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# Searches on star graphs and equivalent oracle problems

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We examine a search on a graph among a number of different kinds of objects (vertices) one of which we want to find. In a standard graph search, all of the vertices are the same, except for one, the marked vertex, and that is the one we wish to find. We examine the case in which the unmarked vertices can be of different types, so that the background against which the search is done is not uniform. We find that the search can still be successful, but the probability of success is lower than in the uniform background case, and that probability decreases with the number of types of unmarked vertices. We also show how the graph searches can be rephrased as equivalent oracle problems.

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## I. INTRODUCTION

One of the main reasons quantum walks were originally introduced into quantum information was the hope that they might provide a useful conceptual framework for the development of new quantum algorithms [1]. This hope has been fulfilled [2]-[5]. In addition, they have recently attracted considerable experimental interest. They have been realized in optical lattices [6, 7] and with photons in waveguide lattices [8]. Analogs of quantum walks, which are based on the wave nature of classical light, make use of the fact that the most important element in the behavior of a quantum walk is interference. These were proposed in [9, 10] and experimentally realized in Ref. [11]. The use of interferometers to realize quantum walks was proposed [10] and recently experimentally implemented [12]. Quantum walks with trapped ions have been theoretically described [13, 14] and experimentally realized as well [15, 16]. Finally, experiments in [17] and [18] provide a framework in which walks with more than one walker can be studied.

Here we will look at quantum walks on star graphs and equivalent oracle problems. A star graph has a single central vertex with  $N$  edges emanating from it. Each of these edges is connected to a single external vertex. Thus the graph resembles spokes extending from a hub. We shall denote the central vertex by 0 and the external vertices by 1 through  $N$ . Star graphs form a natural arena in which to study quantum searches. Some of the external vertices may have different properties from the others, and our task is to find these vertices. In a standard quantum search one object has different properties than all of the others, all of which are the same, and which constitute a kind of background. Here we want to examine searches on more complicated backgrounds. Suppose that we want to find an object with a particular property, but the other objects that we must search among are not all the same. In the case of a star graph, this could be realized by having the external vertices reflect the walking particle with one of three phases. One vertex (the marked vertex) reflects with the phase we want to find, and each of the other vertices reflects with

one of the other two phases. As we shall see, it is possible to find the marked vertex using a quantum walk, but there is a limit to the probability of successfully doing so.

We shall employ the scattering model of a quantum walk in which the particle resides on the edges of the graph and at each time step scatters at the vertices [19, 20]. Each edge has two states, corresponding to the two directions a particle on it can be going. In the case of a star graph, the edge between vertex 0 and  $j$  is associated with the state  $|0, j\rangle$ , which corresponds to the particle being on that edge and going from 0 to  $j$ , and  $|j, 0\rangle$ , which corresponds to the particle being on that edge but going from  $j$  to 0. The states corresponding to the edges form an orthonormal basis for the Hilbert space of the particle on the graph.

In addition to a Hilbert space we need a unitary operator that advances the walk one time step. We shall give a detailed description of this operator shortly, however when a particle scatters at an external vertex the effect of the operator is to reflect the particle with a phase, and the phases can be different at different vertices. It is these phases that determine the behavior of the walk. One case is where all of the phases are 0 except for one which is  $\pi$ . This simply reproduces the standard Grover search, and the particle becomes localized on the edge with the  $\pi$  phase shift after  $O(\sqrt{N})$  steps. As we mentioned, we can do a search on a more varied background by allowing the phases to have more values. If roughly half have phase shifts of  $2\pi/3$ , roughly another half have phase shifts of  $-2\pi/3$ , and a few have phase shifts of zero, the particle will become localized on the edges with phase shift 0 after  $O(\sqrt{N})$  steps, but the maximum probability of finding it there is not close to one, as in the case of the Grover search, but  $3/4$ . Adding further variety to the phases causes this probability to decrease.

We will begin by examining the eigenvalues of the unitary operator that advances the walk one step. One result is that if most of the phase shifts are the same and only a few are different, the eigenstates group themselves into two subspaces, a big, relatively boring, subspace in which not much happens, and a small space in which interesting things can happen. This allows one to focus on what

is happening in the small space and to ignore the large one, which has the effect of simplifying the analysis of the walk.

Next, in order to connect the quantum walk description of a search to the usual oracle-based one, we analyze what happens when the function the oracle evaluates is not a binary-valued one. We examine two different uses of the oracle. The first way of using the oracle reduces the problem to a standard Grover search, and the probability of finding the marked vertex is close to one. This use of the oracle, however, is not equivalent to what happens in our quantum walk. We, therefore, explore a second way of using the oracle, which is equivalent to the quantum walk on the star graph. Using the oracle in this way, for a function that can take more than two values, corresponds to allowing more than two phases in the quantum walk. For this use of the oracle, we further study how increasing the number of possible function values affects the number of steps in the search and the maximum probability of finding the marked element. These results carry over immediately to the quantum walk search on the star graph with multiple phases and give us the number of steps required to localize the particle and the extent to which it can become localized.

