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# Pair breaking versus symmetry breaking: Origin of the Raman modes in superconducting cuprates

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We performed Raman scattering experiments on superconductivity-induced features in  $\text{Bi}_2\text{Sr}_2(\text{Ca}_{1-x}\text{Y}_x)\text{Cu}_2\text{O}_{8+\delta}$  (Bi-2212),  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  (Y-123), and  $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$  (Tl-2201) single crystals. The results in combination with earlier ones enable us to systematically analyze the spectral features in the doping range  $0.07 \leq p \leq 0.24$ . In  $B_{2g}$  ( $xy$ ) symmetry we find universal spectra and the maximal gap energy  $\Delta_0$  to scale with the superconducting transition temperature  $T_c$ . The  $B_{1g}$  ( $x^2 - y^2$ ) spectra in all three compounds show an anomalous increase of the intensity toward overdoping. The energy scale of the corresponding peak is neither related to the pairing energy nor to the pseudogap, but possibly stems from a symmetry breaking transition at the onset point of superconductivity at  $p_{sc2} \simeq 0.27$ .

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## I. INTRODUCTION

The magnitude of the gap in the quasiparticle excitation spectrum determines the energy difference between the superconducting and the normal state. The momentum dependence  $\Delta_{\mathbf{k}} = \Delta_0 f(\mathbf{k})$  reflects properties of the pairing potential. In the case of the copper-oxygen (cuprate) superconductors the  $d$ -wave character of the gap is well established<sup>1</sup>, suggesting the influence of repulsive contributions to the pairing interaction<sup>2-4</sup> while the energy scales, in particular as a function of (hole) doping  $p$  with  $p = 1 - n$  the number of mobile holes per planar Cu atom, remain elusive. There are substantial differences between the overall behavior in the areas close to  $(\pi, 0)$  and equivalent points (anti-node) and  $(\pi/2, \pi/2)$  (node) in the Brillouin zone (BZ) of the  $\text{CuO}_2$  planes which are hard to pin down experimentally. The theoretical interpretation is accordingly controversial. While the nodal gap is believed to reflect more or less the pairing interaction, the anti-nodal energy scale probably originates from a rather complex interrelation of superconductivity and another instability<sup>5-7</sup>. This holds particularly true for the underdoped range,  $p \leq 0.16$ . In some experiments the maximum of the superconducting gap,  $\Delta_0(p)$ , is found to stay essentially constant in wider doping ranges around  $p = 0.16$  and to decrease only close to the end points  $p_{sc1}$  and  $p_{sc2}$  of the superconducting dome<sup>8,9</sup>. Other experiments indicate  $\Delta_0(p)$  to more or less follow the superconducting transition temperature  $T_c$ <sup>10-16</sup>. Finally, studies of the heat transport indicate that  $\Delta_0(p)$  decreases continuously with increasing doping<sup>17-21</sup>.

Below  $T^* > T_c$ , a second energy gap  $\Delta^*(p)$  opens up.  $\Delta^*$  is the typical range over which spectral weight is suppressed in the vicinity of the anti-nodal points of the BZ and is usually referred to as the pseudogap<sup>5,22-25</sup>. In the spectroscopic experiments  $T_c$  and the pseudogap temperature  $T^*$  are found to become indistinguishable in the experiment-dependent range  $0.15 < p^* < 0.20$ . However, there still remain two energy scales below  $T_c$  exhibiting different doping dependences<sup>5,10,12,13,25-27</sup>. The scale observed close to the BZ diagonal appears to follow  $T_c$  quite closely in those experiments which, as pointed out first by Deutscher<sup>13</sup>, probe the condensate such as Andreev reflection or electronic Raman scattering. The anti-nodal scale is approximately proportional to  $(1 - p/p_0)$  with  $0.20 < p_0 < 0.30$  and was first identified in the quasiparticle spectrum<sup>13,28,29</sup> but also anticipated from Raman experiments on the overdoped side<sup>30</sup>. The functional dependence on  $p$ , yet not necessarily the magnitude, of this scale is remarkably similar to that of  $\Delta^*(p)$  (for more details see Refs. 7,25). However, above  $p^*$  there is no longer any suppression of spectral weight in the normal state<sup>8-10,12,14-16,25,27,30-32</sup>, and superconducting coherence peaks are observed everywhere on the Fermi surface by angle-resolved photoemission (ARPES)<sup>9,26,32</sup> and, independent of the location on the sample surface, by scanning tunneling spectroscopy (STS)<sup>8,27,33</sup>.

The wide ranges of the characteristic doping levels  $p^*$  and  $p_0$  as well as of the energy scales  $\Delta^*(p)$  and  $\Delta_{\mathbf{k}}(p)$  are not fully explored yet. There are indications of a dependence on the material class, the samples, and also on the experimental probe. The pseudogap  $\Delta^*(p)$  was first

observed by optical (IR) conductivity in  $\text{YBa}_2\text{Cu}_3\text{O}_{6+x}$  (Y-123) with light polarized along the  $c$ -axis<sup>22</sup>. Given the band structure of Y-123 this experiment emphasizes the anti-node<sup>34,35</sup>. Soon thereafter, an anti-nodal gap above  $T_c$  was observed in ARPES experiments<sup>23,24</sup>. In Raman experiments, two types of gap-like features are found in the  $B_{1g}$  and  $B_{2g}$  symmetries<sup>10,14,36,37</sup> for  $p < 0.21$ . An intensity suppression in  $B_{2g}$  symmetry of order 5 – 10 % occurs below  $T^*$  in an almost doping independent energy range of approximately 100 meV ( $800 \text{ cm}^{-1}$ )<sup>10,14,36</sup> comparable to the exchange energy  $J$ . In  $B_{1g}$  symmetry, the Raman intensity is suppressed in an energy range up to at least 200 meV with the suppression setting in abruptly below  $p \approx 0.21$ <sup>37,38</sup>. In the same doping range the superfluid density starts to decrease<sup>6</sup>, and an asymmetry with respect to zero bias of the scanning tunneling spectra develops for specific locations on the crystals<sup>8,27,33,39</sup>. Also below  $T_c$ , the spectra develop a strong dependence on location. In some areas there are still coherence peaks while in other spots a much larger gap without well-defined peaks is observed. The distribution of gaps becomes wider and ranges from 5 to 12 in units of  $k_B T_c$ .

In a recent Raman experiment the energy of the anti-nodal pair-breaking peak of fully oxygenated  $\text{YBa}_2\text{Cu}_3\text{O}_{\sim 7}$  was studied as a function of applied pressure<sup>40</sup>. The superconductivity-induced peak position decreases by approximately 50 % in the pressure range up to 22.3 GPa while  $T_c$  changes only by 25 %<sup>40,41</sup>. In units of  $k_B T_c$  the position of the peak moves from 6.2 at ambient pressure to 4.2 at 22.3 GPa. Particularly the last result casts doubt on the prevailing interpretation of the  $B_{1g}$  Raman spectra in the superconducting state in terms of a direct relationship to the pairing energy or the pseudogap.

In this paper, we present new electronic Raman scattering experiments on Y-123,  $\text{Bi}_2\text{Sr}_2(\text{Ca}_{1-x}\text{Y}_x)\text{Cu}_2\text{O}_{8+\delta}$  (Bi-2212), and  $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$  (Tl-2201) and put them into context with earlier results. We systematically study the sample dependence and, as an additional variable, the intensity of the superconductivity-induced features for doping levels  $0.07 \leq p \leq 0.24$ . The results in  $B_{2g}$  symmetry show that the momentum dependence of the superconducting gap,  $f(\mathbf{k}) = \Delta_{\mathbf{k}}/\Delta_0$ , hardly depends on doping for either Y-123 or Bi-2212. For  $p > 0.16$ , the anti-nodal spectra of Bi-2212 reflect neither the pseudogap nor the superconducting gap. Rather, the doping dependence of both the intensity and the energy of the superconductivity-induced modes suggests that it is closely related to the onset point of superconductivity at  $p_{sc2} = 0.27$  on the very overdoped side of the phase diagram.

The paper is organized as follows: In Section II we describe details of the experiment, the samples studied here as well as those investigated in our earlier work. In Section III we describe in detail the results of the new experiments with the emphasis placed on the comparison of samples with nearly equal transition temperatures. In Section IV we compile and discuss the results of most of

our experiments obtained over the years. The energy scales are compared to the results from other Raman groups and to those from ARPES and STS. In Section V, we summarize our results and conclusions and formulate open questions.

## II. EXPERIMENT AND SAMPLES

Momentum-dependent electron dynamics such as electron-hole excitations in normal metals, gaps in superconductors, or collective modes can be studied by Raman scattering via the intracell fluctuations of the charge density excited by the photons. As a consequence, different parts of the Brillouin zone (BZ) can be projected out independently by appropriately adjusting the polarizations of the incoming and outgoing photons<sup>42</sup>. In the cuprates,  $B_{1g}$  and  $B_{2g}$  spectra emphasize anti-nodal and nodal electrons, respectively, with the form factors shown in the insets of Fig. 1 (a) and (b). Collective excitations with pure symmetries can be accessed separately. In the superconducting state, the condensate is directly probed, since the anomalous part of the Green function is measured in addition to the normal one<sup>42–45</sup>.

