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Valence transitions in the heavy fermion compound YbCuAl as a function of temperature and pressure

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We report on direct measurements of the Yb valence in the heavy Fermion compound YbCuAl as a function of temperature and pressure using resonant x-ray emission spectroscopy. The increase of the Yb\(^{3+}\) component at \(T < 100 \text{ K}\) and ambient pressure, well described by the single impurity Anderson model, is found to compensate for the thermal contraction of the unit cell volume. Under pressure, the Yb valence increases continuously up to 25 GPa, albeit a marked leveling off close to the critical pressure, at \(P \geq 13 \text{ GPa}\). This finding is reminiscent of a recent report on YbCuSi\(_2\) and further confirms the interplay between electronic and magnetic fluctuations near the magnetic instability point of rare-earth intermediate-valence materials.

I. INTRODUCTION

Quantum criticality in 4\(f\)-based heavy fermion systems is a central problem in condensed-matter physics, because it can cause remarkable emergent properties such as non-Fermi-liquid behavior and unconventional superconductivity. Quantum criticality arises from the interplay between the Ruderman-Kittel-Kasuya-Yosida (RKKY) interaction between the 4\(f\) electrons mediated by hybridized conduction electrons and the Kondo screening of the 4\(f\) moments by the conduction electrons.\(^1\)–\(^3\) The Kondo temperature, or the \(c\)-\(f\) hybridization strength, can be affected by an external parameter such as external pressure, allowing one to probe the behavior of a system in the vicinity of a quantum critical point (QCP).

The symmetry in the pressure dependence of the \(c\)-\(f\) hybridization strength between Ce and Yb-based materials arising from the electron-hole symmetry of the Ce and Yb \(4f\) shell is of particular interest. The electron-hole symmetry results in admirably symmetric pressure-dependent physical properties, such as pressure-induced transitions from antiferromagnetic ordering to mixed valence via the QCP for Ce systems and the reverse transitions for Yb systems. Recently, however, the possibility of unsymmetrical behaviors for Ce and Yb systems has been suggested–Goltsev et al.\(^4\) pointed out that some Yb compounds can undergo a reentrance transition to a mixed-valent state at high pressure that is beyond the pressure-induced magnetic ordered phase for the trivalent state. YbCu\(_2\)Si\(_2\) has been considered as a candidate compound for such an anomaly because its maximum temperature of the resistivity, \(T_{\text{max}}\), shows a minimum at \(P = 15 \text{ GPa}\), in which \(T_{\text{max}}\) increases with pressure. This behavior was attributed to the increase of the Kondo temperature (\(T_K\)) as a function of pressure due to the enhancement of \(c\)-\(f\) hybridization. However, direct measurements of the Yb valence in YbCuSi\(_2\) have been recently conducted at low temperature and high pressure. These results showed a characteristic kink around the QCP in the slope of Yb valence versus pressure, but no indication of a reentrant valence transition was observed up to 20 GPa.\(^5\)

Here we examined the possibility of a valence reentrance at high pressures using the candidate, archetypal heavy fermion, YbCuAl. In an early report by Mignot and Wittig,\(^6\) high-pressure resistivity of the system at room temperature was shown to exhibit a minimum around 8 GPa, followed by a gradual increase up to 20 GPa. If this is an indication of an increase in the Kondo scattering, i. e. of \(T_K\), pressure-dependent measurements for the strength of the \(c\)-\(f\) hybridization or the valence are highly desired. So far, conflicting information about the pressure dependence of the Yb valence in YbCuAl has been reported. The quadrupole coupling in Mössbauer spectroscopy indicated that a valence transition towards Yb\(^{3+}\) was completed at approximately 5 GPa,\(^7\) whereas signs of magnetic ordering transition to the trivalent state were suggested using resistivity results at 1 K and 8 GPa.\(^8\) This mismatch between the electronic and magnetic transitions remains unsolved, and further calls for a reliable study of the Yb \(4f\) electronic occupancy as a function of pressure.