## II. EIGENSTATES AND EIGENVALUES OF STAR GRAPHS

Let us now specify the details of the walk on a star graph. As noted in the introduction, if the star graph has  $N$  edges, we denote the central vertex by 0 and the other vertices by 1 through  $N$ . When the particle making the walk is reflected from the vertex at  $j$ , it picks up a phase factor  $e^{i\phi_j}$ . The action of the unitary operator that advances the walk one step is given by

$$\begin{aligned} U|0, j\rangle &= e^{i\phi_j}|j, 0\rangle, \\ U|j, 0\rangle &= -r|0, j\rangle + t \sum_{k=1, k \neq j}^N |0, k\rangle \\ &= -|0, j\rangle + t \sum_{k=1}^N |0, k\rangle, \end{aligned} \quad (1)$$

where  $r = (N-2)/N$  and  $t = 2/N$ . Let

$$|\psi\rangle = \sum_{j=1}^N (\alpha_j |j, 0\rangle + \beta_j |0, j\rangle). \quad (2)$$

If  $|\psi\rangle$  is an eigenvector of  $U$  with eigenvalue  $e^{i\theta}$ , i.e.  $U|\psi\rangle = e^{i\theta}|\psi\rangle$ , then we have that

$$\begin{aligned} -\alpha_j + t \sum_{k=1}^N \alpha_k &= e^{i\theta} \beta_j \\ e^{i\phi_j} \beta_j &= e^{i\theta} \alpha_j. \end{aligned} \quad (3)$$

Solving the second equation for  $\beta_j$  and inserting it into the first gives

$$-\alpha_j + tS = e^{i(2\theta - \phi_j)} \alpha_j, \quad (4)$$

where  $S = \sum_{j=1}^N \alpha_j$ . There are now two cases. First, suppose  $S \neq 0$ . Then solving the above equation for  $\alpha_j$  and summing it over  $j$  gives

$$S = tS \sum_{j=1}^N \frac{1}{(e^{i(2\theta - \phi_j)} + 1)}, \quad (5)$$

so that the eigenvalues are determined by the equation

$$1 = t \sum_{j=1}^N \frac{1}{(e^{i(2\theta - \phi_j)} + 1)}. \quad (6)$$

If  $S = 0$  we have that

$$-\alpha_j = e^{i(2\theta - \phi_j)} \alpha_j, \quad (7)$$

which implies, if  $\alpha_j \neq 0$ , that

$$e^{i(2\theta - \phi_j)} = -1. \quad (8)$$

The only way this can be satisfied is if two or more of the  $\phi_j$ 's have the same value.

Let's look at an example. Suppose  $\phi_1 = \pi$  and  $\phi_j = 0$  for  $j \geq 2$ . This is, of course, the standard Grover case. Setting  $z = e^{2i\theta}$  Eq. (6) becomes

$$z^2 - 2rz + 1 = 0, \quad (9)$$

giving  $z = r \pm i(1-r^2)^{1/2}$ . Setting  $e^{2i\theta_0} = r + i(1-r^2)^{1/2}$ , we find that this gives a value for  $\theta_0$  that is of order  $1/\sqrt{N}$ . We thus have four eigenvalues of this type,  $\pm e^{i\theta_0}$  and  $\pm e^{-i\theta_0}$ . Eq. (8) yields the eigenvalues  $\pm i$ . Each of these eigenvalues is  $(N-2)$ -fold degenerate. It is also quite straightforward to find the eigenvectors. For  $e^{i\theta_0}$  we find

$$\begin{aligned} |\xi_1\rangle &= \eta \left[ \frac{1}{1 - e^{2i\theta_0}} (|1, 0\rangle - e^{i\theta_0}|0, 1\rangle) \right. \\ &\quad \left. + \sum_{j=2}^N \frac{1}{1 + e^{2i\theta_0}} (|j, 0\rangle + e^{i\theta_0}|0, j\rangle) \right], \end{aligned} \quad (10)$$

and for  $-e^{i\theta_0}$

$$\begin{aligned} |\xi_2\rangle &= \eta \left[ \frac{1}{1 - e^{2i\theta_0}} (|1, 0\rangle + e^{i\theta_0}|0, 1\rangle) \right. \\ &\quad \left. + \sum_{j=2}^N \frac{1}{1 + e^{2i\theta_0}} (|j, 0\rangle - e^{i\theta_0}|0, j\rangle) \right]. \end{aligned} \quad (11)$$

The expression for the eigenvector  $|\xi_3\rangle$ , which corresponds to eigenvalue  $e^{-i\theta_0}$ , is the same as that for  $|\xi_1\rangle$  except with  $\theta_0$  replaced by  $-\theta_0$ , and the expression for  $|\xi_4\rangle$ , which corresponds to eigenvalue  $-e^{-i\theta_0}$ , is the same

as that for  $|\xi_2\rangle$  except with  $\theta_0$  replaced by  $-\theta_0$ . The normalization constant  $\eta$  is given by

$$\eta = \frac{\sin(2\theta_0)}{[N - (N-2)\cos(2\theta_0)]^{1/2}} \cong \frac{1}{\sqrt{N}}. \quad (12)$$

The eigenvectors corresponding to the eigenvalues  $\pm i$  are just determined by the condition  $\sum_{j=2}^N \alpha_j = 0$ , so there is a great deal of freedom in choosing them. One possible choice for the eigenvectors corresponding to  $+i$  is

$$|\zeta_m^{(+)}\rangle = \frac{1}{\sqrt{2(N-1)}} \sum_{j=2}^N e^{2\pi i(j-1)m/(N-1)} (|j, 0\rangle + i|0, j\rangle), \quad (13)$$

for  $m = 1, 2, \dots, N-2$ . The eigenvectors corresponding to  $-i$  can be taken to be of similar form. If we start the walk in the state

$$|\psi_0\rangle = \frac{1}{\sqrt{2N}} \sum_{j=1}^N (|j, 0\rangle + |0, j\rangle), \quad (14)$$

we find that  $|\psi_0\rangle \cong (|\xi_1\rangle + |\xi_3\rangle)/\sqrt{2}$ , assuming  $N \gg 1$ , so that the whole walk takes place in the space spanned

by  $|\xi_1\rangle$  and  $|\xi_3\rangle$ . After  $m$  steps the state of the particle is

$$|\psi_m\rangle = \frac{1}{\sqrt{2}} (e^{im\theta_0} |\xi_1\rangle + e^{-im\theta_0} |\xi_3\rangle) + O(N^{-1/2}), \quad (15)$$

and when  $m\theta_0 = \pi/2$ , we have

$$|\psi_m\rangle = -\frac{1}{\sqrt{2}} (|1, 0\rangle - |0, 1\rangle) + O(N^{-1/2}). \quad (16)$$

Therefore, to good approximation, the particle is located on the edge connected to the vertex with the  $\pi$  phase shift, just as in the Grover algorithm.