The spectra were measured with standard Raman equipment using the  $\text{Ar}^+$  line at 458 nm and, for  $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$ , at 514 nm. The temperatures generally refer to the illuminated spot and are typically between 5 and 10 K above those of the sample holder as determined from the comparison of energy-gain and -loss spectra. All spectra shown here are Raman response functions  $R\chi''_{I,S}(\Omega, T)$  where  $I$  and  $S$  refer to incoming and scattered photons.  $R\chi''_{I,S}(\Omega, T)$  is obtained from the cross section via the fluctuation-dissipation theorem to remove trivial temperature dependences. This means that the experimental intensity is divided by the thermal Bose factor,  $\{1 + n(\Omega, T)\} = [e^{-\hbar\Omega/k_B T} - 1]^{-1}$  after having been corrected for the sensitivity of the instrument and multiplied by  $\omega_I/\omega_S$ .  $\Omega = \omega_I - \omega_S$  is the energy transferred from the photons to the system. The polarizations of the incoming and scattered photons,  $\mathbf{e}_I, \mathbf{e}_S$ , are always indicated symbolically by two arrows and referred to as  $xy$  and  $x'y'$  using Porto notation with  $x' = 1/\sqrt{2}(1, 1)$  and  $y' = 1/\sqrt{2}(1, -1)$ . The constant  $R$  absorbs experimental factors and converting the units of  $\chi'' = \Im\chi$  into ( $\text{counts s}^{-1} \text{ mW}^{-1}$ ). All spectra displaying  $R\chi''_{I,S}(\Omega, T)$  can be compared on an absolute scale.

The samples studied here were prepared in Garching, Stanford, Tokyo, and Vancouver (see Table I).  $(\text{Y}_{0.92}\text{Ca}_{0.08})\text{Ba}_2\text{Cu}_3\text{O}_{6.3}$  (Y-UD28) single crystals were grown in  $\text{BaZrO}_3$  crucibles using starting materials with purities of better than 99.999 % (5N). In this way ultrahigh quality samples can be obtained<sup>50,51</sup>. The  $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$  (Bi-OPT94) and  $\text{Bi}_2\text{Sr}_2(\text{Ca}_{0.92}\text{Y}_{0.08})\text{Cu}_2\text{O}_{8+\delta}$  (Bi-OPT96) single crystals

TABLE I: Complete list of samples used for Raman scattering. In rows (a-g) we describe samples which were measured in this study. In rows (h-t) the properties of samples studied earlier are compiled. Samples (a,b,h,i,j) were prepared by A. Erb (WMI Garching), (c,s) by Shimpei Ono and Yoichi Ando (CRIEPI, Tokyo and Osaka University), (d,e) by A. Damascelli, H. Eisaki, and M. Greven (Stanford, Ref. 46 and Vancouver), (f,g) by D. C. Peets, W. N. Hardy, R. Liang, and D. A. Bonn (Vancouver, Ref. 47), (k,l,t) by H. Berger and L. Forró (EPFL, Lausanne, Ref. 10,37), (m-r) by B. Revaz (University Geneva, Ref. 14,37). The transition temperatures were measured resistively (k,l,t), via magnetometry (h-j,d, m-o), or via the non-linear ac response (a-g,p-s). Hence, in some cases we used 2 methods and found good agreement of  $T_c$  values when appropriately defined. In the cases of resistivity and linear susceptibility measurements the transition widths are defined by the 10 and 90% points. If only the non-linear susceptibility was measured  $\Delta T_c$  is estimated from the shape of the signal close to  $T_c$ , and the uncertainty can be as large as 50%. All doping levels  $p$  are derived from  $T_c$  via the relation  $p = 0.16 \mp 0.11\sqrt{1 - T_c/T_c^{\max}}$  using  $T_c^{\max} = 94 \text{ K}^{48}$ . The resulting error is of order  $\pm 0.01$ . In Y-123 (and, probably, also in Bi-2212), the experimental  $T_c$  is systematically below the parabola in the vicinity of the 1/8 anomaly<sup>49</sup> making the determination of  $p(T_c)$  less precise between 0.10 and 0.15.

label	sample	sample ID	doping	$T_c$ (K)	$\Delta T_c$ (K)
a	$(Y_{0.92}Ca_{0.08})Ba_2Cu_3O_{6.3}$	Y-UD28	0.07	28	2
b	$YBa_2Cu_3O_{6.4}$	Y-UD29	0.07	28	3
c	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OPT94	0.16	94	2
d	$Bi_2Sr_2(Ca_{0.92}Y_{0.08})Cu_2O_{8+\delta 1}$	Bi-OPT96	0.16	96	2
e	$Bi_2Sr_2(Ca_{0.92}Y_{0.08})Cu_2O_{8+\delta 2}$	Bi-OD87	0.19	87	2
f	$Tl_2Ba_2CuO_{6+\delta}$	Tl-OD78	0.21	78	5
g	$Tl_2Ba_2CuO_{6+\delta}$	Tl-OD46	0.24	46	5
h	$YBa_2Cu_3O_{6.5}$	Y-UD60	0.10	58	5
i	$YBa_2Cu_3O_{6.93}$	Y-OPT93	0.16	93	0.5
j	$YBa_2Cu_3O_{6.99}$	Y-OD87	0.18	87	2
k	$Bi_2Sr_2(Ca_{0.62}Y_{0.38})Cu_2O_{8+\delta}$	Bi-UD58	0.10	58	5
l	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-UD92	0.15	91.7	-
m	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OPT92	0.16	92	1
n	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD82	0.20	82	3
o	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD78	0.21	77.8	0.2
p	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD65	0.22	65	1
q	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD62	0.22	62	1
r	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD56	0.23	56	5
s	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD56	0.23	56	5
t	$Bi_2Sr_2CaCu_2O_{8+\delta}$	Bi-OD55	0.23	55	5

were prepared by the floating zone method in the mirror furnace<sup>46</sup>.  $Tl_2Ba_2CuO_{6+\delta}$  (Tl-OD78 and Tl-OD46) were prepared in alumina crucibles<sup>47</sup>. The parameters of the freshly prepared samples are compiled in Table I (a-g). Since one of the central messages of the paper is the dependence of the  $B_{1g}$  results on details of the samples it is crucial to compare the new results to those obtained earlier on different sample sets [Table I (h-t)] under comparable conditions.

### III. RESULTS

In Fig. 1 we plot raw data of new measurements on high-quality  $(Y_{0.92}Ca_{0.08})Ba_2Cu_3O_{6.3}$  (Y-UD28,  $T_c = 28 \text{ K}$ ) (a,b),  $Bi_2Sr_2CaCu_2O_{8+\delta}$  (Bi-OPT94,  $T_c = 94 \text{ K}$ ) (e,f), and  $Bi_2Sr_2(Ca_{0.92}Y_{0.08})Cu_2O_{8+\delta}$  (Bi-OPT96,  $T_c = 96 \text{ K}$ ; Bi-OD87,  $T_c = 87 \text{ K}$ ) (c,d,g,h) single crystals.

Shown are spectra right above and well below the transition temperature  $T_c$ . In spite of the almost identical  $T_c$ s, the two optimally doped Bi-2212 samples [Fig. 1 (c-f)] exhibit substantial differences in the  $B_{1g}$  spectra (c,e) at  $T \rightarrow 0$ . The peak energy  $\Omega_{\text{peak}}^{B_{1g}}$  of sample Bi-OPT96 is approximately 25% higher than that of Bi-OPT94 while  $T_c$  differs only by 2%. The variation of the peak position is accompanied by a change in the amplitude  $A_{sc}$ , i.e. the difference between the superconducting and the normal-state spectra at the peak maximum, by a factor of 2.7. On the other hand, the  $B_{2g}$  spectra exhibit only minor changes in shape, amplitude, and peak energy.

The overdoped sample Bi-OD87 [Fig. 1 (g,h)] was prepared from Bi-OPT96 by oxygen annealing. Both peak frequencies move downward along with  $T_c$ , with a tendency of the  $B_{1g}$  peak to move more rapidly than the  $B_{2g}$  peak as observed earlier in Bi-2212<sup>14,15,30,37</sup>, Y-123<sup>10,14,52,53</sup> and  $HgBa_2CuO_{6+\delta}$ <sup>16</sup>.

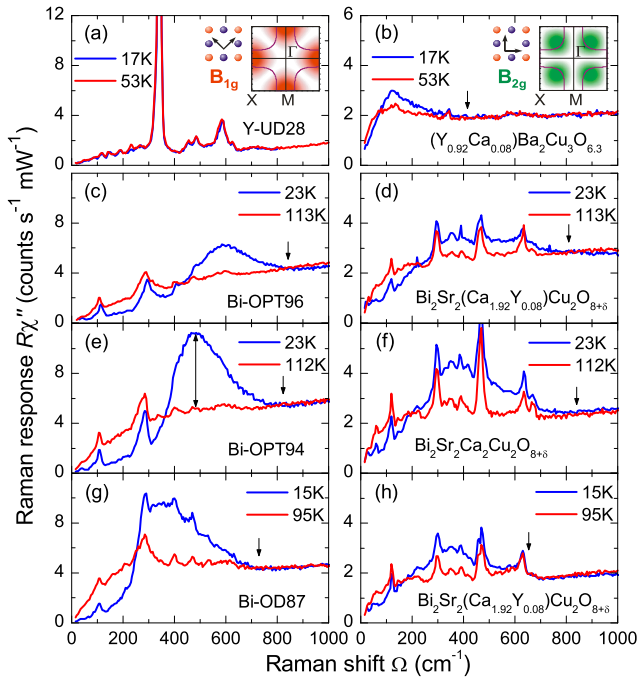


FIG. 1: (Color online) Raman response  $R\chi''(\Omega, T)$  (raw data) of (a,b)  $(Y_{0.92}Ca_{0.08})Ba_2Cu_3O_{6.3}$  (Y-UD28) and (c–h)  $Bi_2Sr_2(Ca_{1-x}Y_x)Cu_2O_{8+\delta}$  (Bi-OPT94, Bi-OPT96, Bi-OD87) in  $B_{1g}$  and  $B_{2g}$  symmetries as indicated. The corresponding light polarizations and sensitivities in the Brillouin zone are shown in the insets with copper and oxygen atoms displayed in red and blue, respectively. In (e) a double-headed arrow indicates the amplitude  $A_{sc}$  of the superconductivity-induced peak. Whenever applicable, a down-pointing arrow gives the approximate position where normal-state and superconducting spectra merge.