In this paper, we present experimental results on the Yb valence in YbCuAl as a function of temperature as low as 8 K and pressure up to 25 GPa by employing bulk-sensitive techniques to probe its electronic structures, including partial fluorescence yield x-ray absorption spectroscopy (PFY-XAS) and resonant x-ray emis-
sion spectroscopy (RXES). We have observed that the valence increases steadily with pressure until around 13 GPa, and then levels off at higher pressures. This change in the slope in the valence manifests a possible change in the magnetic fluctuations, indicating the occurrence of a magnetic instability point. While the valence is still slightly below 3 at 25 GPa, Yb adopts the magnetic trivalent state when the valence reaches ~2.95 at 8 GPa. However, no evidence of a reentrant valence transition was found up to 25 GPa.

II. EXPERIMENTS

A polycrystalline sample of YbCuAl was prepared in a closed tantalum-tube filled with argon by a high-frequency induction furnace and then annealed. YbCuAl exhibits the Fe3P-type hexagonal crystal structure (see Fig. 1(c)) with the space group P62m. Magnetic susceptibility of the sample was measured with a superconducting quantum interference device (SQUID) magnetometer at an applied field of 1000 Oe. The temperature dependence of the lattice constants was measured down to 10 K using a Rigaku RINT2000 x-ray diffractometer. We used the Be gasket in-plane geometry where both incoming and outgoing x-ray beams passed through the Be gasket due to its higher transmissivity to the x-rays in comparison to other high-Z materials. Pressures of the sample chamber were monitored using the ruby $R_1$ line fluorescence.

III. RESULTS AND ANALYSES

A. Magnetic susceptibility

The temperature dependence of the magnetic susceptibility, $\chi$, of YbCuAl is shown in Figs. 1(a) and (b). $\chi$ decreases at temperatures above $T \sim 30$ K. We fitted the $1/\chi$ curve with the formula $1/\chi = (T - \Theta_p)/C$ at $T \geq 150$ K as shown in Fig. 1(a), where $\Theta_p$ and $C$ are the Weiss temperature and the Curie constant, respectively. We obtained $C = 2.33$ and $\Theta_p = -29.7$ K. The Curie constant is expressed as $C = N_A \mu_{eff}^2/3k_B$, where $\mu_{eff}$, $N_A$, and $k_B$ are the effective magnetic moment, Avogadro’s number, and Boltzmann’s constant, respectively. In the Yb systems, where valence fluctuations are expected, the effective magnetic moment can be used as a measure of the degree of the valence admixture, since the Yb$^{3+}$ ($4f^1$), $J = 7/2$ ion displays a magnetic moment while the Yb$^{2+}$ ($4f^2$, $J = 0$) is nonmagnetic. $\mu_{eff}$ is estimated to be 4.32$\mu_B$, which is slightly smaller than 4.54$\mu_B$ for Yb$^{3+}$ ion which was calculated using the formula according to Hund’s rule, $\mu_{eff} = g\sqrt{J(J+1)}\mu_B$, where $g$ is the Landé $g$ factor. The effective magnetic moment indicates that Yb is nearly trivalent with a small fraction of the divalent component.

