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Confinement and correlation effects in the Xe@C_{60} generalized oscillator strengths

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Abstract

The impact of both confinement and electron correlation on generalized oscillator strengths (GOS's) of endohedral atoms, $A@C_{60}$, is theoretically studied choosing the Xe@C_{60} $4d$, $5s$ and $5p$ fast electron impact ionization as the case study. Calculations are performed in the transferred to the atom energy region beyond the $4d$ threshold, ω = 75–175 eV. The calculation methodology combines the plane wave Born approximation, Hartree-Fock approximation, and random phase approximation with exchange in the presence of the C_{60} confinement. The confinement is modeled by a spherical δ-function-like potential as well as by a square well potential to evaluate the effect of the finite thickness of the C_{60} cage on the Xe@C_{60} GOS's. Dramatic distortion of the $4d$, $5p$ and $5s$ GOS's by the confinement is demonstrated, compared to the free atom. Considerable contributions of multipolar transitions beyond dipole transitions in the calculated GOS's is revealed, in some instances. The vitality of accounting for electron correlation in calculation of the Xe@C_{60} $5s$ and $5p$ GOS's is shown.

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I. INTRODUCTION

Nano-objects $A@C_n$, consisting of an atom A encapsulated inside the hollow inner space of a carbon cage C_n, known as endohedral fullerenes, or endohedral atoms, or, simply, endohedrals, or confined atoms, have attracted much attention of investigators. This is because of their importance to various basic and applied sciences and technologies. To name a few, one could emphasize their significance for astrophysics [1], invention of quantum computers [2], development of unique superconductors [3, 4], cancer therapy [5], etc. Understanding of their quantum structure as well as interaction with various incoming beams of particles - photons, electrons, ions, etc. - is imperative. From a theoretical side, the problem is formidable in complexity due to its multi-faceted nature. A unique theory that solves this problem once and for all is yet to be developed. Meanwhile, with the help of simpler, physically transparent theoretical models, theorists have been unraveling most unusual aspects of $A@C_n$ confined atoms, thereby identifying the most useful experimental studies, which could be performed. Much of attention has been turned to various aspects of photo-ionization of endohedral atoms. The interested reader is referred to review papers [6, 7] as well as some recent papers [8–13] and references therein in addition to other references presented in this paper, for a detailed introduction into the subject. Many important insights into $A@C_{60}$ photoionization have been obtained on the basis of the Δ-potential [6, 7] and δ-potential [14, 15] models. In the δ-potential model the C_{60} cage is assumed to have the zero thickness and is modeled by a spherical δ-function potential $V(r) = U_0\delta(r - R_0)$ of an inner radius R_0 and depth U_0. In contrast, the Δ-potential model accounts for the finite thickness Δ of the C_{60} cage. It models the cage by a square well potential of the width Δ. One of spectacular findings, obtained on the basis of these models, has been the discovery of resonances, termed confinement resonances (CR's)[15, 16] and correlation confinement resonances (CCR's) [17], in the photoionization spectrum of an endohedral atom. CR's (also referred to as *ordinary* CR's in this paper) occur in photoionization spectra of endohedral atoms due to interference of the photoelectron waves emerging directly from the confined atom A, and those scattered off the C_{60} carbon cage. CCR's differ from these ordinary CR's in that they occur in the spectrum of an *outer* subshell of the confined atom A due to interference of transitions from this subshell with ordinary CR's emerging in *inner* shell transitions, via interchannel coupling [17]. CCR's represent a novel class of resonances that can exist neither without

confinement nor electron correlation. Both, ordinary CR's and CCR's have attracted much interest of researchers. In particular, of great importance were theoretical predictions of a dramatic distortion of the atomic Xe $4d$ giant resonance by CR's in the Xe@C_{60} $4d$ photoionization made on the basis of the δ-potential model [18], Δ-potential model [17], and time-dependent local density approximation (TDLDA) [11] calculations. This has stimulated a photoionization experiment that led to a recent experimental discovery of CR's in the Xe@C_{60}^{+} $4d$ photoionization spectrum [19]. The results obtained were in a much better agreement with the δ-potential model calculated data [18] than with those obtained in the framework of the Δ-potential model or TDLDA. (Note, beforehand, that for this reason primarily the δ-potential model is employed in the present paper study).