On the underdoped side we studied Y-123 for its superior crystal quality<sup>50</sup>. We find superconductivity to be observable only in  $B_{2g}$  symmetry. The peak energy is at approximately one third of that observed at optimal doping and follows  $T_c$ . The absence of superconductivity-induced peaks in  $B_{1g}$  symmetry appears to be a generic feature of underdoped cuprates with  $p \leq 0.13$  (for a discussion see Ref. 42) which occurs approximately in the same doping range as the loss of coherence close to the anti-nodal points<sup>26,54</sup> and the loss of spectral weight in the oxygen  $K$  edge absorption<sup>55,56</sup>.

In Fig. 2 we show results on a freshly prepared set<sup>47</sup> of high-quality  $Tl_2Ba_2CuO_{6+\delta}$  single crystals at  $p = 0.21$  and  $p = 0.24$  on the overdoped side, Tl-OD78 and Tl-OD46, with transition temperatures of 78 and 46 K, respectively. The crystals were relatively small with maximal dimensions of less than 1 mm. The spectra were measured on as-grown surfaces since cleaving was not attempted for the small number of samples available. Residual flux and adsorbates lead to a contribution from the laser line (at  $\Omega = 0$ ) extending to energies as high as  $50\text{--}100\text{ cm}^{-1}$  [see in particular Fig. 2 (b)] and temperature-dependent variations of the overall Raman

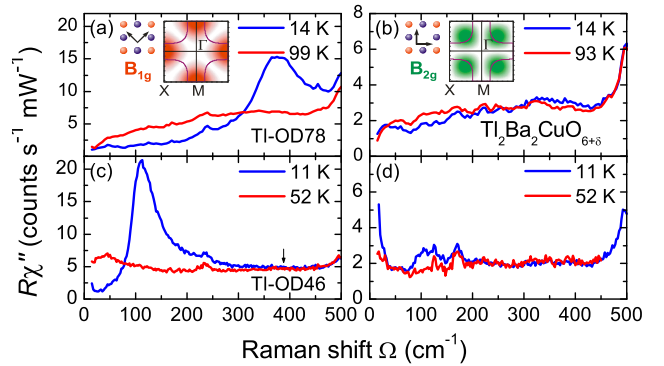


FIG. 2: (Color online) Raman response  $R\chi''(\Omega, T)$  of  $Tl_2Ba_2CuO_{6+\delta}$  with  $T_c = 78$  K (Tl-OD78) [panels (a), (b)] and  $T_c = 46$  K [panels (c), (d)] (Tl-OD46). Note that the energy scales are different from those in Fig. 1. The spectra were measured with the laser line at 514 nm. The superconducting spectra were multiplied by constant scaling factors  $s_0$  in the range  $0.9 \leq s_0 \leq 1.1$  to match the intensities above the pair-breaking range with those in the normal state.

intensity in the range  $\pm 20\%$ . Consequently, the overall intensities are not as quantitative as those of Y-123 and Bi-2212 with a stability of better than  $\sigma = 5\%$ . By using the laser line at 514 nm the cross sections of Tl-2201 become comparable to those of Y-123 and Bi-2212 and do not depend significantly on doping. With excitation at 458 nm resonance effects enhance the absolute intensity in Tl-OD78 by almost a factor of 2 over that measured with the line at 514 nm. Similar changes for both the continuum and the superconductivity-induced peaks have been reported earlier for excitation between 416 and 755 nm<sup>57</sup>. Kang *et al.* showed also that the resonances essentially disappear for wavelengths  $\lambda \geq 514$  nm. Therefore, it is safe to conclude that the relative change of the normalized  $B_{1g}$  pair-breaking intensity is close to 2 for the doping levels studied here.

The  $B_{1g}$  peaks are clearly resolved and found at 370 and  $120\text{ cm}^{-1}$  for Tl-OD78 and Tl-OD46 corresponding to 6.8 and  $3.7 k_B T_c$ , respectively. The  $B_{2g}$  peaks are very weak, with maxima close to those in  $B_{1g}$ . While the energy for Tl-OD78 is in the same range as in Bi-2212, that of Tl-OD46 is smaller than in any other sample studied here. However, low peak energies at doping levels above 0.22 have been observed before<sup>11,57,58</sup>. Only Gasparov *et al.* measured both the  $B_{1g}$  and the  $B_{2g}$  channel, and they observed peak energies and relative intensities close to those here. For Tl-OD46, the analysis of the  $B_{2g}$  peak yields a line-shape similar to that of the  $B_{1g}$  mode. This suggests a polarization leakage originating from either a misorientation by approximately  $4^\circ$  or from internal strain fields and an orthorhombic distortion in Tl-2201<sup>59–61</sup>. We note that the orthorhombic distortion is not larger than 1% and is not enough to explain the leakage. It is concluded that a quantitative analysis of the  $B_{2g}$  spectra of Tl-2201 is premature. The change

in energy and intensity of the  $B_{1g}$  mode as a function of doping seems to be robust and in qualitative agreement with the published literature. The variation of the intensity will be of particular interest later.

## IV. DISCUSSION

### A. Energy scales derived from the $B_{2g}$ spectra

Fig. 1 shows that the  $B_{2g}$  peak energies in the superconducting state follow  $T_c$  as pointed out earlier in various publications<sup>10,13–15,37,42,52,62,63</sup>. Beyond that we demonstrate here that the entire  $B_{2g}$  spectra can be scaled by normalizing the energy axis of each sample to the respective experimentally determined  $T_c$  (N.B., not the doping  $p$  which is a derived quantity) and the intensity to unity at energies in the range 800–1000  $\text{cm}^{-1}$ . As shown in Fig. 3, the superconducting  $B_{2g}$  spectra collapse on universal curves for both Y-123 and Bi-2212. We note that the  $B_{2g}$  spectrum of Tl-OD46 is not consistent with this picture [see Fig. 2] while  $\text{HgBa}_2\text{CuO}_{4+\delta}$  (Hg-1201) at  $p = 0.24$  fits reasonably well<sup>63</sup>. The reasons for the discrepancies of Tl-OD46 are given in section III and are not considered sufficiently significant to challenge the scaling argument. Further work is needed here.

The low-energy part of the normalized spectra can be described quantitatively in terms of a  $d_{x^2-y^2}$  gap<sup>64</sup>. Naturally, the description fails at higher energies since only the weak coupling limit is considered, which neglects the strong interactions responsible for the large self energy of the electrons<sup>26</sup> and, hence, the Raman spectra at high energies<sup>42,65</sup>. With the gap represented by

$$\Delta_{\mathbf{k}} = \frac{\Delta_0}{2} [\cos k_x - \cos k_y] \quad (1)$$

we find agreement between theory and experiment up to and slightly beyond the pair-breaking peak (see Fig. 3). As shown recently for optimal doping, the range of agreement between the Raman spectra and the  $d$ -wave prediction can be extended if the self-energy effects such as observed by ARPES<sup>26,66</sup> are included using a strong coupling approach<sup>65</sup>.

While the  $B_{2g}$  maximum  $\Omega_{\text{peak}}^{B_{2g}}(p)$  itself scales approximately as  $6k_B T_c$  consistent with previous reports (see references above) the gap maximum  $\Delta_0$  from the  $d$ -wave fit in Fig. 3 follows  $T_c$  as

$$2\Delta_0(p) = (9.3 \pm 0.5) k_B T_c(p) \quad (2)$$

corresponding to approximately twice the canonical weak-coupling BCS result of 4.21 for a two-dimensional  $d$ -wave gap. Thus the  $B_{2g}$  spectra of high-quality Y-123 and Bi-2212 samples provide a wealth of evidence that both the gap ratio  $\Delta_0/k_B T_c$  as seen by Raman scattering and the momentum dependence  $f(\mathbf{k})$  remain unchanged throughout the entire doping range studied

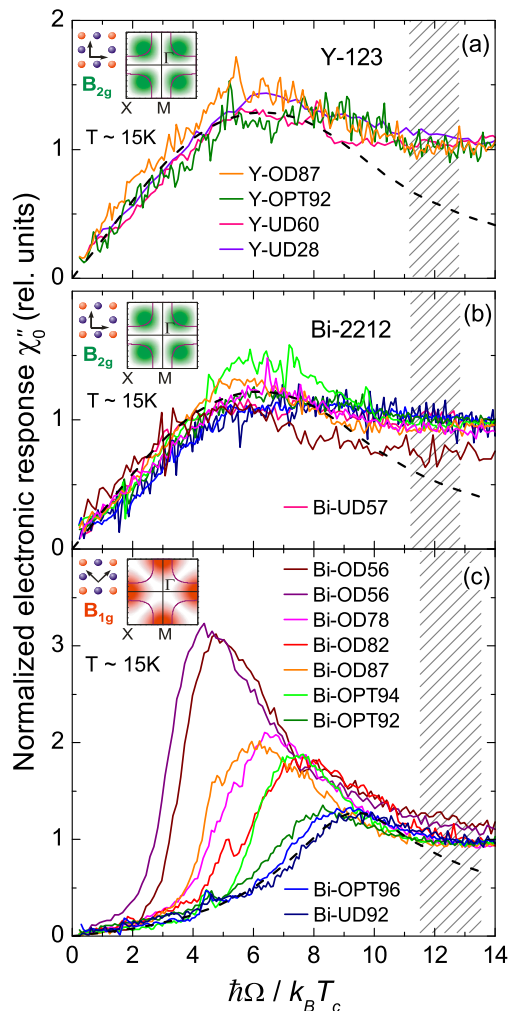


FIG. 3: (Color online) Normalized electronic Raman response  $\chi''_0(\Omega, p)$  (phonons subtracted) of  $(Y_{1-y}\text{Ca}_y)\text{Ba}_2\text{Cu}_3\text{O}_{6+x}$  in  $B_{2g}$  symmetry (a) and  $\text{Bi}_2\text{Sr}_2(\text{Ca}_{1-x}\text{Y}_x)\text{Cu}_2\text{O}_{8+\delta}$  in  $B_{2g}$  (b) and  $B_{1g}$  symmetry (c). Spectra from samples other than those shown in Fig. 1 are taken from our published work<sup>10,14,37,53</sup>. For clarity, the phonons have been subtracted. The energy axes are normalized to the individual transition temperatures. All superconducting spectra merge with the normal-state response in the shaded range. The theoretical weak coupling pair-breaking spectra (dashed lines) are the Tsuneto function on a realistic band structure weighted with the vertices for  $B_{1g}$  and  $B_{2g}$  symmetry.  $2\Delta_0 = 9.3 k_B T_c$  and a phenomenological broadening of 20% was used.