Now let us move on to a more complicated situation. Suppose that each of the  $\phi_j$ 's takes on one of three values  $\phi_1, \phi_2$  or  $\phi_3$ , and, in particular,  $n_1$  are  $\phi_1$ ,  $n_2$  are  $\phi_2$ , and  $n_3$  are  $\phi_3$ , where  $n_1 + n_2 + n_3 = N$ . Then Eq. (6) gives us

$$\frac{N}{2} = \sum_{j=1}^3 \frac{n_j}{ze^{-i\phi_j} + 1}, \quad (17)$$

which becomes

$$z^3 - 2 \left[ \sum_{j=1}^3 x_j e^{i\phi_j} \right] z^2 + 2 \left[ x_1 e^{i(\phi_2+\phi_3)} + x_2 e^{i(\phi_1+\phi_3)} + x_3 e^{i(\phi_1+\phi_2)} \right] z - e^{i(\phi_1+\phi_2+\phi_3)} = 0, \quad (18)$$

where  $x_j = n_j/N$ .

One case in which the equation can be solved is if  $\phi_1 = 2\pi/3$ ,  $\phi_2 = -2\pi/3$  and  $\phi_3 = 0$ . In addition, we shall assume that  $N$  is odd, and  $n_1 = n_2 = (N-1)/2$ , and  $n_3 = 1$ . In that case, the above equation becomes

$$z^3 + \left( \frac{N-3}{N} \right) z^2 - \left( \frac{N-3}{N} \right) z - 1 = 0. \quad (19)$$

The solutions to this equation are  $z = 1$  and

$$z = -1 + \frac{3}{2N} \pm i \left( \frac{3}{N} - \frac{9}{4N^2} \right)^{1/2}. \quad (20)$$

Note that the last pair of solutions for  $z$  yields eigenvalues of the form  $\pm[i \pm O(N^{-1/2})]$ . Comparing to the previous case we can surmise that this situation will again lead to a search in which we can find the vertex corresponding to the phase of zero in  $O(\sqrt{N})$  steps. This is, in fact the case, but rather than use the exact eigenvalues to show this, we shall make use of an approximation method which allows us to find a solution in a more general case.

We want to look at the situation in which one of the  $x_j$ 's is small, typically of order  $1/N$ , and  $N \gg 1$ . In that case the above equation can be solved perturbatively. If

the full third degree polynomial above is  $f(z)$ , then let  $f_0(z)$  be the same polynomial but with the terms of order  $1/N$  set equal to zero, that is,  $f_0(z)$  is the  $N \rightarrow \infty$  limit of  $f(z)$ . We then get a zeroth order solution to  $f(z) = 0$ , which we shall denote by  $z_0$ , by solving  $f_0(z) = 0$ .

We can then get corrections to the zeroth order solution, by setting  $z = z_0 + \delta z$  substituting back into the cubic equation, keeping the lowest order surviving terms, and then solving for  $\delta z$ . This will give us

$$\begin{aligned} 0 &= f(z_0 + \delta z) \\ &\cong f(z_0) + [f'_0(z_0) + O(1/N)]\delta z \\ &\quad + [f''_0(z_0) + O(1/N)](\delta z)^2 + \dots \end{aligned} \quad (21)$$

Note that we have substituted  $f_0 + O(1/N)$  for  $f$  in the derivative terms, e.g.  $f'_0(z_0) + O(1/N)$  for  $f'(z_0)$ , because the  $f_0$  terms are typically of order one. Because  $f(z_0)$  is usually of order  $1/N$  and the  $(\delta z)^2$  term can be neglected in comparison to the  $\delta z$  term, we will usually have  $\delta z \sim 1/N$ . This situation is not interesting, and let us see why. If  $z_0 = e^{i\theta_0}$ , and  $z_0 + \delta z = e^{i(\theta_0 + \delta\theta)}$ , then  $\delta z \sim 1/N$  will give us  $\delta\theta \sim 1/N$ . After  $m$  steps we will have a factor of  $e^{im(\theta_0 + \delta\theta)}$  multiplying the eigenvector. The part of this phase factor that contains the information about the

unusual vertices (the ones with the rare phase shift) is  $e^{im\delta\theta}$  and this phase factor will differ substantially from one when  $m$  is of order  $N$ . That means that in this case, the unusual vertices would have a significant effect on the dynamics after  $N$  steps. However, we could find the unusual vertices classically just by checking each vertex to see if its phase shift was one of the rare ones, and this procedure would take  $N$  steps too, so we have gained nothing.

What we really want, to get some kind of a quadratic speedup, is for  $\delta\theta \sim 1/\sqrt{N}$ . Examining Eq. (21), we see that this can happen if the order one part of the coefficient of  $\delta z$ , which is  $f'_0(z_0)$ , vanishes. We then get a quadratic equation for  $\delta z$

$$0 = f(z_0) + [O(1/N)]\delta z + f''_0(z_0)(\delta z)^2, \quad (22)$$

which typically results in  $\delta z$  being of order  $N^{-1/2}$ . The condition  $f'_0(z_0) = 0$  implies that  $z_0$  is a double root of the equation  $f_0(z_0) = 0$ . Summarizing, in order to find a situation in which we can have a quadratic speed up, we want the solution of the zeroth order problem to be a double root.