According to the Bethe-Ansatz solution of the Coqblin-Schrieffer model based on the SIAM, the physical properties of a Kondo material are scaled by a single energy parameter, $T_0$. We thus estimated $T_0$ using the numerical results of Rajan’s and Bicker’s approaches as shown in Fig. 1(b). Using the Rajan’s approach, we presented modeled fits with the angular momenta of $J = 3/2$ ($N = 4$), $J = 5/2$ ($N = 6$), and $J = 7/2$ ($N = 8$). The best fit with smaller misfits is obtained for $N = 6$, implying the lift of the degeneracy due to the crystal electric field (CEF). We note that the inelastic neutron scattering experiments on YbCuAl did not have sufficient evidence to clarify the well-defined CEF splittings, because the CEF splittings and the strength of the Kondo interaction were of a similar order. In addition, the 4f-electron contribution to the specific heat, $C_{4f}$, showed a maximum at $T \approx 28$ K with the value $C_{4f} \sim 7$ J mol$^{-1}$K$^{-1}$. Calculated using the Coqblin-Schrieffer model, the peak lies in the value of the maximum specific heat in between $J = 7/2$ ($C_{max} = 8.2$ J mol$^{-1}$K$^{-1}$) and $J = 5/2$ ($C_{max} = 6.2$ J mol$^{-1}$K$^{-1}$). Thus the dominant effect on what in the YbCuAl is the Kondo interaction, although the effect of the CEF likely plays an important role as well. We obtained $T_0 = 93$ K for the Rajan’s fits having $J = 5/2$ and 83 K for the Bicker’s fit, respectively. These values are comparable to $T_K = 100$ K which was previously estimated using a similar method.
FIG. 1. (Color online). (a) Inverse of the magnetic susceptibility (open squares) as a function of temperature with the Curie-Weiss fit at $T > 150$ K. (b) Magnetic susceptibility (open squares) and fits based on the single impurity Anderson models. Estimated characteristic temperatures ($T_0$) are also indicated. (c) Crystal structure. (d) Lattice constants $a$ and $c$ as a function of temperature. (e) Volume as a function of temperature.

B. Lattice constants

The temperature dependences of the lattice constants and the volume are shown respectively in Figs. 1(d) and (e) whereas the crystal structure is presented in Fig. 1(c). While the lattice parameters slightly decrease with temperature down to 70 K, they remain almost constant between 70 and 20 K. As shown in Fig. 4(c) below, this anomalous behavior at low temperatures correlates with the change in the Yb valences.

C. Electronic structure

Figure 2(b) shows the Yb 2p-3d RXES spectra as a function of the incident photon energy across the Yb $L_3$ edge at 300 K and 25 GPa for YbCuAl, while PFY-XAS spectra are shown in Fig. 2(a). The Yb valence is nearly trivalent and a small fraction of Yb$^{2+}$ is present in the spectra. Energy transfer in Fig. 2(b) is defined as the difference between the incident and emitted photon energies. The Raman component of the RXES spectrum is clearly observed in the double peaks which remain at constant transferred energies below the absorption edge at $h\nu \sim 8944$ eV, corresponding to the Yb$^{2+}$ ($4f^{13}$) and Yb$^{3+}$ ($4f^{14}$) configurations. The PFY-XAS peak at $h\nu = 8935$ eV may correspond to the quadrupole (QP) transition, that is insensitive to changes in temperature or pressure in our measurements. Above the absorption edge, the energy transfer of the fluorescence peaks changes linearly with the incident energy. These trends are also visible in the two-dimensional contour maps in Fig. 2(c).

Examples of the fitting results for the PFY-XAS spectrum at 25 GPa and 300 K and the RXES spectrum at $h\nu = 8936$ eV are shown in Fig. 3. We fitted the PFY-XAS and RXES spectra with Voigt functions assuming three components of Yb$^{2+}$ ($4f^{13}$), Yb$^{3+}$ ($4f^{14}$), and fluorescence with the QP component. A proper fit of the Yb$^{3+}$ feature required use of two components. The split of the Yb$^{3+}$ component arises from the splitted 5$d$ band above the Fermi level. In YbInCu$_4$ and related
compounds, we observed a similar split of the $Yb^{3+}$ component that was previously noted.\textsuperscript{21} This spectrum has been well reproduced theoretically in YbInCu$_5$ by assuming two semi-elliptic DOS of the 5$d$ band.\textsuperscript{22} These results indicate that the split of the $Yb^{3+}$ component is mainly due to the nature of the 5$d$ band. Furthermore, the PFY-XAS and RXES spectra reflect the final state effect, but do not involve the ground state directly.\textsuperscript{23} However, the spectra for Yb$_2$O$_3$ displayed the $Yb^{3+}$ peak around 8948 eV, while no evidence of a $Yb^{4+}$ ($4f^{12}$) related feature caused by the final state effect was observed.\textsuperscript{24} We note that the $Yb$ valence estimated from our results agrees well with the values obtained by other methods as described below. This evidence suggests that the final-state effects are, if existing at all, very weak.