( $0.08 \leq p \leq 0.23$ ). Similar conclusions may apply to Hg-1201<sup>63</sup> although the scaling was not attempted.

The weak dependence of the  $B_{2g}$  spectra on sample details can be explained straightforwardly by small changes in the concentration of scattering centers<sup>67</sup> resulting from defects or quenched disorder, for instance<sup>46</sup>. The doping independence of the normalized  $B_{2g}$  spectra includes the intensity, the position and the shape of the superconducting peaks (see Figs. 1 and 3).

## B. Interrelation of $B_{1g}$ and $B_{2g}$ spectra

The doping dependence of the  $B_{2g}$  spectra imposes constraints on the interpretation of the  $B_{1g}$  spectra since the two symmetries are linked by the form factors<sup>42</sup>. A potential change in the momentum dependence  $f(\mathbf{k})$  of the gap would inevitably leave imprints on both the  $B_{1g}$  and  $B_{2g}$  spectra demonstrating that the disappearance of the  $B_{1g}$  gap structures for  $p < 0.15$  [see, e.g., Fig. 1 (a)] is an effect which, to the resolution possible in this experiment, is occurring in the  $B_{1g}$  symmetry alone. This conclusion is supported by the results in the normal state. For  $p > 0.21 \pm 0.01$  the spectra and, in particular, the electronic relaxation rates are essentially isotropic<sup>54,62</sup>. For  $p \leq 0.21$  the anisotropy between the symmetries develops almost abruptly<sup>54</sup>. With less resolution in  $p$ , this was corroborated recently for Hg-1201<sup>16</sup>, Bi-2212<sup>38</sup>,  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  (LSCO) and Tl-2201<sup>68</sup>.

Below  $p \simeq 0.21$ , the oxygen  $K$  edge absorption, which, supposedly, is related to the number of holes, starts to decrease<sup>55,56</sup> suggesting a transition between metal and Mott physics. Similarly,  $p = 0.21$  is close to  $p = 0.23$  where recent ARPES experiments on Y-123 indicate that not only the anti-nodal but also the nodal quasiparticle weight  $Z_N$  starts to decrease<sup>69</sup>. In earlier ARPES studies below  $T_c$ , the spectral weight was observed to be lost predominantly at the anti-node for  $p \leq 0.18$  (Refs. 32, 70). More recently, ARPES studies on Tl-2201 indicated a further sharpening of the anti-nodal quasiparticle peak for doping levels above  $p = 0.21$  which, however, is more likely to originate from a reduced quasiparticle scattering rate  $\Gamma_{AN}$  than from a further increase of the weight  $Z_{AN}$  (Ref. 71). For  $p < 0.15$ , the Fermi surface shrinks to arcs in the vicinity of the nodal direction<sup>69,72</sup> although, at least for LSCO, neither is the spectral weight at the anti-node completely lost for  $p > 0.03$  nor are there rapid changes in the weight at specific  $k$ -points<sup>73,74</sup>. A loss of coherent quasiparticle weight is also observed in the specific heat<sup>6</sup> and in the tunneling spectra<sup>8,27,33</sup>.

The reduction of  $Z_{AN}$  is in at least qualitative agreement with the variation of the  $B_{1g}$  spectra in the normal state which, for  $p$  decreasing from 0.21 to 0.16, lose 30 to 40% of their spectral weight in the range up to  $2000\text{ cm}^{-1}$  and change their shape<sup>54,63</sup>. In contrast, in  $B_{2g}$  symmetry the shape is by and large conserved and the overall intensity even increases by approximately 20% (Ref. 14,75). This observation is at variance with the reduction of the nodal quasiparticle weight  $Z_N$ <sup>69</sup> at least when the Raman continuum is considered to originate from particle-hole excitations in the lowest order approximation. We emphasize that the spectral changes in  $B_{1g}$  symmetry have to be distinguished clearly from the influence of the pseudogap on the  $B_{2g}$  spectra for  $T < T^*$  which is a 5 – 10% effect in the range up to  $800\text{ cm}^{-1}$  (Refs. 10,14,36,76). Already the different energy ranges indicate that the gap-like phenomena in  $B_{1g}$  and  $B_{2g}$  symmetry cannot have the same origin. Apparently, the dichotomy between the energy ranges and

spectral weight changes is the generic behavior if (i) the samples are sufficiently clean and (ii) resonance effects are unimportant (see above and Ref. 68). Hence, from the viewpoint of a very large set of Raman data the substantial loss of spectral weight both above and below  $T_c$  is a peculiarity of the  $B_{1g}$  symmetry.

Our interpretation of this dichotomy between  $B_{1g}$  and  $B_{2g}$  symmetry in terms of two independent phenomena is not generally shared. Chen<sup>52</sup>, Le Tacon<sup>16</sup>, Blanc<sup>63</sup> and coworkers interpret the suppression of the  $B_{1g}$  and part of the  $B_{2g}$  pair breaking features in terms of a gradual loss of quasiparticle weight starting at the anti-node as qualitatively observed by ARPES. In this scenario 100% of the  $B_{1g}$  and 50 to 80% of the  $B_{2g}$  intensities are lost<sup>63</sup> towards underdoping. However, this interpretation cannot easily be reconciled with the experimental observations found consistently in four cuprate families<sup>68</sup> since, as mentioned above, the  $B_{2g}$  spectra in the normal state (including  $T < T^*$ ) are largely doping independent. We believe that one has to go beyond the lowest-order particle-particle correlation function for the Raman response to resolve the apparent discrepancies between single- and two-particle properties.

To summarize this part, the Raman spectra in  $B_{1g}$  and  $B_{2g}$  symmetry both in the normal and the superconducting state exhibit distinctly different doping dependences occurring in rather different energy ranges which make it difficult to explain that they have the same origin. We rather argue that the suppression of spectral weight in  $B_{1g}$  symmetry is linked to the correlation or Mott gap<sup>54,77</sup> while the effects of the pseudogap are observed predominantly in  $B_{2g}$  symmetry and are relatively small<sup>10,14,36,76</sup>. In the narrow doping range  $0.12 < p < 0.16$  there may also be effects of the pseudogap on the  $B_{1g}$  spectra<sup>38</sup> before any relationship to the carriers is quenched by the correlation gap below  $p \sim 0.12$ . On the other hand, symmetry protects the Fermi liquid-like state of the nodal electrons down to  $p \simeq 0.05$  where superconductivity disappears. For  $p \lesssim 0.05$  the overall  $B_{2g}$  intensity decreases essentially linearly with doping before it vanishes at  $p = 0$ <sup>75</sup>. Presently, we cannot provide a microscopic model for the selective suppression of the  $B_{1g}$  spectra neither in the superconducting nor in the normal state<sup>54</sup>. Yet, the inclusion of higher order diagrams with symmetry-dependent cancelation effects for spin and charge channels may be a viable way toward an explanation<sup>78</sup>.

## C. Properties of the mode in $B_{1g}$ symmetry

Now, the focus will be placed on those properties of the  $B_{1g}$  mode appearing below  $T_c$  which have not been analyzed previously. In Fig. 3 (c) electronic  $B_{1g}$  spectra of Bi-2212 are plotted in units of  $k_B T_c$ . As a general trend, the peaks soften from 9.0 to 4.5  $k_B T_c$  for  $p$  increasing from 0.15 to 0.23. The variation of the peak energies is not monotonic. For example, the peaks of samples Bi-

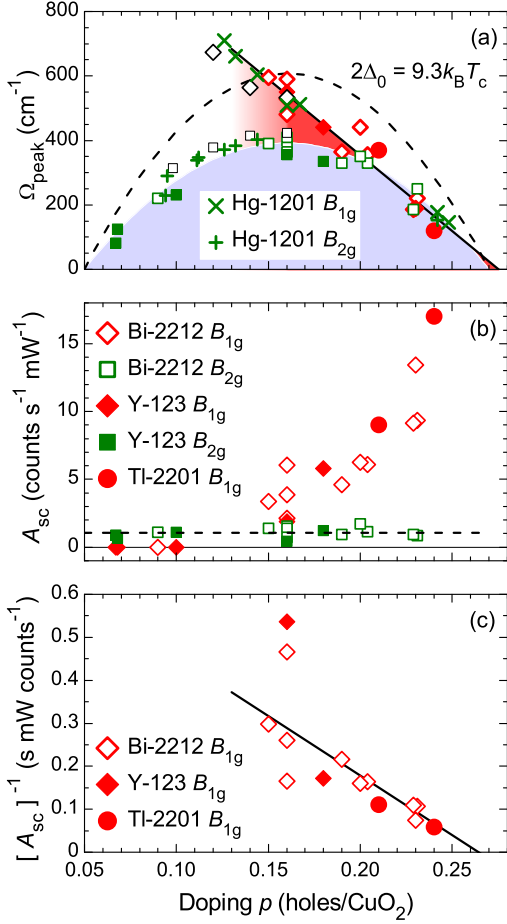


FIG. 4: (Color online) Doping dependence of the superconductivity-induced features in Y-123 (full symbols), Bi-2212 (open symbols), HgBa<sub>2</sub>CuO<sub>4+ $\delta$</sub>  (Hg-1201, crosses) and Tl-2201 in  $B_{1g}$  symmetry (full circles). The data for Hg-1201 and those of Bi-2212 at  $0.10 \leq p \leq 0.16$  (slightly smaller symbols) are taken from Ref. 63. (a) Peak energies  $\Omega_{\text{peak}}$ .  $\Omega_{\text{peak}}^{B_{2g}}$  [The crosses are for Hg-1201, other symbols are given in panel (b)] is smaller than  $2\Delta_0(p) = 9.3k_B T_c^{\text{max}}[1 - 82.6(p - 0.16)^2]$  with  $T_c^{\text{max}} = 94$  K (dashes). The same holds true for  $\Omega_{\text{peak}}^{B_{1g}}$  (diamonds, circles and crosses) at  $p > 0.16$ . A linear fit to the Bi-2212 data (straight full line) extrapolates to  $p_{\text{sc}2} \simeq 0.27$ . (b) Amplitudes  $A_{\text{sc}}$  in  $B_{1g}$  and  $B_{2g}$  symmetries (only our own data). The horizontal line at 1.03 cps/mW is the average of the amplitudes in  $B_{2g}$  symmetry. (c) Inverse  $B_{1g}$  amplitudes  $[A_{\text{sc}}(p)]^{-1}$  of Bi-2212, Y-123 and Tl-2201. The linear fit extrapolates to zero at  $p \simeq 0.26$  close to  $p_{\text{sc}2}$ .