Let's give an example where that happens. We are going to assume that almost half of the phases are  $\phi_1 = 2\pi/3$ , almost half are  $\phi_2 = -2\pi/3$  and the remainder are  $\phi_3 = 0$ . We will take most of the phases from  $j = 1$  to  $j = N/2$  to be  $2\pi/3$ , and most of the phases between  $j = (N/2) + 1$  to  $j = N$  to be  $-2\pi/3$ , and we will split the phases that are zero putting some between  $j = 1$  to  $j = N/2$  and some between  $j = (N/2) + 1$  to  $j = N$ . In particular, let us assume that for  $j = 1, \dots, n_{31}$  and for  $j = (N/2) + 1, \dots, (N/2) + n_{32}$  the phases are 0, for  $j = n_{31} + 1, \dots, (N/2)$  the phases are  $2\pi/3$  and for  $j = (N/2) + n_{32} + 1, \dots, N$  the phases are  $-2\pi/3$ . Therefore, the total number of edges with phase 0 is  $n_3 = n_{31} + n_{32}$ , the total number with phase  $2\pi/3$  is  $n_1 = (N/2) - n_{31}$  and the total number with phase  $-2\pi/3$  is  $n_2 = (N/2) - n_{32}$ . We will consider the case where  $n_3 \ll N$ .

Let us first get the zeroth order solution to Eq. (18). To do so, we neglect all quantities of order  $n_3/N$ , which means we set  $x_1 = x_2 = 1/2$  and  $x_3 = 0$ . Doing so, we get the equation

$$z^3 + z^2 - z - 1 = (z^2 - 1)(z + 1) = 0, \quad (23)$$

which has a double root at  $z = -1$ . Let us now calculate corrections to this root. First we set  $z = -1 + \delta z$  and substitute it back into the equation. We then keep terms of up to second order in small quantities, with both  $\delta z$  and anything of order  $n_3/N$  being considered small. Toward this end, we set  $x_{31} = n_{31}/N$  and  $x_{32} = n_{32}/N$ , and note that

$$\sum_{j=1}^3 x_j e^{i\phi_j} = -\frac{1}{2} + x_{31}(1 - e^{2\pi i/3}) + x_{32}(1 - e^{-2\pi i/3}), \quad (24)$$

and

$$\begin{aligned} & x_1 e^{i(\phi_2 + \phi_3)} + x_2 e^{i(\phi_1 + \phi_3)} + x_3 e^{i(\phi_1 + \phi_2)} \\ &= -\frac{1}{2} + x_{31}(1 - e^{-2\pi i/3}) + x_{32}(1 - e^{2\pi i/3}). \end{aligned} \quad (25)$$

Upon making our substitution and keeping only quantities of up to second order, we first find that, as expected, the term proportional to  $\delta z$  vanishes, and we are left with

$$\begin{aligned} & (\delta z)^2 - \{x_{31}[2(1 - e^{2\pi i/3}) + (1 - e^{-2\pi i/3})] \\ & + x_{32}[2(1 - e^{-2\pi i/3}) + (1 - e^{2\pi i/3})]\}\delta z + 3x_3 = 0. \end{aligned} \quad (26)$$

To solve this one can either use the quadratic formula, or note that if we neglect the  $\delta z$  term this gives  $\delta z \sim 1/\sqrt{N}$ , and then making use of this we see that both the constant term in the equation and  $(\delta z)^2$  are of order  $1/N$ , while the term linear in  $\delta z$  is of order  $1/N^{3/2}$ , i.e. smaller than the other two, so that it represents an even smaller order correction. This implies that to lowest order we have

$$\delta z = \pm i\sqrt{3x_3}. \quad (27)$$

Define  $\theta_0$  so that

$$e^{2i\theta_0} = 1 + i\sqrt{3x_3}, \quad (28)$$

which implies that  $z = -e^{\pm 2i\theta_0}$  and that  $2\theta_0 = \sqrt{3x_3}$ . The eigenvalues of  $U$  corresponding to these values of  $z$  are  $\pm i e^{\pm i\theta_0}$ .

We can now find the eigenstates corresponding to these four eigenvalues by making use of the relations

$$\alpha_j = \frac{\eta}{\lambda^2 e^{-i\phi_j} + 1} \quad \beta_j = \lambda e^{-i\phi_j} \alpha_j, \quad (29)$$

where  $\lambda$  is one of the eigenvalues, and  $\eta$  is a normalization constant. This gives us the four eigenvectors: for  $\lambda = i e^{i\theta_0}$

$$\begin{aligned} |\xi_1\rangle = \eta & \left[ \sum_{j=1}^{n_{31}} \frac{|j, 0\rangle + i e^{i\theta_0} |0, j\rangle}{1 - e^{2i\theta_0}} \right. \\ & + \sum_{j=n_{31}+1}^{N/2} \frac{|j, 0\rangle + i e^{i\theta_0} e^{-2\pi i/3} |0, j\rangle}{1 - e^{2i\theta_0} e^{-2\pi i/3}} \\ & + \sum_{j=(N/2)+1}^{(N/2)+n_{32}} \frac{|j, 0\rangle + i e^{i\theta_0} |0, j\rangle}{1 - e^{2i\theta_0}} \\ & \left. + \sum_{j=(N/2)+n_{32}+1}^N \frac{|j, 0\rangle + i e^{i\theta_0} e^{2\pi i/3} |0, j\rangle}{1 - e^{2i\theta_0} e^{2\pi i/3}} \right], \end{aligned} \quad (30)$$

for  $\lambda = -ie^{i\theta_0}$

$$\begin{aligned}
|\xi_2\rangle = \eta & \left[ \sum_{j=1}^{n_{31}} \frac{|j, 0\rangle - ie^{i\theta_0} |0, j\rangle}{1 - e^{2i\theta_0}} \right. \\
& + \sum_{j=n_{31}+1}^{N/2} \frac{|j, 0\rangle - ie^{i\theta_0} e^{-2\pi i/3} |0, j\rangle}{1 - e^{2i\theta_0} e^{-2\pi i/3}} \\
& + \sum_{j=(N/2)+1}^{(N/2)+n_{32}} \frac{|j, 0\rangle - ie^{i\theta_0} |0, j\rangle}{1 - e^{2i\theta_0}} \\
& \left. + \sum_{j=(N/2)+n_{32}+1}^N \frac{|j, 0\rangle - ie^{i\theta_0} e^{2\pi i/3} |0, j\rangle}{1 - e^{2i\theta_0} e^{2\pi i/3}} \right], \quad (31)
\end{aligned}$$