The fitting results of the RXES spectra as a function of the incident photon energy in Fig. 2(b) are summarized in Fig. 2(d) along with a PFY-XAS spectrum. The fits enable us to confirm the presence of the quadrupole, fluorescence, and two Raman components, and to track their energy-dependent intensities. The estimated $Yb$ valence, $v$, from the fit for the Raman components in Fig. 2(d) is $2.98 \pm 0.01$ at 25 GPa, which is almost the same as the value derived from the fit to the PFY-XAS spectrum at 25 GPa (see further discussion below).

The temperature dependence of the PFY-XAS spectra at ambient pressure is shown in Fig. 4(a). Figure 4(b) shows the temperature-dependent RXES spectra measured at the $Yb^{2+}$ resonance energy of $\nu = 8938$ eV. Figure 4(c) shows the temperature dependence of the $Yb$ valence estimated from the fit for the PFY-XAS spectra. Temperature dependence of the $Yb$ valence derived from the fit of the RXES spectra after the normalization with the absolute values derived from PFY-XAS, is also shown in the Fig. 4(c). The trends of the $Yb$ valence obtained from both spectra are similar, which further substantiates the reliability of our analytical method. The estimated $Yb$ valence of $2.96 \pm 0.01$ at 300 K and ambient pressure is close to the previously reported values of 2.94 (Ref. 18) from the SIAM fit to the susceptibility and 2.95 (Ref. 25) at ambient pressures. A characteristic temperature of $T_0 = 90$ K and a bandwidth of $\Gamma = 60$ meV were derived from Mössbauer spectroscopic results.\textsuperscript{26} From the relation $n_f/(1 - n_f) = (2J + 1)\Gamma/(\pi T_0)$ ($n_f$ is 4$f$...
fractional occupancy), the result was $v = 2.96$, which is the same value to our estimate.

The Yb valence decreases with decreasing temperature and the most notable change in the valence occurs below 100 K. The temperature dependence of the valence is expressed as the following equation based on the SIAM with the non-crossing approximation\textsuperscript{15}: $v(T) = 2 + n_f(\infty) - (\Delta n_f(T)/\Delta n_f(0))\Delta n_f(0)$, where $n_f(\infty)$ and $\Delta n_f(0)$ are the intermediate temperature limit of the valence and the decrease in the total valence, respectively. We used a universal curve for $\Delta n_f(T)/\Delta n_f(0)$ as a function of the normalized temperature, $T/T_0$ in Fig. 14 of Ref. 15. In the analyses, we derived the band width ($D$) from the electronic specific coefficient to obtain $n_f(\infty)$ using $\gamma = 5$ mJ mol$^{-1}$K$^{-2}$ for LuCuAl.\textsuperscript{17} $D$ and $n_f(\infty)$ are estimated to be approximately 3.1 eV and 0.97, respectively. The temperature dependence of the valence was obtained from fitting the PFY-XAS spectra using two fitting methods: (1) the first is with $\Delta n_f(0)$ as a free parameter assuming $T_0 = 93$ K, which is derived from the fit to the susceptibility; (2) the second is with $n_f(\infty)$, $\Delta n_f(0)$, and $T_0$ as free parameters. The fitted results are shown in Fig. 4(c). The derived value from the first method is slightly higher than the values of the Yb valence obtained experimentally. The derived value from the second method is $T_0 = 159$ K, which is higher than the value previously estimated by Schlottmann.\textsuperscript{18} The second model also initializes $n_f(\infty)$ at 0.97, the same value as the first case, showing that reasonable values of the band width can be obtained from these analyses.