OPT92, Bi-OPT94 and Bi-OPT96 are at 8.3, 7.3, and 9.0  $k_B T_c$ , respectively, with  $T_c$  and the  $B_{2g}$  peaks staying pinned with an accuracy of a few percent. In spite of these differences the normal and superconducting spectra still merge in the same range of 11.5 – 13.5  $k_B T_c$  just as in  $B_{2g}$  symmetry.

In order to make a connection to previous work we plot the peak energies  $\Omega_{\text{peak}}(p)$  for  $B_{1g}$  and  $B_{2g}$  symmetry in Fig. 4 (a). For the range  $0.10 \leq p \leq 0.16$  and

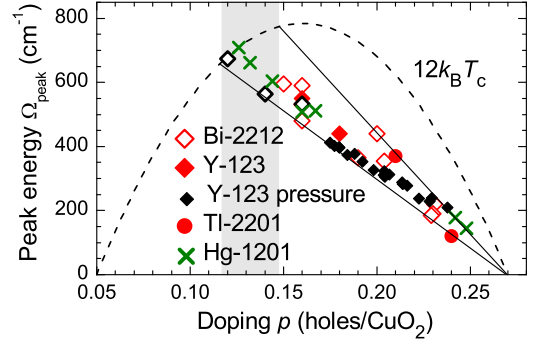


FIG. 5: (Color online) Energy and doping range of the superconductivity-induced  $B_{1g}$  features for materials as indicated. The dashed line corresponds to  $12k_B T_c$ , where superconducting and normal spectra approximately merge as indicated in Figs. 1, 2, and 3. The black filled diamonds represent results of Raman measurements at high pressure (Ref. 40) with the doping calculated as described there (see text). All points for Hg-1201 and 3 points for Bi-2212 ( $p = 0.12, 0.14, 0.16$ ; black open diamonds) are taken from Ref. 63. The lower (upper) envelope roughly corresponds to samples with lower (higher) defect concentration and/or internal strain. The shaded area indicates the doping range where the differences between normal and superconducting spectra fade away as observed for practically all homogeneous systems<sup>10,14–16,37,52,62,63</sup>.

for  $p \simeq 0.24$  we added data from Ref. 63 to get a more complete picture. Also shown are  $2\Delta_0(p) = 9.3k_B T_c(p)$  and a linear fit to  $\Omega_{\text{peak}}^{B_{1g}}(p)$ . Here and throughout the paper  $T_c(p)/T_c^{\text{max}} = 1 - 82.6(p - 0.16)^2$  as given by Presland *et al.*<sup>48</sup> with  $T_c^{\text{max}} = 94$  K. Clearly,  $\Omega_{\text{peak}}^{B_{1g}} = 1294[1 - p/p_0] \text{ cm}^{-1}$  with  $p_0 = (0.275 \pm 0.020)$  is unrelated to  $2\Delta_0(p)$  but extrapolates linearly to the upper critical doping  $p_{\text{sc}2} \simeq 0.27$  terminating the superconducting dome. In contrast, the peak energies in  $B_{2g}$  symmetry scale approximately as  $1.3\Delta_0(p)$  where the peak is expected if the lowest order  $d$ -wave variation of the gap according to Eq. (1) applies<sup>42</sup>. We note that the effect of impurities on the  $B_{1g}$  peak energies would be larger than in  $B_{2g}$  symmetry with  $\Omega_{\text{peak}}^{B_{1g}} \simeq 2\Delta_0 + \Gamma(\Gamma/\Delta_0)$ , where  $\Gamma$  is the impurity scattering rate,<sup>67</sup> and just opposite to what is observed.

Now the amplitudes  $A_{\text{sc}}(p)$  in  $B_{1g}$  and  $B_{2g}$  symmetry will be analyzed. For doping levels  $p \leq 0.22$  Blanc and coworkers<sup>38</sup> presented a similar analysis for the  $B_{1g}$  spectra and interpreted their results in terms of a gradual loss of coherence close to the  $(\pi, 0)$  point of the BZ inspired by ARPES results<sup>73</sup> and the discussion of the superfluid density<sup>6,70</sup>. We observe that at doping levels above approximately 0.20 the intensity does not saturate as would be expected from the superfluid density but rather increases at an even higher rate than below 0.20. In Fig. 4 (b) we compile results for  $A_{\text{sc}}(p)$  from the present study on Bi-2212 and Tl-2201 and our earlier results in Y-123<sup>10,14,53</sup> and Bi-2212<sup>10,14,37</sup> with all am-



plitudes given in absolute units. (Data from other than our own measurements are not included since intensities from different laboratories cannot be compared.) The differences between Y-123 and Bi-2212 are small, indicating little individual variation for these two double-layer compounds and little influence of resonantly enhanced scattering with excitation in the visible spectral range<sup>37</sup>. For  $B_{2g}$  symmetry  $A_{sc}(p)$  is practically doping independent with an average close to 1 count  $(\text{mW s})^{-1}$ . The variations of order  $\pm 50\%$  between individual samples (even at similar doping levels) not only reflect impurity effects<sup>67</sup> but also variations of the overall cross section which are not fully understood yet. Similar sample-dependent changes are also observed in  $B_{1g}$  symmetry. However, the large basis of results allows us to derive two significant trends: (i) below  $p \simeq 0.13$ ,  $A_{sc}(p)$ , i.e. any superconductivity-induced spectral change, vanishes in  $B_{1g}$  symmetry [cf. Fig. 1 (a)] in accordance with earlier studies<sup>10,14,16,38,62,63</sup>. This coincides with the rapid decrease of the coherence peaks in tunneling<sup>27,33,79</sup> and in ARPES at the anti-nodal Fermi surface crossing<sup>9,26,31,32</sup>. (ii) In Bi-2212 for  $p > 0.18$   $A_{sc}(p)$  increases to a degree which has not been appreciated yet. The trend is fully corroborated for samples from one source when studied at various doping levels (see, e.g., Ref. 37). The points from Y-123 and Tl-2201 follow the same trend although issues of oxygen order in the chains lead to strong intensity variations for optimally doped vs fully oxygenated Y-123. If we plot  $[A_{sc}(p)]^{-1}$  [Fig. 4 (c)] we find a divergence point at  $0.26 \pm 0.03$  close to  $p_{sc2} = 0.27$  where superconductivity disappears (or appears, depending on the point of view).

#### D. Comparison with other results

Before discussing possible explanations of the increasingly strong  $B_{1g}$  mode, it seems instructive to have a closer look at its range of existence and to compare the doping dependence of the Raman results in general with ARPES and tunneling results. We first note that superconductivity-induced features in  $B_{1g}$  symmetry exist only in a well-defined range of the doping-energy plane as shown in Fig. 5. For the extremal doping levels ( $p \leq 0.16$  and  $p \sim 0.24$ ) data from Ref. 63 are included. In addition, the peak energies derived from the superconducting  $B_{1g}$  spectra at high pressure are shown<sup>40</sup>. The pressure  $P$  in the range up to 22.3 GPa is converted into a doping level  $p$  using the linear relation  $p = 0.1750 + 0.0028P$  as derived by Goncharov and Struzhkin<sup>40</sup>.

First, there are no  $B_{1g}$  modes at energies in excess of approximately  $12 k_B T_c(p)$  where the normal and superconducting spectra are found to merge for those doping levels, where differences can be observed (see Fig. 1). Secondly, none of the peak energies is found outside the range given by the two straight lines obeying  $\Omega_L(p) = (1155 - 4279p) \text{ cm}^{-1}$  and  $\Omega_U(p) = (1715 -$

$6352p) \text{ cm}^{-1}$  with L (U) denoting the lower (upper) envelope. Although determined by several data points (4-6 each) the lines should be considered more a guide to the eye rather than a fit. Given this qualification it is nevertheless remarkable that both of them extrapolate to the upper critical doping  $p_{sc2}$  in a similar (approximate) fashion as the linear fit to our own data shown in Fig. 4. Up to approximately 10 GPa the peak energies derived from the pressure experiment follow  $\Omega_L(p)$  then the doping dependence slightly decreases. As possible reasons, the pressure does not directly correspond to doping or the coefficient may decrease slightly for higher pressure. For instance, a reduction by 20% aligns the peaks with  $\Omega_L(p)$ . The shaded range in which, depending on the individual sample, the superconductivity-induced features fade away is approximately determined by the intersection points of  $\Omega_{L(U)}(p)$  with the energy  $12 k_B T_c(p)$  as expected. It is reemphasized that no changes in the shape and the intensity of the  $B_{2g}$  spectra can be observed in this range. Thirdly,  $\Omega_{\text{peak}}^{B_{1g}}(p)$  decreases faster than  $T_c(p)$  for  $0.16 < p \lesssim 0.23$  as demonstrated clearly in the pressure experiment, where the ratio  $\Omega_{\text{peak}}^{B_{1g}}(P)/T_c(P)$  decreases from 6.2 to 4.2 in the range 0 to 22.3 GPa corresponding to  $p = 0.175$  and  $0.24$ <sup>40</sup>. We consider this observation along with the sample dependence at fixed doping as one indication that the mode in the  $B_{1g}$  Raman spectrum is not directly linked to the superconducting gap.