for  $\lambda = ie^{-i\theta_0}$

$$\begin{aligned}
|\xi_3\rangle = \eta & \left[ \sum_{j=1}^{n_{31}} \frac{|j, 0\rangle + ie^{-i\theta_0} |0, j\rangle}{1 - e^{-2i\theta_0}} \right. \\
& + \sum_{j=n_{31}+1}^{N/2} \frac{|j, 0\rangle + ie^{-i\theta_0} e^{-2\pi i/3} |0, j\rangle}{1 - e^{-2i\theta_0} e^{-2\pi i/3}} \\
& + \sum_{j=(N/2)+1}^{(N/2)+n_{32}} \frac{|j, 0\rangle + ie^{-i\theta_0} |0, j\rangle}{1 - e^{-2i\theta_0}} \\
& \left. + \sum_{j=(N/2)+n_{32}+1}^N \frac{|j, 0\rangle + ie^{-i\theta_0} e^{2\pi i/3} |0, j\rangle}{1 - e^{-2i\theta_0} e^{2\pi i/3}} \right], \quad (32)
\end{aligned}$$

and, finally, for  $\lambda = -ie^{-i\theta_0}$

$$\begin{aligned}
|\xi_4\rangle = \eta & \left[ \sum_{j=1}^{n_{31}} \frac{|j, 0\rangle - ie^{-i\theta_0} |0, j\rangle}{1 - e^{-2i\theta_0}} \right. \\
& + \sum_{j=n_{31}+1}^{N/2} \frac{|j, 0\rangle - ie^{-i\theta_0} e^{-2\pi i/3} |0, j\rangle}{1 - e^{-2i\theta_0} e^{-2\pi i/3}} \\
& + \sum_{j=(N/2)+1}^{(N/2)+n_{32}} \frac{|j, 0\rangle - ie^{-i\theta_0} |0, j\rangle}{1 - e^{-2i\theta_0}} \\
& \left. + \sum_{j=(N/2)+n_{32}+1}^N \frac{|j, 0\rangle - ie^{-i\theta_0} e^{2\pi i/3} |0, j\rangle}{1 - e^{-2i\theta_0} e^{2\pi i/3}} \right]. \quad (33)
\end{aligned}$$

We find that for these vectors to be normalized, we must have  $\eta = (1/2)\sqrt{3/N}$ .

Now in order to help make sense of this, define the

following four vectors

$$\begin{aligned}
|\zeta_1^{(\pm)}\rangle = \frac{i}{\sqrt{2n_3}} & \left[ \sum_{j=1}^{n_{31}} (|j, 0\rangle \pm i|0, j\rangle) \right. \\
& \left. + \sum_{j=(N/2)+1}^{(N/2)+n_{32}} (|j, 0\rangle \pm i|0, j\rangle) \right], \quad (34)
\end{aligned}$$

and

$$\begin{aligned}
|\zeta_2^{(\pm)}\rangle = \eta\sqrt{2} & \left[ \sum_{j=n_{31}+1}^{N/2} \frac{|j, 0\rangle \pm ie^{-2\pi i/3} |0, j\rangle}{1 - e^{-2\pi i/3}} \right. \\
& \left. + \sum_{j=(N/2)+n_{32}+1}^N \frac{|j, 0\rangle \pm ie^{2\pi i/3} |0, j\rangle}{1 - e^{2\pi i/3}} \right]. \quad (35)
\end{aligned}$$

Now we have, approximately, that

$$\begin{aligned}
|\xi_1\rangle &= \frac{1}{\sqrt{2}} (|\zeta_1^{(+)}\rangle + |\zeta_2^{(+)}\rangle) \\
|\xi_3\rangle &= \frac{1}{\sqrt{2}} (-|\zeta_1^{(+)}\rangle + |\zeta_2^{(+)}\rangle) \\
|\xi_2\rangle &= \frac{1}{\sqrt{2}} (|\zeta_1^{(-)}\rangle + |\zeta_2^{(-)}\rangle) \\
|\xi_4\rangle &= \frac{1}{\sqrt{2}} (-|\zeta_1^{(-)}\rangle + |\zeta_2^{(-)}\rangle). \quad (36)
\end{aligned}$$

Now let us consider for the initial state of the walk the state

$$|\psi_{init}\rangle = \frac{1}{\sqrt{N}} \sum_{j=1}^N |j, 0\rangle. \quad (37)$$

Defining the state

$$|\zeta_0\rangle = \frac{1}{\sqrt{2}} (|\zeta_2^{(+)}\rangle + |\zeta_2^{(-)}\rangle) \quad (38)$$

we find that  $|\langle\zeta_0|\psi_{init}\rangle| = \sqrt{3}/2$ , i.e. these states have a substantial overlap. Noting that

$$|\zeta_0\rangle = \frac{1}{2} \sum_{k=1}^4 |\xi_k\rangle \quad (39)$$

we see that after  $m$  steps

$$\begin{aligned}
|\zeta_0\rangle &\rightarrow \frac{1}{2} [(ie^{i\theta_0})^m |\xi_1\rangle + (ie^{-i\theta_0})^m |\xi_3\rangle \\
&\quad + (-ie^{i\theta_0})^m |\xi_2\rangle + (-ie^{-i\theta_0})^m |\xi_4\rangle] \\
&= \frac{1}{\sqrt{2}} \left\{ i^m [i \sin(m\theta_0) |\zeta_1^{(+)}\rangle + \cos(m\theta_0) |\zeta_2^{(+)}\rangle] \right. \\
&\quad \left. + (-i)^m [i \sin(m\theta_0) |\zeta_1^{(-)}\rangle + \cos(m\theta_0) |\zeta_2^{(-)}\rangle] \right\}. \quad (40)
\end{aligned}$$