In contrast to photoionization cross sections of endohedral atoms, little is known about their generalized oscillator strengths (GOS's). GOS's reflect the atomic response to fast electron impact ionization. They are more complicated and informative parameters than corresponding photoionization cross sections. This is because, generally, many more multipolar transitions contribute to electron impact ionization of atoms versus primarily dipole transition contributions to the photoionization process, see, e.g., Ref. [20]. The electron spectroscopy of quantum objects thus serves as another powerful tool for the study of their structures. However, to date, GOS's of endohedral atoms were investigated only in the theoretical work [21], and there has been no associated experimental studies performed. In Ref. [21], the ionization of the innermost $1s$ subshells of endohedral He@C_{60} and Ne@C_{60} by fast electrons was chosen as the case study. Both the δ- and Δ-potential models were employed in the study. Electron correlation was not accounted for. It was shown that, much as due to photoionization, noticeable ordinary CR's emerge in GOS's of endohedral atoms as well.

The above mentioned finding on GOS's of endohedral atoms is the only known to date confinement-related-feature of their electron impact ionization. The knowledge on a possible significance of electron correlation in GOS's of A@C_{60} atoms is ultimately absent, despite its obvious importance. The present paper fills in this vacancy in one's knowledge. It advances the initial understanding of GOS's of such atoms by accounting for a mutual impact of confinement and correlation on the ionization process. With the general impetus of the successful experimental study of CR's in the Xe@C_{60} $4d$ photoionization [19], we, too, choose to explore the Xe@C_{60}, as the case study. We focus on the $4d$, $5p$ and $5s$ GOS's upon fast electron impact ionization of the confined Xe assuming that the transferred to

the atom energy exceeds the $4d$ threshold which is approximately 74 eV. This is in order to avoid dealing with the C_{60} dynamical polarization impact on an $A@C_{60}$ atom, which is known to be strong at lower energies [22–24], to simplify matters. Confinement effects are then accounted for primarily in the framework of the δ-potential model without regard for said dynamical polarization of C_{60}. The Δ-potential model is employed in some instances as well, to evaluate the finite-potential-width-impact on the $Xe@C_{60}$ GOS's. In the calculations, the plane wave Born approximation (PWBA) is used for the fast incoming and scattered electrons. A Hartree-Fock (HF) approximation in the presence of the C_{60} confinement is employed relative to the confined atom itself. This completes GOS's calculations in the zero-order approximation (omitting correlation). Electron correlation is then accounted for in the framework of the random phase approximation with exchange (RPAE) [20, 25], as the final step in the GOS's study.

II. THEORY

In this section, we provide the outline of the general approach to calculations of GOS's of free (A) and endohedral $A@C_{60}$ atoms.

In PWBA, GOS of an nl atomic subshell, $f_{nl}(q,\omega)$, is defined by [in atomic units (a.u.), but with the energy being measured in Rydbergs (Ry)] [20]

$$f_{nl}(q,\omega) = \frac{2(2\lambda+1)N_{nl}\omega}{q^2(2l+1)} \sum_{l'\lambda} |Q^{\lambda}_{nl,\epsilon' l'}(q)|^2. \tag{1}$$

Here, N_{nl} is the number of electrons in the ionizing atomic subshell nl, q is the magnitude of the transferred linear momentum to the atom upon the collision, ω and λ are the corresponding transferred energy and orbital momentum, respectively, $Q^{\lambda}_{nl,\epsilon' l'}(q)$ is a reduced matrix element for the ionization amplitude (in length-form), ϵ' is the energy of an ejected electron ($\epsilon' = \omega - I_{nl}$, I_{nl} being the nl subshell ionization potential).

In a HF approximation,

$$\begin{aligned} Q^{\lambda(\mathrm{HF})}_{nl,\epsilon' l'}(q) \;=\; & \sqrt{(2l'+1)(2\lambda+1)} \begin{pmatrix} l & l' & \lambda \\ 0 & 0 & 0 \end{pmatrix} \\ & \times \int_0^{\infty} P_{nl}(r) j_{\lambda}(qr) P_{\epsilon' l'}(r) dr. \end{aligned} \tag{2}$$

Here, $P_{nl}(r)$ and $P_{\epsilon' l'}(r)$ are radial parts of corresponding HF atomic wave-functions of the initial and final states of the atom, and $j_\lambda(qr)$ is the spherical Bessel function.