A compilation of tunneling, ARPES, and Raman results is shown in Fig. 6. In the entire overlapping doping range ( $0.07 \leq p \leq 0.19$ ),  $\Omega_{\text{peak}}^{B_{2g}}(p)$  is very similar to the quantity  $2\Delta_{sc}$  derived from the nodal part of the single-particle gap as measured by ARPES.<sup>9</sup>  $2\Delta_{sc}$  corresponds to  $2\Delta_0$  times the fraction of the Fermi surface  $\ell_{\text{arc}}/\ell_{\text{FS}}$  on which the observed gap follows the lowest-order  $d$ -wave variation as given in Eq. (1). The definitions of  $\Delta_0$  and the arc length  $\ell_{\text{arc}}$  are shown schematically in the upper left inset of Fig. 6.  $\ell_{\text{FS}}$  is the entire length of the Fermi surface. Obviously,  $\Delta_0$  is supposed to coincide with maximal gap according to Eq. (1) while  $\Delta^*$  is larger than  $\Delta_0$ . We note that deviations from the lowest-order  $d$ -wave variation of the gap were first observed by Mesot *et al.*<sup>80</sup>. Yoshida and coworkers<sup>9</sup> interpreted their recent results in terms of two independent gaps  $\Delta_0$  and  $\Delta^*$  which we followed here. Recently, Chatterjee and collaborators<sup>81</sup> found the gap to obey the lowest-order  $d$ -wave form for all doping levels indicating that there is still no agreement among the ARPES results.

The tunneling results cover a doping range from 0.100 to 0.225. (All energy scales are doubled as in the case of ARPES.) The results of Alldredge *et al.*<sup>8</sup> and Pasupathy *et al.*<sup>82</sup> are spatially resolved STS data whereas those of Miyakawa and coworkers<sup>83</sup> are not. In the latter case the energies are derived from the peak positions. Since the STS experiments yield location dependent  $I - V$  curves there are various approaches to extract a gap as discussed in the respective publications. For the comparison here, we took the maximum of the gap distribution from

Ref. 82 and, from Alldredge *et al.*<sup>8</sup>, the gap  $\Delta_1$  derived from a  $d$ -wave model with impurities. The STS energies are slightly higher but close to those without spatial resolution. Two doping ranges can be distinguished: (i) For  $p < 0.13$ , the tunneling data follow the line  $6k_B T^*$ . Hence, they are close to those of the anti-nodal gap  $\Delta^*$  which, in the doping range,  $0.09 \leq p \leq 0.12$  are slightly above  $\Delta_0$  as derived from ARPES. (ii) For  $p \geq 0.15$ , the tunneling results are close to or slightly above  $9.3k_B T_c(p)$  except for the point at  $p = 0.225$  which is approximately at  $7k_B T_c$ .

It appears that more than one half of the ARPES data for  $\Delta_0$  and the tunneling results for  $p \geq 0.15$  follow the line  $9.3k_B T_c(p)$ , as derived from Raman scattering in  $B_{2g}$  symmetry, to within the experimental resolution which, for Bi-2212, can be estimated for ARPES close to  $p = 0.19$  and for tunneling close to  $p = 0.21$  to be of the order of  $\pm 10\%$ . Importantly, the ARPES results for  $\Delta_0$  follow  $9.3k_B T_c(p)$  not only close to optimal doping,  $0.16 \pm 0.04$ , but also for  $p \simeq 0.07$ . Major discrepancies are observed in the range  $0.08 \leq p \leq 0.11$  and at 0.22. The ARPES results on Tl-2201,<sup>61</sup> although determined in a different way and with somewhat larger error bars, corroborate those on Bi-2212 and extend the doping to  $p = 0.25$ . In conclusion, for  $0.16 \leq p \leq 0.21$ , the analysis of more recent tunneling, ARPES, and Raman results reveals a significant difference between the maximal gap  $\Delta_0(p)$  and the  $B_{1g}$  peak energy  $\Omega_{\text{peak}}^{B_{1g}}(p)$  but consistent values for  $\Delta_0(p)$  if the  $B_{2g}$  Raman spectra are used. The agreement includes even wider doping ranges if  $B_{2g}$  Raman and ARPES data are compared.

Since the glass is apparently more than half full (from the viewpoint of the Raman results) we point out possible reasons for the discrepancies. Most importantly, only in the case of clean isotropic BCS superconductor can the gap be measured directly. In all other cases the gap parameter or its maximum are derived quantities. Raman scattering measures projections of a coherent superposition of the normal and the anomalous part of the electron's Nambu-Green function,  $G_{1,1}$  and  $G_{1,2}$ , respectively,<sup>42-45</sup> while ARPES (and tunneling spectroscopy) essentially observe  $G_{1,1}$  only. Therefore, Raman experiments predominantly see the condensate similarly to, for instance, Andreev reflection or optical spectroscopy, and there are good reasons to conclude that the Raman results are closer to the properties of the condensed electrons than single-particle methods. Nevertheless, prominent features in  $G_{1,1}$  and  $G_{1,2}$  are expected at similar energies at least in the weak coupling approximation. However, in the case of underdoped cuprates with additional instabilities next to superconductivity or, more generally, in strongly interacting anisotropic systems, this assumption is unlikely to hold. For example, there may be an interrelation between the superconducting gap  $\Delta_{\mathbf{k}}$  and the pseudogap  $\Delta_{\mathbf{k}}^*$  having different dependences on momentum, and the two scales cannot sufficiently be disentangled by only analyzing  $G_{1,1}$ . This may lead to an overestimation of  $\Delta_0$  if  $\Delta^*$  is larger. Similarly,

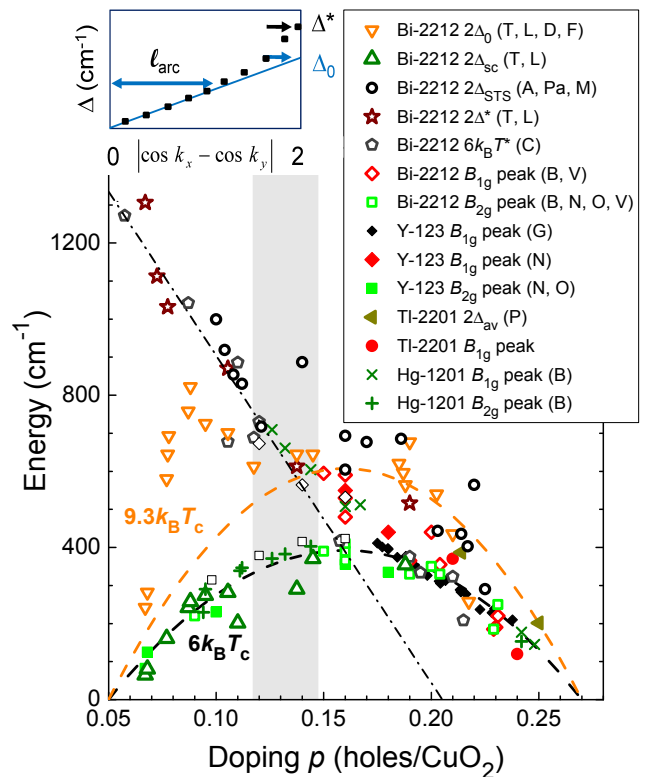


FIG. 6: (Color online) Comparison of tunneling, ARPES, and Raman experiments with the tunneling and ARPES energy scales doubled. All open symbols correspond to Bi-2212. Tunneling results are plotted as open circles. They are compiled from Refs. 8,82,83 and labeled as A,Pa,M, respectively. The stars represent the leading edge midpoints (LEM) measured above  $T_c$  at the anti-nodal Fermi surface crossing and correspond to the pseudogap  $2\Delta^*$  (T, Ref. 31 and L, Ref. 32). The pentagons correspond to the crossover temperatures  $T^*$  (C, Ref. 29) and are plotted in energy units as  $6k_B T^*$ . The down-pointing triangles represent the gap  $2\Delta_0$  derived from ARPES as sketched in the inset according to Ref. 9. The data are compiled from Refs. 31,32,84,85. The up-pointing triangles represent  $2\Delta_{\text{sc}} = 2\Delta_0 l_{\text{arc}}/l_{\text{FS}}$ , where  $l_{\text{arc}}$  is defined in the inset and  $l_{\text{FS}}$  is the full Fermi surface length (T, Ref. 31 and L, Ref. 32). On the overdoped side, ARPES results for Tl-2201 ( $p = 0.21$  and 0.25; triangles pointing to the left) are included (P, Ref. 61). Here, the averages of the peak position and the LEM are shown, and the error is on the order of  $\pm 20\%$ . The Raman data (approximate peak energies) correspond to those of the previous figures (N, O, V, B, Refs. 10,14,54,63). They are similar to within  $\pm 20\%$  to those of earlier publications as summarized in Ref. 42 and, more recently, Ref. 63. The high-pressure results are taken from Ref. 40 (G). The dash-dotted line is a fit to  $6k_B T^*$  and  $2\Delta^*$  in the range  $p \leq 0.16$ .

superconducting fluctuations as reported recently<sup>86</sup> may have an impact on  $G_{1,1}$ . In the case of the Raman response,  $G_{1,2}$  projects mainly the coherent part below  $T_c$  and  $\Delta^*$  loses influence. The solution to this problem is in fact more a theoretical challenge than an experimental one. For the case where the pseudogap results from

a charge-ordering instability there exist already some studies<sup>87–90</sup> but the way  $\Delta_0$  should be extracted from the single-particle spectra needs to be worked out further. This holds also true for other type of excitations which can be at the origin of the pseudogap or appear in the pseudogap range such as magnetic order<sup>91–94</sup>. At high doping,  $p \simeq 0.22$ ,  $2\Delta_0$  as derived from ARPES is smaller than  $9.3 k_B T_c(p)$  and appears in the same range of energies as the  $B_{1g}$  peaks. The reduction of the energy below  $9.3 k_B T_c(p)$  may start already at 0.21 and is probably related to the presence of the second energy scale observed in the  $B_{1g}$  Raman data which starts close to  $2\Delta^*(p)$  near optimal doping and then drops faster than  $T_c$ .