Our results in Fig. 4(c) indicate that the temperature-induced valence change is on the order of 1.5% between 300 and 10 K, and that a change of less than 3% occurs down to 0 K. These are in reasonable agreement with a previous estimate based on the difference of the thermal expansion between YbCuAl and LuCuAl, which yielded 3.5% decrease between 300 and 1.5 K.\textsuperscript{17} The relatively abrupt decrease of the valence at $T < 100$ K inferred from our measurements indicates an increase of the Kondo effect and the $c$-$f$ hybridization in this temperature range, which favors the Yb$^{2+}$ state. Whereas a slight thermal contraction of the unit cell is typically expected in materials upon cooling, as shown in Fig. 1(e) the volume of YbCuAl is nearly temperature independent at $T < 100$ K. In YbInCu$_4$, a first-order valence transition occurs at 42 K.\textsuperscript{27–29} The change in the Yb valence is on the order of 0.1 (Ref. 30) and the corresponding volume change is approximately 0.5%.\textsuperscript{31} In YbCuAl, on the other hand, the change in the volume is approximately 0.3% between 20 and 300 K while the change in the valence is approximately 0.05. Thus, the rate of the volume change roughly corresponds to that of the Yb valence. Our results for YbCuAl interestingly suggest that the effects from the thermal contraction of the unit cell

![FIG. 4. (Color online) Temperature or pressure-dependence PFY-XAS and RXES spectra of YbCuAl at ambient pressure. (a) PFY-XAS spectra of YbCuAl at ambient pressure. (b) RXES spectra. (c) Yb valence estimated from the fit to the PFY-XAS spectra (closed circles). The Yb valence obtained from the RXES spectra in (b) (closed squares) is also shown. These values are scaled with those from the PFY-XAS spectra for comparison. Fitted curves based on the SIAM for the Yb valence are shown with the $\chi T$ curve (long-dashed-line, right vertical axis). (d) Pressure dependence of the PFY-XAS spectra at 300 K. (e) Pressure dependence of the the RXES spectra. (f) Pressure dependent Yb valence (closed circles) estimated from the fit to the PFY-XAS spectra. The Yb valence obtained from the RXES spectra in (e) (closed squares) is also shown. These values in (c) and (f) are corrected to those from the PFY-XAS spectra.](image-url)
and the increase in the $c$-$f$ hybridization and thus the Yb valence compensate each other at temperatures between 20 to 80 K.

We also measured the PFY-XAS spectra and the RXES spectra at $h\nu = 8938$ eV as a function of pressure at 300 K, respectively shown in Figs. 4(d) and 4(e). The pressure dependence of the Yb valence up to 25 GPa is shown in Fig. 4(f). The mean value of $v$ at ambient pressure and room temperature is slightly different in Figs. 4(c) and 4(f), suggesting that the absolute value of $v$ depends on the sample or experimental conditions, including the sample environments inside or outside of the high-pressure cell. However, with the relative changes in the valence within each data set, temperature and pressure are unaffected and should remain robust. We note that at first glance, pressure-induced changes in the lineshapes of the PFY-XAS and RXES data seem small, whereas the amplitude of pressure and temperature-induced valence variations, as estimated from the fits, is comparable (cf. Figs. 4(c), 4(f)).

Under pressure, these spectral bands become broadened, resulting in larger spectral features, and also weaker maximum intensities for constant peak areas. This causes a smearing out of the apparent pressure-induced lineshape changes. For example, band broadening and valence increase under pressure result in a decrease and an increase of the Yb$^{3+}$ peak maximum, respectively, which explains the little variation of the experimental peak intensity. The Yb$^{3+}$ band broadening also obscures the pressure-induced change in the intensity of Yb$^{2+}$ component. As seen in Fig. 4(f), there is a steady, monotonous increase of the valence up to $\sim$13 GPa, although the valence increase is markedly milder with further pressure increase. Even at the maximum pressure of 25 GPa in our study, Yb does not exhibit purely trivalent state. This is at odds with a previous Mössbauer study which suggested that the trivalency is completed at 5 GPa,\textsuperscript{7} showing that the accuracy and reliability of valence estimates based on $L$-edge x-ray spectroscopic measurements of the lanthanide compounds remain to be resolved. Furthermore, given that the onset pressure of the magnetic order is $P_m = 8$ GPa,\textsuperscript{8} our results imply that Yb already begins to behave as a magnetic trivalent ion before reaching a pure trivalent state at $\sim$2.95.