In RPAE, the equation for the GOS reduced matrix element $Q^{\lambda}_{nl,\epsilon' l'}(q)$ is more complicated due to the specific accounting for intershell coupling of the $nl \to \epsilon l'$ transition with electronic transitions from other subshells of the atom, see Eq. (10.14) in Ref. [20]:

$$\begin{aligned} Q^{\lambda(\mathrm{RPAE})}_{nl,\epsilon' l'}(q) \;=\; & Q^{\lambda(\mathrm{HF})}_{nl,\epsilon' l'}(q) \\ & + \left(\sum_{\substack{k''l'''>F,\\ k'l''\le F}} - \sum_{\substack{k'l''>F,\\ k''l'''\le F}} \right) \\ & \times \frac{U_{nl,\epsilon' l';k'l'',k''l'''} Q^{\lambda(\mathrm{RPAE})}_{k'l'',k''l'''}(q)}{\omega - \epsilon_{k''l'''} + \epsilon_{k'l''} + i\eta}. \end{aligned} \tag{3}$$

Here $kl \le F$ denotes summation over all occupied atomic states, $kl > F$ marks summation over discrete excited states including integration over continuous spectrum with the assumption of $\eta \to +0$, ϵ_{kl}'s are the HF energies of corresponding vacant or occupied atomic states, $U_{nl,\epsilon' l';n'l'',k'l'''} = (nl, \epsilon' l'|V|n'l'', k'l''') - (nl, \epsilon' l'|V|k'l''', n'l'')$ is the difference between direct and exchange Coulomb matrix elements of intershell interaction, respectively. The interested reader is referred to Ref. [20] for more details of the RPAE methodology.

We now turn to the description of GOS's of $A@C_{60}$ endohedral atoms.

Let us first employ the δ-potential model [14, 15] to account for the C_{60} confining cage. This model exploits the following two key assumptions. First, it is assumed that the size of a confined atom is much smaller than the C_{60} radius R_0, $R_0 = 6.64$ a.u. [26]. This allows one to equal the ground state energies and electronic wave-functions of the confined atom to those of the free atom. Second, it is assumed that the thickness Δ of the C_{60} cage is much smaller than the wavelength of the outgoing electron released upon ionization of the confined atom. Hence, the thickness can be disregarded at all, i.e., $\Delta \to 0$, to a good approximation. Correspondingly, one can model the C_{60} cage by the δ-function-like potential $U_\delta(r)$:

$$U_\delta(r) = -U^{(0)}_\delta \delta(r - R_0). \tag{4}$$

Here, $U^{(0)}_\delta = 0.442$ a.u. is the potential depth which was found [14, 15] by matching the calculated electron affinity (EA) of C_{60} to the known one, $EA = 2.65$ eV [26]. The confinement brought impact on the ionization process is then associated only with modification of

the wave-function $P_{\epsilon l'}(r)$ of an outgoing electron due to its scattering off the confining cage. The modification results [15] in $P_{\epsilon l'}(r)$ which differs from that of the free atom $P^{\rm free}_{\epsilon l'}(r)$ only by a multiplicative factor $D_{l'}(k)$ (k being the momentum of an outgoing electron):

$$P_{\epsilon l'}(r) = D_{l'}(k) P^{\rm free}_{\epsilon l'}(r), \tag{5}$$

where

$$D_{l'}(k) = \cos\eta_{kl'} \left[1 - \tan\eta_{kl'} \frac{G_{kl'}(R_0)}{P_{kl'}(R_0)}\right]. \tag{6}$$

Here, $G_{kl'}$ is the irregular-at-zero solution of the HF equation for the isolated atom, whereas $\eta_{kl'}$ is the additional to the free atom phase shift due to the δ-potential well:

$$\tan\eta_{kl'}(k{\rm r}) = \frac{P^2_{kl'}(R_0)}{P_{kl'}(R_0) G_{kl'}(R_0) + k/2B}. \tag{7}$$

With the help of Eq. (5), the HF matrix element for the confined atom GOS amplitude, labeled as $Q^{\lambda(\delta{\rm HF})}_{nl,\epsilon' l'}(q)$, differs [21] from that of the free atom, $Q^{\lambda({\rm free})}_{nl,l'}(q)$, Eq. (2), only by the factor $D_{kl'}$:

$$Q^{\lambda(\delta{\rm HF})}_{nl,\epsilon' l'}(q) = D_{l'}(k) Q^{\lambda({\rm free})}_{nl,\epsilon' l'}(q). \tag{8}$$

We will be referring to the described HF approximation for calculating GOS's of confined atoms as the δHF approximation; the symbol δ emphasizes that the approximation employs the δ-potential concept.