$2\Delta^*(p)$ , defined as the leading-edge midpoint (LEM) of the energy distribution curve at  $k_F$  close to the antinodal point<sup>9</sup>, coincides with  $\Omega_{\text{peak}}^{B_{1g}}(p)$  in the range  $0.12 \leq p \leq 0.19$  to within the experimental accuracy. Note that  $2\Delta^*(p = 0.19)$  is rather high and close to  $2\Delta_0$ . The crossover temperature  $T^*(p)$  below which the pseudogap is observed<sup>29</sup> coincides with  $\Omega_{\text{peak}}^{B_{1g}}(p)$  in an even wider doping range if it is converted into energies as  $6k_B T^*$ . Hence, the energy scale  $2\Delta^*(p)$  and to some extent also  $T^*(p)$ , both associated with the pseudogap, have a similar functional dependence on  $p$  as  $\Omega_{\text{peak}}^{B_{1g}}(p)$ . However, two important differences should be noticed: (i) The pseudogap becomes unobservable above  $p \simeq 0.21$ <sup>9,29</sup> whereas the  $B_{1g}$  Raman peak gains most of its intensity for  $p \geq 0.20$ . (ii) The doping dependences of  $2\Delta^*(p)$  and  $6k_B T^*$  change in the range where the Raman peak disappears. The low-doping part of both extrapolates to  $p = 0.205 \pm 0.020$  as indicated by a dash-dotted line in Fig. 6 with the functional dependence being represented by  $2\Delta^*(p) = (1765 - 8606p) \text{ cm}^{-1}$ .  $p \sim 0.2$  is usually identified with  $p^*$  and considered a quantum critical point<sup>5,6,95,96</sup>. We note that the onset-temperatures of the hidden magnetic order<sup>91,92</sup> and of the Kerr rotation<sup>94</sup> as well as the charge-ordering tendencies observed in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ <sup>68,97</sup> extrapolate to zero at a similar hole concentration  $0.18 < p < 0.20$  (for an interpretation see Ref. 96). Both doping levels are significantly below  $p_0 \sim 0.27$  derived from  $\Omega_{\text{peak}}^{B_{1g}}(p)$ . Although these general trends have been known for a while the details described here were not yet appreciated. However, we regard them relevant for the clarification of the  $B_{1g}$  Raman response.

It is particularly unexpected that the thermal conductivity as a typical low-energy transport probe yields a maximal gap  $\Delta_0(p)$  with a functional doping dependence close to that of  $\Delta^*(p)$  or  $T^*(p)$ <sup>19,21</sup>. Here,  $\Delta_0(p)$  is determined from the ratio of the perpendicular to the tangential electron velocities,  $v_F/v_2$ , at the Fermi surface<sup>98</sup>. Part of the puzzle may be buried in the derivation of  $v_F$  from the dispersion  $\epsilon_{\mathbf{k}}$  as seen by ARPES. For a long time  $v_F$  was believed to be universal and close to  $1.5 \text{ eV \AA}^{-1}$  ( $2 \cdot 10^7 \text{ cm s}^{-1}$ )<sup>99</sup>. In a recent high-resolution experiment, however,  $v_F$  was found to be smaller and to depend on doping<sup>100</sup>. Another part depends on the

definition of the superconducting gap and, in particular, its momentum dependence. Finally, if the nodal quasi-particle weight decreases<sup>69</sup> or if part of the carriers do not contribute any further to the heat conduction at low doping the gap appears to be too large. In any case the proper interrelation and interpretation of the various experiments and energy scales cannot be finally solved here. From the viewpoint of Raman scattering being sensitive in a spectral range above  $\Omega \simeq 15 \text{ cm}^{-1}$  (2 meV) one finds an energy which scales with  $T_c$  and another one with a more linear doping dependence which we focus on now.

### E. Origin of the $B_{1g}$ mode

Given the negligible doping dependence of  $f(\mathbf{k})$  and  $2\Delta_0/k_B T_c$  as particularly clearly demonstrated by the scaling of the  $B_{2g}$  response [Fig. 3 (a) and (b)], the variation of  $\Omega_{\text{peak}}^{B_{1g}}(p)/k_B T_c$  by a factor of two, and the tendency of spectral weight  $A_{\text{sc}}^{B_{1g}}(p)$  to diverge, it is hard to identify the  $B_{1g}$  maximum with  $\Delta_0$ . What are the alternatives?

An explanation in terms of an exciton-like bound state below  $2\Delta_0$  as described first by Bardasis and Schrieffer<sup>101</sup> has been proposed recently<sup>102,103</sup>. It is predicted that a  $\delta$ -like mode appears below the maximal gap  $2\Delta_0$  at  $\Omega_{\text{peak}}^{B_{1g}} = 2\Delta_0 - E_b$  with the binding energy  $E_b$  and the intensity increasing simultaneously with coupling strength  $\alpha_L^2$ <sup>104</sup>. In Fig. 7 we plot the dependence of the spectral weight of the in-gap mode on the bound state energy  $E_b$  given in units of  $k_B T_c$ ,  $E_b^0 = E_b/k_B T_c = (2\Delta_0 - \Omega_{\text{peak}}^{B_{1g}})/k_B T_c$ .  $E_b$  is the binding energy of the ‘‘Cooperon’’, an exciton of two electrons bound by particle-particle interactions in a channel  $L$  orthogonal to the pairing channel<sup>45,104–106</sup>. This representation demonstrates that the split-off is very small at optimal doping and increases towards higher doping levels. At first glance, the energy and intensity variations predicted on the basis of a spin-fluctuation model are similar to those observed here with a simultaneous increase of both amplitude and split-off below  $2\Delta_0$ . It is particularly luring that this type of plot partially compensates for the differences between the samples leading to an improved linear scaling, the origin of which has to remain unexplained. However, the doping dependence of the  $B_{1g}$  Raman mode is just opposite to what one expects for the spin channel. Similar arguments apply for a bound state induced by charge ordering<sup>89</sup> since the coupling for both spin and charge instabilities increases towards low, not high, doping. Similarly, the electron-phonon coupling is expected to increase towards lower carrier concentration due to the reduced screening<sup>107</sup> also in the presence of correlations<sup>108,109</sup>. At present we are not aware of an interaction with dramatically increasing coupling strength towards high doping.

Alternatively, band structure effects may play a role. A significant increase of the  $B_{1g}$  intensity is expected if the

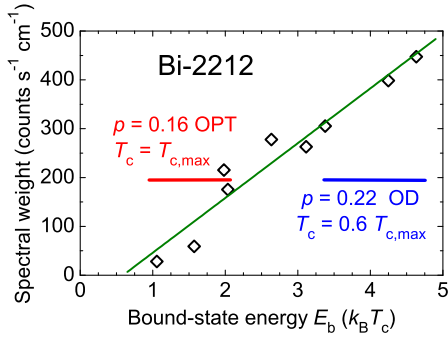


FIG. 7: (Color online) Dependence of the bound state’s spectral weight as a function of  $E_b^0 = (2\Delta_0 - \Omega_{\text{peak}}^{B_{1g}})/k_B T_c$  of Bi-2212. The ranges close to optimal doping ( $p = 0.16 \pm 0.01$ ) and at high doping ( $p = 0.22 \pm 0.01$ ) are indicated. Here, the spectral weight is that of the in-gap mode alone after subtracting the weak-coupling pair-breaking response from the full response. This can be directly derived from the data shown in Fig. 3 (c).

van Hove singularity at  $(\pi, 0)$  approaches and crosses the Fermi level  $E_F$  since the Raman cross section is proportional to the density of states at  $E_F$ . In LSCO and probably also in Bi-2212 such a crossing is indeed observed<sup>110</sup>. However, Tl-2201 has a hole-like Fermi surface at all doping levels as derived from both ARPES<sup>71</sup> and quantum oscillations<sup>111</sup>, and the data on Tl-2201 do not show significant differences to those of Bi-2212. Rather, the doping dependence on the overdoped side turns out to be quite universal<sup>68</sup> and the influence of the van Hove singularity on the superconducting Raman spectra is rather weak<sup>65</sup>. Finally, the quite complicated multi-sheeted Fermi surface of Y-123 seems to have only little influence on the spectra in the superconducting state. Hence, the synopsis of the results in Bi-2212, Y-123, and Tl-2201 allows us to conclude that band structure effects can be excluded as the origin of the increasing intensity at doping levels above  $p = 0.18$ . Similarly, the quasiparticle weight saturates above 0.20 towards the mean-field expectation and there is no indication, that the superfluid density increases further for  $p > 0.19$  (Refs. 70,112).

Since these more traditional possibilities fail to provide a qualitatively correct description of the experiments we explore a scenario which rests on the unconventional evolution with doping of the  $B_{1g}$  intensity. If individual variations between the samples are neglected,  $A_{\text{sc}}^{B_{1g}}(p)$  diverges approximately as

$$A_{\text{sc}}^{B_{1g}}(p) \propto [1 - p/p_{\text{sc}2}]^{-1}. \quad (3)$$

Although the quasiparticles at the anti-node become substantially sharper upon overdoping, such as observed by ARPES particularly in the case of Tl-2201 at  $p = 0.25$ <sup>71</sup>, the evolution of the  $B_{1g}$  Raman response can hardly be explained in this way. If this were the case, the  $B_{1g}$  maximum would just become narrower while conserving the integrated area.