What this means is that when  $m\theta_0 = \pi/2$ ,  $|\zeta_0\rangle$  becomes a state with the particle located on the edges with  $\phi = 0$ . Because the overlap between  $U^m|\psi_{init}\rangle$  and  $U^m|\zeta_0\rangle$  is the same as the overlap between  $|\psi_{init}\rangle$  and  $|\zeta_0\rangle$ , if we measure the location of the particle in the state  $U^m|\psi_{init}\rangle$ , we will find that it is, with probability  $3/4$  on one of the edges connected to a vertex with  $\phi = 0$ . Note that in order to find edges with  $\phi = 0$  with this procedure, we do not have to know which edges have  $\phi = 2\pi/3$  and which have  $\phi = -2\pi/3$ . Also note that unlike in the usual Grover algorithm, the probability to find a vertex with  $\phi = 0$  does not go to one as  $N \rightarrow \infty$ .

### III. ORACLE VERSION

Now let us show that we can rephrase our walk search problem as an oracle search problem. We consider the situation in which the function that the oracle evaluates is multivalued. The function  $f(j)$  with  $j = 1, 2, \dots, N$  can take on integer values between 0 and  $(d-1)$ , where  $d$  is an integer greater than 2. It is 0 at  $j = j_0$ , while it takes on one of the integer values between 1 and  $(d-1)$  at other values of  $j$ . We are interested in finding  $j_0$ . We will do this in two different ways. The first reduces to a standard Grover search, and the probability of finding  $j_0$  using this method is close to 1. The second results in a search that is analogous to what happens in the quantum walk on the star graph, and in that case the probability of finding  $j_0$  is a decreasing function of  $d$ .

Let us first show how the oracle for a multivalued function can be used to perform a standard Grover search [21, 22]. Our oracle has the action  $O|j\rangle_1|k\rangle_2 = |j\rangle_1|k \oplus f(j)\rangle_2$ , where  $|j\rangle_1$  is the state of the input register that contains the argument of the function,  $|k\rangle_2$  is a qudit state, and  $\oplus$  denotes addition modulo  $d$ . We need to manufacture an operation that flips the sign of  $|j_0\rangle$  but leaves other states  $|j\rangle$ , where  $j \neq j_0$  alone. The operator  $O^{-1}(I_1 \otimes I_2 - 2I_1 \otimes |0\rangle_2\langle 0|)O$ , where  $I_m$  for  $m = 1, 2$  is the identity on space  $m$ , does exactly that. We have

$$\begin{aligned} & O^{-1}(I_1 \otimes I_2 - 2I_1 \otimes |0\rangle_2\langle 0|)O|j\rangle_1|0\rangle_2 \\ &= \begin{cases} |j\rangle_1|0\rangle_2 & j \neq j_0 \\ -|j\rangle_1|0\rangle_2 & j = j_0 \end{cases}, \end{aligned} \quad (41)$$

because  $j_0$  is the only value of  $j$  for which  $f(j)$  is 0. Note that because  $O^d$  is the identity, we have that  $O^{-1} = O^{d-1}$  so that we can realize  $O^{-1}$  simply by applying  $O$   $d-1$  times. If we now apply this operator followed by the usual inversion-about-the-mean operator, we have the standard Grover iteration. This will result in a probability of close to one of finding  $j_0$ .

That is not how our quantum walk worked, however. There, we had  $d$  different phase shifts and different values of  $f(j)$  led to different phase shifts. Let us now formulate an oracle problem in which that happens.

We assume for simplicity that  $f(j) = 0$  at only one particular value of  $j$ , say  $j_0$ . To search for  $j_0$ , we

adopt a straightforward extension of the standard Grover algorithm [21, 22]. The starting point of our search is to take a  $d \times d$  matrix

$$H_d = \begin{pmatrix} 1 & 1 & 1 & \cdots & 1 \\ 1 & \beta & \beta^2 & \cdots & \beta^{d-1} \\ 1 & \beta^2 & \beta^4 & \cdots & \beta^{2(d-1)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \beta^{d-1} & \beta^{2(d-1)} & \cdots & \beta^{(d-1)^2} \end{pmatrix} \quad (42)$$

as a  $d$ -dimensional equivalent of the Hadamard operator  $H \equiv H_2 = \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$ , where

$$\beta = e^{i2\pi/d} \quad (\beta^d = 1), \quad (43)$$

and

$$1 + \beta + \beta^2 + \cdots + \beta^{d-1} = 0. \quad (44)$$

With  $n$  qudits in the main register ( $d^n = N$ ) and one ancilla qudit in the secondary register, the search is then conducted in the following order:

1. Initial preparation. The  $n$  qudits are prepared in state

$$(H_d|0\rangle)^{\otimes n} = \left( \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} |k\rangle \right)^{\otimes n} = \frac{1}{\sqrt{N}} \sum_{j=1}^N |j\rangle, \quad (45)$$

and the ancilla qudit in state  $H_d|1\rangle = \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \beta^k |k\rangle$ .