As follows from Eq. (6), the coefficient $D_l(k)$ has an oscillatory character versus k (and, hence, versus the transferred to the atom energy ω). Therefore, there are resonances - confinement resonances - emerging in the transition matrix elements for $A@C_{60}$ atoms. They translate into resonances either in their photoionization cross sections or, what is more important to us, generalized oscillator strengths [21].

In the framework of the alternative Δ-potential model, the potential $U_\delta(r)$, Eq. (4), is replaced by a short-range square-well potential $U_\Delta(r)$ of the width Δ and depth $U^{(0)}_\Delta$:

$$U_\Delta(r) = \begin{cases} -U^{(0)}_\Delta, \text{ at } R_0 - \frac{1}{2}\Delta \le r \le R_0 + \frac{1}{2}\Delta \\ 0, \qquad\qquad \text{otherwise} \end{cases} \tag{9}$$

Now, excited wavefunctions of the confined atoms $P^{\Delta}_{\epsilon l'}(r)$ are not proportional to wavefunctions $P^{\rm free}_{\epsilon l'}(r)$ of the free atom, Eq. (5). Instead, the new $P^{\Delta}_{\epsilon l'}(r)$ are to be found by a straightforward solution of a "confined" HF equation (refereed to as the ΔHF approximation), i.e., the HF equation which includes the potential $U_{\Delta}(r)$ in addition to the free atom potential. In the present work, we employ the values of $U^{(0)}_{\Delta} = 0.422$ a.u. and $\Delta = 1.25$ a.u. since they were proven [27] to result in the best possible Δ-model description of experimentally observed CR's in the Xe@C^{+}_{60} $4d$ photoionization cross section [19]. Corresponding HF GOS's amplitudes $Q^{\lambda(\Delta{\rm HF})}_{nl,\epsilon' l'}(q)$ are then calculated with the help of the thus found wavefunctions $P^{\Delta}_{\epsilon l'}(r)$.

To account for RPAE electron correlation in GOS's of A@C_{60} atoms in either of the discussed models, the standard free-atom-RPAE-equation is turned into the "confined-atom-RPAE-equation" by a straightforward replacement of all free atomic excited state wavefunctions $P^{\rm free}_{\epsilon l'}(r)$ by the above discussed wavefunctions $P^{\delta}_{\epsilon l'}(r)$ or $P^{\Delta}_{\epsilon l'}(r)$, respectively. Similar to the used δHF and ΔHF abbreviations, we refer to such confined-atom-RPAE methodology as δRPAE and ΔRPAE, respectively.

In particular, the δRPAE equation transforms into

$$\begin{aligned} Q^{\lambda(\delta{\rm RPAE})}_{nl,\epsilon' l'}(q) \;=\; & Q^{\lambda(\delta{\rm HF})}_{nl,\epsilon' l'}(q) \\ & + \left(\sum_{\substack{k' l'' \le F, \\ k'' l''' > F}} D^2_{l'''}(k'') - \sum_{\substack{k' l'' > F, \\ k'' l''' \le F}} D^2_{l''}(k') \right) \\ & \times \frac{U_{nl,\epsilon' l';k' l'',k'' l'''} Q^{\lambda(\delta{\rm RPAE})}_{k' l'',k'' l'''}(q)}{\omega - \epsilon_{k'' l'''} + \epsilon_{k' l''} + i\eta}. \end{aligned} \tag{10}$$

The intershell interaction term in the δRPAE equation (the second term on the right-hand-side of the equation) explicitly depends on the oscillatory parameter $D_l(k)$. Hence, the δRPAE approximation is capable of accounting for CCR's in the GOS's spectra. It accounts for ordinary CR's as well, owing to the term $Q^{\lambda(\delta{\rm HF})}_{nl,\epsilon' l'}(q)$ (the first term on the right-hand-side of the δRPAE equation) which itself is determined by Eq. (8).