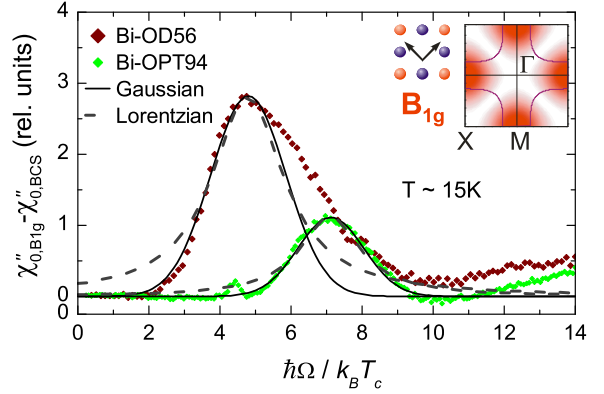


FIG. 8: (Color online) Analysis of the spectral shape of the extra contribution to the superconducting  $B_{1g}$  response in Bi-2212. The spectra are derived from those in Fig. 3 (c) by subtracting the weak-coupling result [dashed line in Fig. 3 (c)]. The solid line represents a Gaussian, the dashed one a Lorentzian. Except for the highest doping level of 0.23 all spectra are similarly well-fit by Gaussians as the one at optimal doping.

The observed intensity increase along with the reduction of  $\Omega_{\text{peak}}^{B_{1g}}(p)$  [Fig. 4 (a)] is instead more compatible with the behavior of a Goldstone mode appearing when a continuous symmetry is broken<sup>113,114</sup>. Such a scenario was actually suggested to be a possible explanation of the  $\pi$ -resonance at low doping in the SO(5) scenario<sup>115</sup>. Qualitatively, the  $B_{1g}$  Raman mode exhibits an analogous doping dependence towards  $p_{\text{sc}2}$ . In fact, we not only find  $A_{\text{sc}}^{B_{1g}}(p)$  to diverge at  $p = 0.26 \pm 0.03$  but also  $\Omega_{\text{peak}}^{B_{1g}}(p)$  to extrapolate linearly to zero at  $p = 0.27 \pm 0.02$  as expected for a symmetry-breaking mode.

In this scenario the  $B_{1g}$  spectrum is a superposition of the weak coupling pair-breaking feature (having, however, a substantially increased  $2\Delta_0/k_B T_c$  ratio) and an additional mode with  $B_{1g}$  symmetry originating from a broken continuous symmetry. This would, in a natural way, explain the strong energy dependence of the transient amplitude of the  $B_{1g}$  pair-breaking feature observed in the time domain<sup>116</sup> which was in fact interpreted in terms of a two-component response. The superposition can be explored quantitatively for Bi-2212 since the weak-coupling BSC prediction can directly be subtracted from the spectra as shown in Fig. 8. The resulting peak has a purely Gaussian lineshape for  $0.16 \leq p \leq 0.21$ . For higher doping levels deviations on the high-energy side are observed resulting either from inhomogeneities of the doping level which are very likely for  $p > 0.21$  or from an interaction of the mode with the in-gap states. In none of the cases is a Lorentzian line shape compatible with the data. Usually, a Gaussian is indicative of inhomogeneous broadening rather than a finite lifetime. The additional response is then a superposition of lines at statistically distributed energies with relatively narrow if not resolution-limited width. This would trace back the

strong sample dependences to strain fields from disorder or insufficient relaxation of the structure<sup>46</sup>. In fact, the pressure experiments directly show the susceptibility of  $\Omega_{\text{peak}}^{B_{1g}}/k_B T_c$  to stress<sup>40</sup>.

The microscopic origin remains elusive. Light scattering from a spin-density modulation with  $\mathbf{q} = (\pi, \pi)$  or from the celebrated  $\pi$ -resonance<sup>117-119</sup> would not appear in  $B_{1g}$  but, rather, in  $A_{1g}$  symmetry<sup>120,121</sup> and is, therefore, unlikely to be the origin of the  $B_{1g}$  mode. Similar symmetry arguments apply for the weakly dispersive magnetic modes which were discovered recently below  $T^*$  (Ref. 93,122) and interpreted in terms of loop currents<sup>123,124</sup>. Scattering on one of these excitations would actually appear in  $A_{2g}$  symmetry because of their chiral character.<sup>123</sup> Combined neutron and Raman scattering experiments on  $\text{HgBa}_2\text{CuO}_4$  could help clarifying this issue. A Pomeranchuk instability of the Fermi surface<sup>125,126</sup> or spin and/or charge ordering fluctuations with ordering vector  $\mathbf{Q} = (0.2\pi, 0)$ <sup>127,128</sup> would have the proper symmetry. However, only the coupling of two fluctuations would guarantee to satisfy the  $q = 0$  selection rule of light scattering in the normal and in the superconducting state. Recently, the direct coupling between the charge density wave and the superconducting gap has been studied by ARPES.<sup>90</sup> This could be a possible solution for doping levels below  $p \sim 0.2$  as pointed out already by Benfatto and coworkers<sup>88</sup> and Zeyher and Greco<sup>89</sup> whereas higher doping levels with very short range order have not been explored yet. Finally, the mode could be a critical mode of a nematic phase being established at the onset of superconductivity at  $p_{\text{sc}2} \simeq 0.27$ <sup>129,130</sup>. For symmetry reasons charge-ordering fluctuation or electronic nematic order are the most probable candidates. While indications of charge order were observed in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  (Ref. 127) their relevance for the high- $T_c$  compounds is not clear. In all cases, the mode is closely related to superconductivity without, however, being linked to the energy gap.

The experiments suggest a smooth transition between the energies  $\Omega_{\text{peak}}^{B_{1g}}$  and  $2\Delta^*$  close to optimal doping. Whether or not the similarity of the energy scales is just accidental or indicative of a relationship between the pseudogap and the putative symmetry breaking occurring at  $p_{\text{sc}2}$  is among the tantalizing open questions highlighted by this study. Yet, the present theoretical understanding is too limited to arrive at an answer.

## V. CONCLUSIONS

We have measured and analyzed the superconducting Raman response in Bi-2212, Y-123 and Tl-2201 single crystals. The present results are compared to existing data from our own work and from the literature. In Bi-2212 close to optimal doping the material dependence of the  $B_{1g}$  spectra was scrutinized revealing a significant variation with the samples' origin and history. Tl-2201 and Y-123 were used to get insight into

more extreme doping levels on the over- and under-doped sides. Superconductivity-induced features could be observed in  $B_{1g}$  and  $B_{2g}$  symmetries in the doping ranges  $0.15 \leq p \leq 0.24$  and  $0.07 \leq p \leq 0.23$ , respectively.

The scaling with the experimentally determined transition temperature  $T_c$  of the *entire*  $B_{2g}$  pair-breaking spectra, which by and large measure the condensate, is one of the central results of this publication and pins down the superconducting gap's momentum dependence to be close to the lowest order  $d_{x^2-y^2}$  form in the doping range  $0.07 \leq p \leq 0.23$ . The gap ratio  $2\Delta_0/k_B T_c \simeq 9.3$  is much higher than the BCS weak-coupling results and is found to not depend on doping. The scaling of not only the peak frequencies but also the complete  $B_{2g}$  spectra and the weak doping dependence of the overall intensity do not support a compensation effect between the gap magnitude and the loss of coherence on the anti-nodal parts of the Fermi surface as proposed recently<sup>63</sup>. Rather, the Raman results suggest that the near-nodal electronic states of the condensate exhibit BCS-type properties and seem to be protected by symmetry. Yet, there is no convincing explanation why the light scattering response does not reflect the reduction of the nodal quasiparticle weight  $Z_N$  observed by ARPES<sup>69</sup>. Similarly, there is no conclusive answer as to why the heat transport as a typical low-energy probe reflects the large gap following  $(1 - p/p_{\text{sc}2})$ <sup>19,20</sup> and why the gap derived from the single particle spectra<sup>9,100</sup> has a quite complicated doping dependence with a maximum close to  $p = 0.09$ . More generally, there remain considerable discrepancies between different methods probing  $\Delta_{\mathbf{k}}$  which apparently go beyond the conventional understanding, even in the limit of strong coupling, and need further attention.

The variations of energy and amplitude of the superconductivity-induced  $B_{1g}$  spectra cannot originate from a doping dependence of the gap beyond the approximate scaling with  $T_c$  of  $\Delta_0$ , since there should also be an influence on the  $B_{2g}$  spectra. For  $p > 0.16$  we are apparently dealing with a mode of well defined  $B_{1g}$  symmetry (typical for a collective mode) rather than a projection of the gap as in  $B_{2g}$  symmetry. We speculate that the mode indicates a broken continuous symmetry at the onset point of superconductivity at  $p_{\text{sc}2} \simeq 0.27$ . Possible candidates are fluctuating spin and charge textures with finite  $q$  (Refs. 128,131), a Pomeranchuk instability of the Fermi surface<sup>125</sup> or nematic charge order at  $q = 0$  (Ref. 129). In all cases the damping is reduced for energies smaller than  $2\Delta_0$ . Possible signatures of antiferromagnetism or orbital currents are expected in  $A_{1g}$  and  $A_{2g}$  symmetry, respectively, and are therefore unlikely to be an explanation of the  $B_{1g}$  Raman response. If the coincidence of the energies  $\Omega_{\text{peak}}^{B_{1g}}$  and  $2\Delta^*$  occurring close to optimal doping is not just accidental the results indicate a relationship between the broken symmetry and the pseudogap. It is an intriguing possibility that the  $B_{1g}$  mode is related to the fluctuation spectrum at the origin of superconductivity.

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