2. Oracle transformation. The  $n$  qudits along with the ancilla qudit are directed to an oracle in which the transformation  $|j\rangle|k\rangle \rightarrow |j\rangle|f(j) \oplus k\rangle$  is performed, where  $|j\rangle$  and  $|k\rangle$  represent the state of the  $n$  qudits in the main register and of the ancilla qudit in the secondary register, respectively. When the ancilla qudit is in state  $\frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \beta^k |k\rangle$ , the oracle transformation is represented by

$$|j\rangle \left( \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \beta^k |k\rangle \right) \rightarrow \beta^{-f(j)} |j\rangle \left( \frac{1}{\sqrt{d}} \sum_{k=0}^{d-1} \beta^k |k\rangle \right). \quad (46)$$

The oracle transformation can thus be written, in the  $N$ -dimensional Hilbert space of the  $n$ -qudit state, as  $O = \sum_{j=1}^N \beta^{-f(j)} |j\rangle\langle j|$ .

3. Inversion about average. This is achieved by

$$D = \frac{2}{N} \sum_{i=1}^N \sum_{j=1}^N |i\rangle\langle j| - \sum_{j=1}^N |j\rangle\langle j|. \quad (47)$$

4. Repeat the operation  $G \equiv DO$  until the state of the  $n$  qudits in the main register comes sufficiently close to the desired state  $|j_0\rangle$ .

5. Measure the state of the  $n$  qudits in the main register to get the solution. Evaluate  $f(j)$  to check that the solution is correct. If the solution is correct, then the search succeeds. If not, go back to (1) and start the search over again. This last step (5) is necessary because, as will be shown later, the probability for the search to produce the correct solution is less than 1 for  $d \geq 3$ .

In the  $N$ -dimensional Hilbert space, the  $N \times N$  matrix  $G = DO$  is given by

$$G = \frac{2}{N} \sum_{i=1}^N \sum_{j=1}^N \beta^{-f(j)} |i\rangle\langle j| - \sum_{j=1}^N \beta^{-f(j)} |j\rangle\langle j|. \quad (48)$$

After applications of the operator  $G$   $k$  times, the state of the  $n$  qudits in the main register becomes  $|\psi_k\rangle = G^k |\psi_0\rangle$ , where  $|\psi_0\rangle = \frac{1}{\sqrt{N}} \sum_{j=1}^N |j\rangle$ . The probability to find the desired state  $|j_0\rangle$  after  $k$  iterations is then given by

$$P_k = |\langle j_0 | G^k | \psi_0 \rangle|^2 = \left| \sum_{j=1}^N z_j^k \langle j_0 | \chi_j \rangle \langle \chi_j | \psi_0 \rangle \right|^2, \quad (49)$$

where  $z_j$ 's and  $|\chi_j\rangle$ 's are eigenvalues and eigenvectors of the matrix  $G$ . The probability  $P_k$  can thus be determined by calculating  $z_j$ 's and  $|\chi_j\rangle$ 's. When  $d = 2$ , the method described here reproduces the results of the standard Grover search with a two-valued function.

Let us now find the eigenvalues  $z_k$  ( $k = 1, 2, \dots, N$ ) and corresponding eigenvectors  $|\chi_k\rangle$  of the  $N \times N$  matrix  $G$  given by Eq. (48), where

$$G|\chi_k\rangle = z_k |\chi_k\rangle, \quad k = 1, 2, \dots, N \quad (50)$$

The eigenvectors  $|\chi_k\rangle$  can be expressed in a linear superposition of the basis vectors  $|j\rangle$  as

$$|\chi_k\rangle = \sum_{j=1}^N c_{kj} |j\rangle \quad (51)$$

Substituting Eqs. (48) and (51) into Eq. (50) and solving the resulting equation for  $c_{kj}$ , we obtain

$$c_{kj} = \frac{\frac{2}{N} S_k}{\beta^{-f(j)} + z_k}, \quad j = 1, 2, \dots, N \quad (52)$$

where

$$S_k = \sum_{j=1}^N \beta^{-f(j)} c_{kj} \quad (53)$$

Substituting our expression for  $c_{kj}$  back into the above equation, we obtain

$$\sum_{j=1}^N \frac{1}{1 + \beta^{f(j)} z_k} = \frac{N}{2} \quad (54)$$

This is identical to Eq. (6). We note that one application of the Grover operator  $G$  is equivalent to two steps of the walk, and therefore the eigenvalue  $z$  of  $G$  is equal to the square of the eigenvalue,  $\lambda^2$ , of the operator  $U$ . We proceed to solve Eq. (54) for the case where there is one  $j$ , i.e.,  $j_0$ , at which  $f(j) = 0$ . We further assume for simplicity that the integer values between 1 and  $(d-1)$  of the function  $f(j)$  are distributed evenly among other  $(N-1)$   $j$ 's. There is thus one  $j$ , i.e.,  $j_0$ , for which  $f(j) = 0$  and  $\frac{N-1}{d-1}$   $j$ 's each at which  $f(j) = 1, 2, \dots, d-1$ . We assume that  $\frac{N-1}{d-1}$  is an integer. Note that if we choose  $N = d^n$ , where  $n$  is a positive integer, then  $\frac{N-1}{d-1}$  is an integer. Using the relation

$$1 + x + x^2 + \dots + x^{d-1} = \frac{1 - x^d}{1 - x} \quad (55)$$

with  $x = -\beta^j z_k$  and noting that  $\beta^{jd} = 1$ , Eq. (54) can be rewritten as

$$\frac{1}{1 + z_k} + \left( \frac{N-1}{d-1} \right) \frac{1}{1 - (-z_k)^d} \left[ d-1 - \sum_{j=1}^{d-1} (-z_k)^j \right] = \frac{N}{2}, \quad (56)$$

or

$$(-z_k)^d - \left( \frac{2}{d-1} \right) \left( \frac{N-d}{N} \right) \sum_{j=1}^{d-1} (-z_k)^j + 1 = 0. \quad (57)$$

In the limit  $N \rightarrow \infty$ , the above equation becomes

$$(-z_k)^d - \left( \frac{2}{d-1} \right) \sum_{j=1}^{d-1} (-z_k)^j + 1 = 0. \quad (58)$$