As for the alternative ΔRPAE equation, the latter accounts for CR's and CCR's implicitly, via final-state and intermediate-state functions $P^{\Delta}_{\epsilon l}(r)$.

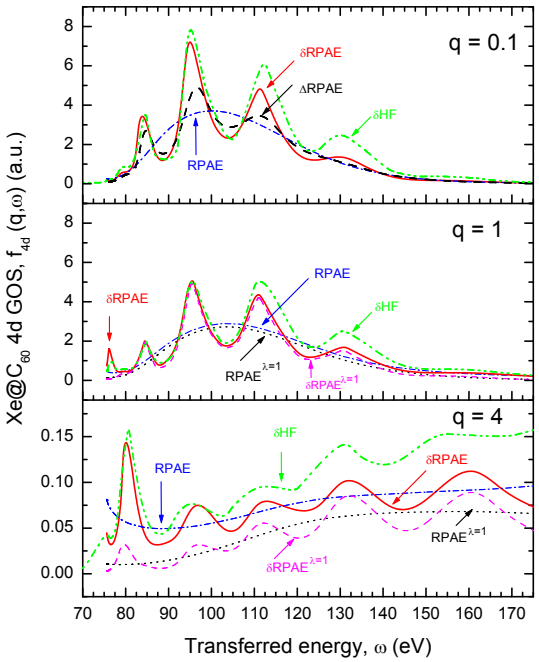

FIG. 1: (Color online) The $4d$ generalized oscillator strengths $f_{4d}(q,\omega) = \sum_\lambda f^\lambda_{4d}(q,\omega)$ upon fast electron impact ionization of Xe@C_{60} (δHF and δRPAE) and free Xe (HF and RPAE). Calculated results marked as RPAE$^{\lambda=1}$ and δRPAE$^{\lambda=1}$ relate to $f^{\lambda=1}_{4d}(q,\omega)$ due to only dipole transition contributions.

III. RESULTS AND DISCUSSION

Our δHF, ΔHF, δRPAE and ΔRPAE calculated results for the $4d$ GOS's $f_{4d}(q,\omega)$ of Xe@C_{60}, upon fast electron impact ionization, are presented in Fig. 1 along with corresponding HF and RPAE calculated data for free Xe. Calculations were performed for transferred momenta $q = 0.1$, 1, and 4 and accounted for major monopole ($\lambda = 0$), dipole ($\lambda = 1$), quadrupole ($\lambda = 3$), and octupole ($\lambda = 3$) multipolar contributions to $f_{4d}(q,\omega)$. RPAE and δRPAE calculations included dominant intershell interaction between transitions from $4d$, $5s$, and $5p$ subshells. A number of spectacular trends in $f_{4d}(q,\omega)$ versus ω, q, and λ is seen.

One of the trends is the presence of strong oscillations in the $f_{4d}(q,\omega)$ of Xe@C_{60}. They are absent in the $4d$ GOS's of free atom. Thus, the oscillations are due to the C_{60} confinement, thereby featuring the emergence of confinement resonances in the $4d$ GOS's of Xe@C_{60}.

Furthermore, interesting, relative intensities and positions of these CR's appear to be noticeably changing with increasing value of q in the whole energy region, including near threshold. As a results, the calculated $f_{4d}(q = 0.1,\omega)$ and $f_{4d}(q = 1,\omega)$ have little to do with