IV. DISCUSSION

The pressure dependence of the physical properties of YbCuAl is reminiscent to that in YbCu$_2$Si$_2$ (Ref. 32) and YbInAu$_2$ (Ref. 8). In YbCu$_2$Si$_2$, we also note that the magnetic order also sets in at $P_m=8$ GPa.\textsuperscript{33} A change in the slope of the pressure-dependent valence occurs at 11 GPa and 300 K and at 9 GPa and 7 K, beyond which the pressure effect is much less significant.\textsuperscript{5} Interestingly, the changes at 9 and 11 GPa pressures are just above $P_m$. Our results also show a clear change in the slope of the pressure-dependent valence for YbCuAl at approximately 13 GPa and 300 K. Based on the analogy with YbCu$_2$Si$_2$,\textsuperscript{5} it is conceivable that the change in the slope of pressure-dependent Yb valence in YbCuAl reduces by about 2 GPa at low temperatures. This transition pressure is also slightly above the onset pressure for the magnetic order.\textsuperscript{8}

The temperature-dependent resistivity revealed Fermi liquid behavior of YbCuAl occurs at approximately 8 GPa.\textsuperscript{8} Although clear evidence for the quantum criticality has not yet been observed for YbCuAl up to 8 GPa, it is generally expected that intersite magnetic fluctuations are enhanced around the magnetic instability point and likely play an important role in the thermodynamical properties of the system. Thus, the change in the slope of the valence increase may suggest that valence fluctuations in YbCuAl, as in YbCu$_2$Si$_2$, are somewhat connected to the magnetic fluctuations induced in the vicinity of the magnetic instability point. Consequently, this also raises the possibility that the flattening in the pressure-dependent valence is a general characteristic nature of the intermediate-valence rare-earth systems at $P$-$T$ conditions close to their cross over magnetic instability points. This calls for additional investigation on the $T^2$-dependence at higher pressures in the YbCuAl system in order to confirm the existence of the QCP. We note that in YbRh$_2$Si$_2$, the QCP was previously confirmed using the divergent character of $C/T$,\textsuperscript{35} where $C$ is specific heat, although the $T^2$ dependence of the resistivity was observed at the QCP. Therefore, in situ observation of the divergence of $C/T$ would provide sufficient experimental evidence to confirm the occurrence of the QCP.

