>N<sub>2</sub>O, an example Ti<sub>*n*</sub>N<sub>2</sub>O<sub>2*n*–3</sub> compound, is found to have a small band gap that is suitable to absorb visible light. The band edge potentials of Ti<sub>2</sub>N<sub>2</sub>O are much better aligned to the water reduction and oxidation levels. We anticipate the synthesis of Ti<sub>*n*</sub>N<sub>2</sub>O<sub>2*n*–3</sub> structures and their applications for visible-light-driven water-splitting photocatalytic systems.

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