$f_{4d}(q=4,\omega)$. Moreover, this impact also results in that the GOS's of free Xe and, on the average, of Xe@C_{60} are found to exhibit a strong, broad resonance for $q=0.1$ and $q=1$ in contrast to $q=4$, in the energy region under discussion. To understand this, we also plotted in Fig. 1 calculated data of a trial δRPAE calculation [labeled as $f_{4d}^{\lambda=1}(q,\omega)$] accounting only for dipole contributions to the $4d$ GOS's. The comparison of the total GOS $f_{4d}(q,\omega)$ with $f_{4d}^{\lambda=1}(q,\omega)$ shows the dominant role of dipole transitions in the ionization process for smaller q's, $q=0.1$ and 1, both for free Xe and Xe@C_{60} [note, for $q=0.1$ dipole channels exceed other channels by several orders of magnitude, for which reason $f_{4d}(q,\omega)$ is undistinguished from $f_{4d}^{\lambda=1}(q,\omega)$]. The Xe dipole $4d\rightarrow\epsilon f$ transition is known to exhibit a strong resonance versus ω, known as the $4d$ giant resonance [25]. Consequently, the nature of the strong, broad resonance seen in the $4d$ GOS's at $q=0.1$ and 1 is the same as in the case of the Xe $4d$ photoionization, i.e., it is the familiar $4d$ dipole giant resonance. For a larger value of q, $q=4$, other channels beyond the dipole channel acquire considerable strengths compared to the dipole channel. This can be judged by comparing total $f_{4d}(q,\omega)$ with partial $f_{4d}^{\lambda=1}(q,\omega)$ for $q=4$ displayed in Fig. 1. This explains why $f_{4d}(q=4,\omega)$ has, in general, little in common with the $4d$ GOS's for smaller q's, as well as why $f_{4d}(q=4,\omega)$ does not exhibit the $4d$ giant resonance - dipole transitions matter little. Thus, the transferred momentum as well as multipolar impacts on the $4d$ GOS's of both free Xe and Xe@C_{60} are found to be considerable.

Moreover, Fig. 1 additionally features a varying role of multipolar contributions to GOS's near threshold. A trial calculation showed that the first resonance maximum in $4d$ GOS's near threshold is primarily due to monopole channels for $q=1$, whereas it is mainly due to octupole channels for a larger $q=4$.

Finally, it also obvious from Fig. 1 that the alternative ΔRPAE calculation of the Xe@C_{60} $4d$ GOS results in lower and yet clearly prominent intensities of emerged CR's compared to the δRPAE data; this is in line with results of the previous theoretical study [21]. Note, since GOS's of endohedral atoms have been experimentally unexplored, the question of which of the used models is most appropriate remains open.

We now proceed to the discussion of the Xe and Xe@C_{60} $5s$ GOS's. Corresponding calculated HF and RPAE (for free Xe) as well as δHF and δRPAE data for $f_{5s}(q,\omega)=\sum_\lambda f_{4d}^{\lambda}(q,\omega)$ are depicted in Fig. 2 for $q=0.1$ and 1 (the $5s$ GOS's for $q=4$ appear to be negligible compared to those for $q=0.1$ or 1, thus presenting little interest for discussion). As in the above

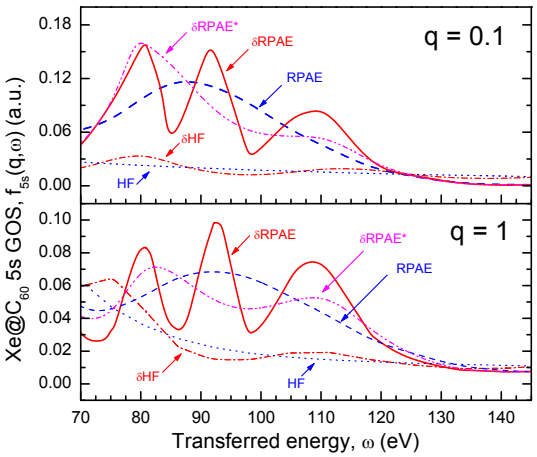

FIG. 2: (Color online) Calculated data for generalized oscillator strengths $f_{5s}(q,\omega)$ of free Xe (HF and RPAE) and Xe@C_{60} (δHF and δRPAE). δRPAE* labels the fictitious calculated data for the Xe@C_{60} $5s$ GOS's (see text).