Here we further discuss another perception for the change in the slope of the valence-pressure curve as a possible indication of the transition from the intermediate-valence (IV) regime to the Kondo regime. In the IV regime, the effect of CEF is obscured because of the high Kondo temperature, while in the Kondo regime, $T_K$ is lower than CEF; thereby, a well-defined CEF splitting should be observed at $P$-$T$ conditions in the system. This transition has indeed been observed in YbNi$_2$Ge$_2$. It has been reported that the temperature dependence of the electrical resistivity of YbNi$_2$Ge$_2$ displays a single maximum related to the Kondo effect for pressures below 3.4 GPa, while two maxima are developed at pressures above 3.4 GPa.\textsuperscript{36} This is suggestive of a transition from the IV to the Kondo regime. Indeed, we have observed an inflection in the valence-pressure curve at approximately 5 GPa and room temperature.\textsuperscript{11} Furthermore, from the pressure dependent resistivity data, it has been suggested that the ratio $A/\gamma^2$, known as the Kadowaki-Woods ratio,\textsuperscript{37} increases by nearly two orders of magnitude. This increase in the Kadowaki-Woods ratio is in accordance with the decrease of the ground state degeneracy $N$ due to the CEF effect as $T_K$ is reduced by pressure.\textsuperscript{38} Therefore, the change in the slope of the pressure-dependent valence possibly indicates the transition from the IV to the Kondo regime. This transition is accompanied by the reduction of the ground state degeneracy and a rapid
The volume contraction of what system in temperature cooling favors the magnetic ground state of the Yb$^{3+}$ ion because of its smaller ionic radius compared with the Yb$^{2+}$ ion. The pressure-induced volume contraction at applied pressures up to 10 GPa is typically about one order magnitude larger than in temperature cooling. Our results show that a volume contraction of about 1% occurs with decreasing temperature from 300 to 20 K, corresponding to the application of external pressure slightly less than 1 GPa. At ambient pressure, the c-f hybridization is enhanced by cooling at temperatures below the characteristic temperature of approximately 100 K, resulting in the increase of the Yb$^{2+}$ state which compensates the volume contraction effect.

We finally address on the absence of the reentrance into the mixed-valence state within the pressure range of the present study. As stated in the introduction, the electrical resistivity of YbCuAl at room temperature shows a minimum at 8 GPa, while the electrical resistivity increases with pressure at above 8 GPa. From our results, the possibility for an increase in $T_K$ at pressures above 8 GPa is unlikely, ruling out the scenario proposed by Gol’sev and Abd-Elmeguid, because the Yb valence constantly enhances at applied pressures. Based on this evidence, we would argue that the increase of resistivity at pressures above 8 GPa results from the enhancement in electron scattering due to the increase of the magnetic local moment of the Yb$^{3+}$ ion. We also note that in YbCuAl and YbCu$_2$Si$_2$, the CEF effect can significantly influence their physical properties even at ambient pressure. Although inelastic neutron scattering experiments did not reveal a well-defined CEF excitation, the best fit to our magnetic susceptibility result can be obtained using the Rajan’s model taking into account the degeneracy, indicating the presence of CEF effects even at ambient pressure. Since the reentrance of the valence transition was not observed at high pressure, the anomaly in the pressure-dependent resistivity may arise from the increased local magnetic moment of Yb$^{3+}$ as well as the CEF effect.

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**V. CONCLUSION**

Temperature- and pressure-induced changes in the electronic structure of the heavy fermion compound YbCuAl have been measured using PFY-XAS and RXES at temperatures below 8 K and ambient pressure, and at pressures up to 25 GPa and room temperature. The temperature dependence of the Yb valence is well represented within the SIAM and we have obtained a characteristic temperature comparable to the previous evaluations of the Kondo temperature obtained from the magnetic susceptibility measurements. Our results suggest that the thermal contraction is compensated for by the volume increase caused by enhancement of the c-f hybridization at low temperature due to the Kondo effect. Our pressure-dependent results reveal a continuous increase of the Yb valence up to 25 GPa, with a decrease of the slope around 13 GPa. The proximity of this change in the slope to the magnetic instability confirms the interplay between electronic and magnetic fluctuations. The change in the slope of the pressure-dependent valence data can likely be associated with the transition from the intermediate valence state with no CEF effect to the Kondo regime where CEF effect overcomes $T_K$. As a result, the RKKY interaction dominates the low-temperature properties. Such interplay has already been suggested to occur in YbNi$_2$Ge$_2$ and YbCu$_2$Si$_2$ under pressure, suggesting that this may be an intrinsic nature in all trivalent Yb-bearing systems. A better fit to the magnetic susceptibility is obtained using the Rajan’s model taking into account the degeneracy, indicating the presence of CEF effects even at ambient pressure. Since the reentrance of the valence transition was not observed at high pressure, the anomaly in the pressure-dependent resistivity may arise from the increased local magnetic moment of Yb$^{3+}$ as well as the CEF effect.
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