study, calculations accounted for contributions of major monopole, dipole, quadrupole, and octupole ionization channels to $f_{5s}(q,\omega)$ as well as intershell coupling between transitions from the Xe $4d$, $5s$, and $5p$ subshells both in RPAE and δRPAE equations. The calculated data show that, as in the known case of the Xe $5s$ photoionization [25], the $5s$ GOS's of free Xe are ultimately affected by electron correlation, both for smaller and bigger values of q. This clearly follows from the comparison of corresponding HF and RPAE calculated data. The same tendency is found to preserve in endohedral Xe@C_{60} as well, cf. δHF and δRPAE calculated data. Furthermore, similar to $4d$ GOS's, the Xe@C_{60} $5s$ GOS's are found to be dramatically distorted by confinement. Indeed, the presence of the C_{60} confinement results in the emergence of three strong oscillations (resonances) in $f_{5s}(q,\omega)$ at given ω's, cf. RPAE and δRPAE calculated data. Noting that the resonance positions are about the same as the position of ordinary CR's in the $4d$ GOS's at approximately 83, 95, and 102 eV, one might be tempted to interpret the resonances in $5s$ GOS's as CCR's, i.e., being induced in $f_{5s}(q,\omega)$ by the three ordinary CR's in $4d$ ionization channels, via intershell interaction, as in the case of the Xe@C_{60} 5s photoionization [17]. This, however, would not be entirely correct. Indeed, e.g., such interpretation would fail to explain why a weak CR in $f_{4d}(q,\omega)$ at about 83 eV induces as strong CCR in $f_{5s}(q,\omega)$ as the two other stronger neighboring resonances. The actual origin of the three emerged resonances in $5s$ GOS's is thus intriguing.

To unravel the nature of the three resonances in question in $f_{5s}(q,\omega)$ of Xe@C_{60}, we per-

formed a fictitious trial δRPAE calculation of the $5s$ GOS's. There, we artificially eliminated ordinary CR's from corresponding coupling $4d$ ionization channels in δRPAE, Eq. (10). This was achieved by substituting the excited wavefunctions of the $4d$ electrons of *free* Xe instead of those of Xe@C_{60} into Eq. (10). This fictitious methodology will be referred to/labeled as δRPAE*. Corresponding δRPAE* calculated data are depicted in Fig. 2 as well. One interesting important observation is that δRPAE* intershell interaction noticeably increases CR's presented in δHF calculated $f_{5s}(q,\omega)$ for both values of q and also shifts a lower energy CR in $f_{5s}(q=1,\omega)$ from about 75 to about 83 eV. Such δRPAE* calculation, however, does not bring a third (middle) resonance in $f_{5s}(q,\omega)$. This resonance emerges only in a true δRPAE calculation (which accounts for CR's in $4d$ ionization channels) of $f_{5s}(q,\omega)$, and its position coincides with that of a middle resonance in $4d$ GOS's. Thus, our first conclusion in unraveling the nature of the three resonances in $f_{5s}(q,\omega)$ is that the middle resonance is undoubtedly correlation confinement resonance, CCR, by nature [it does not exist without simultaneous impact of intershell interaction and confinement on $f_{5s}(q,\omega)$]. What about the left and right resonances in $f_{5s}(q,\omega)$? Their nature is more complicated. Qualitatively, they exist in $f_{5s}(q,\omega)$ even without coupling with CR's in $4d$ ionization channels (see δRPAE* calculated data), i.e., the resonances uncovered by the performed δRPAE* calculation are ordinary CR's. In a rare, unique occasion, these CR's (see δRPAE* calculated data) peak at about the same energies as the left and right *ordinary* CR's in $4d$ GOS's, Fig. 1. Therefore, when the $4d$ CR's are accounted for in a true δRPAE calculation they strongly affect the two originally existing δRPAE* ordinary CR's in $f_{5s}(q,\omega)$, thereby enhancing them, via intershell interaction. Thus, the left and right resonances in true δRPAE calculated data for $f_{5s}(q,\omega)$ are the result of intershell coupling of the ordinary CR's in $4d$ channels with *existing* ordinary CR's in $5s$ ionization channels. Therefore, the left and right resonances in $f_{5s}(q,\omega)$ are neither purely ordinary CR's nor purely CCR's. Rather, they may be termed as *CR-CR-correlation-interference-resonances* - a new type of resonances which have not been met earlier, to the best of our knowledge. This interpretation, in particular, explains why the lower left resonance in $f_{5s}(q,\omega)$ is as strong as the middle CCR. To start with, it was relatively strong from the very beginning (see δRPAE* calculated data). Next, the interaction with a weaker left $4d$ CR makes this $5s$ resonance somewhat stronger, so that it now matches the middle CCR which is brought about by the strongest middle CR in the $4d$ channel. To summarize, the discussed $5s$ GOS resonance spectrum of Xe@C_{60} has neither

a purely CR nor CCR nature. Rather, it consists of one CCR and two CR-CR-correlation-interference-resonances. Note, this makes the $5s$ GOS spectrum of Xe@C_{60} be unique and different in its origin from corresponding $5s$ Xe@C_{60} photoionization spectrum [17] which consists of purely CCR's.

Of specific interest are the GOS's of $5p$ electrons in the considered ω region. Since it is well above the $5p$ ionization threshold, $f_{5p}(q,\omega)$ for Xe@C_{60} could have been expected to have negligible or very weak, at best, CR's. This seems to be in line with a theory of scattering of particles off a potential well or barrier. Indeed, at a sufficiently high energy of the outgoing electron, the corresponding coefficient of reflection off a finite potential well or barrier is small. As a result, the interference effect between the outgoing and scattered electron waves becomes weak, and so are the associated CR's. This, however, is true only in terms of an independent particle approximation. As was shown in Ref. [28], where photoionization of endofullerenes was chosen as a case study, CR's can reappear - resurrect - and be strong at the transferred to the atom energy far exceeding (by thousands of eV) the nl ionization threshold, as a general phenomenon. This will happen at transferred energies, which correspond to opening of inner-shell photoionization channels, whose intensities exceed by far the intensity of transitions from the outer subshell of the confined atom and which are strongly coupled with the innershell transitions. This is just the case with the Xe@C_{60} $5p$ ionization above the $4d$ threshold. Indeed, as known from photoionization studies [25], the Xe $5p$ ionization is affected strongly by intershell interaction with $4d$ transitions. As was found above, the Xe@C_{60} $4d$ GOS's are (a) strong and (b) have pronounced CR's, see Fig. 1. Therefore, one can predict the emergence of strong CCR's in the Xe $5p$ GOS when intershell interaction between the $5p$ and $4d$ ionization amplitudes is accounted for in δRPAE calculation, as in the above detailed case of $5s$ GOS's. Corresponding HF, RPAE, δHF and δRPAE calculated data for $f_{5p}(q,\omega)$ for $q = 0.1$ are depicted in Fig. 3. The presented results are self-explanatory. In brief, there is the only weak CR related oscillation in the δHF calculated $5p$ GOS for $q = 0.1$, in contrast to three strong resonance features present in the δRPAE calculated data. The latter are found to be induced in the $5p$ GOS by CR's in the $4d$ GOS amplitudes, i.e., they are CCR's. They are strong, in a full accordance with the above made prediction. In addition, calculated data show that, as in the known case of the free Xe $5p$ photoionization [25], the $5p$ GOS's of both free Xe and Xe@C_{60} are completely determined by intershell correlation with the $4d$ subshell. This clearly follows from the comparison between HF and

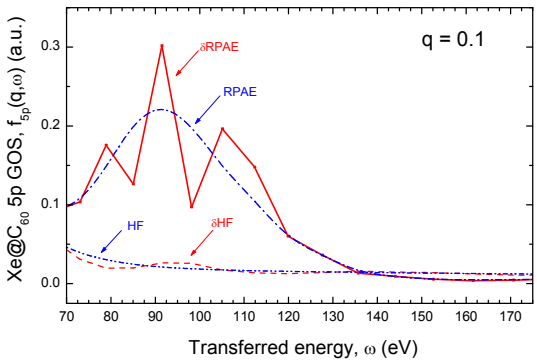

FIG. 3: (Color online) Calculated data for generalized oscillator strengths $f_{5p}(q,\omega)$ of free Xe (HF and RPAE) and Xe@C_{60} (δHF and δRPAE), as marked.

RPAE calculated data on the one hand, and δHF and δRPAE calculated data on the other hand.

IV. CONCLUSION

In the present work we focused on the study of the impact of the C_{60} confinement on the $4d$, $5s$ and $5p$ generalized oscillator strengths of Xe@C_{60}, in the energy region above the $4d$ threshold, where, in our opinion, the most interesting effects occur. We hope that the discovered impact of the transferred momentum q, electron correlation, and confinement on generalized oscillators strengths of Xe@C_{60} will challenge experimentalists to verify our predictions. Theorists, we hope, will be driven by the desire to improve the made predictions with the help of more sophisticated theories. All this would indisputably result in uncovering of a richer variety of possible effects outside of the made predictions, thereby advancing this field of endeavor.

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