This equation, because it has real coefficients, will have as its roots either real numbers or complex conjugate pairs. By examining the above expression and its derivative, both of which must be zero at a double root, one finds that the only double root of this equation is  $z_k = -1$ . This can be seen as follows. Setting the derivative of the above equation equal to zero and then multiplying by  $-z_k$  gives

$$(-z_k)^d - \frac{2}{d(d-1)} \sum_{j=1}^{d-1} j (-z_k)^j = 0. \quad (59)$$

Substituting into this equation for  $(-z_k)^d$  from Eq. (58), we find

$$\frac{2}{d-1} \sum_{j=1}^{d-1} \left( 1 - \frac{j}{d} \right) (-z_k)^j = 1. \quad (60)$$

Noting that

$$\frac{2}{d-1} \sum_{j=1}^{d-1} \left( 1 - \frac{j}{d} \right) = 1, \quad (61)$$



we see that the only way Eq. (60) can be satisfied is if all of the terms add in phase. The only way that can happen is if  $z_k = -1$ , and so, as stated above, that is the only double root.

Going back to the equation for finite  $N$ , we find that this root splits into two roots. In order to obtain an expression for these two roots, we substitute  $z_k = -1 + \delta z$  into Eq. (57) and find  $\delta z \approx \pm i \sqrt{\frac{12}{N(d+1)}}$ . We designate these two roots as  $z_1$  and  $z_2$ . We can now write

$$z_1 \approx -1 + i \sqrt{\frac{12}{N(d+1)}} \equiv -e^{-i2\theta} \quad (62)$$

$$z_2 \approx -1 - i \sqrt{\frac{12}{N(d+1)}} \equiv -e^{i2\theta} \quad (63)$$

where

$$\sin(2\theta) \approx 2\theta \approx \sqrt{\frac{12}{N(d+1)}}. \quad (64)$$

The finite  $N$  corrections of these roots are of order  $N^{-1/2}$ , whereas the finite  $N$  corrections to the other roots of Eq. (58) are of order  $N^{-1}$ .

Let us now find the eigenvectors  $|\chi_1\rangle$  and  $|\chi_2\rangle$  corresponding to  $z_1$  and  $z_2$ . Using Eqs. (51) and (52), we write

$$|\chi_1\rangle = \eta_1 \sum_{j=1}^N \frac{1}{\beta^{-f(j)} - e^{-i2\theta}} |j\rangle \quad (65)$$

and similarly for  $|\chi_2\rangle$ . The normalization condition  $1 = \langle \chi_1 | \chi_1 \rangle$  determines the constant  $\eta_1$ . Using again Eq. (55) with  $x = \beta^{-j} e^{i2\theta}$ , the normalization condition is reduced to

$$1 = |\eta_1|^2 \left\{ \frac{1}{2(1 - \cos 2\theta)} + \left( \frac{N-1}{d-1} \right) \frac{1}{2[1 - \cos(2d\theta)]} [d(d-1) - 2 \sum_{j=1}^{d-1} (d-j) \cos(2j\theta)] \right\} \quad (66)$$

For  $(n/N) \ll 1$ , we have that  $\cos 2n\theta \approx 1 - n^2 \frac{6}{N(d+1)}$ , we obtain  $\eta_1 \approx \sqrt{\frac{6}{N(d+1)}}$ , and thus

$$|\chi_1\rangle \approx \sqrt{\frac{6}{N(d+1)}} \sum_{j=1}^N \frac{1}{\beta^{-f(j)} - e^{-i2\theta}} |j\rangle \quad (67)$$

and similarly

$$|\chi_2\rangle \approx \sqrt{\frac{6}{N(d+1)}} \sum_{j=1}^N \frac{1}{\beta^{-f(j)} - e^{i2\theta}} |j\rangle \quad (68)$$

Other eigenstates are given by

$$|\chi_k\rangle = \eta_k \sum_{j=1}^N \frac{1}{\beta^{-f(j)} + z_k} |j\rangle, \quad k = 3, 4, \dots, N \quad (69)$$

and the normalizations for these states are

$$|\eta_k|^2 \sum_{j=1}^N \left| \frac{1}{\beta^{-f(j)} + z_k} \right|^2 = 1. \quad (70)$$

For any  $j$  and  $k$ , we have  $|\beta^{-f(j)} + z_k| < 2$ . Therefore,  $\eta_k$ 's should satisfy

$$|\eta_k|^2 \frac{N}{4} < 1 \quad (71)$$

or

$$\eta_k < \frac{2}{\sqrt{N}}. \quad (72)$$

The probability  $P_k$  to find the desired state  $|j_0\rangle$  after  $k$  iterations is given by Eq. (49). In order to evaluate this, we begin by noting that

$$\langle j_0 | \chi_1 \rangle \approx -i \frac{1}{\sqrt{2}}, \quad \langle j_0 | \chi_2 \rangle \approx i \frac{1}{\sqrt{2}}, \quad (73)$$

which implies that  $\langle j_0 | \chi_k \rangle \approx 0$ . Therefore, we have that

$$P_k \approx |z_1^k \langle j_0 | \chi_1 \rangle \langle \chi_1 | \psi_0 \rangle + z_2^k \langle j_0 | \chi_2 \rangle \langle \chi_2 | \psi_0 \rangle|^2 \quad (74)$$

From Eq. (67), we have

$$\langle \chi_1 | \psi_0 \rangle \approx \sqrt{\frac{6}{N^2(d+1)}} \sum_{j=1}^N \frac{1}{\beta^{f(j)} - e^{i2\theta}} \quad (75)$$

and similarly for  $\langle \chi_2 | \psi_0 \rangle$ . Evaluating the summation in this equation using Eq. (54), we obtain

$$\langle \chi_1 | \psi_0 \rangle = \langle \chi_2 | \psi_0 \rangle \approx -\sqrt{\frac{3}{2(d+1)}} \quad (76)$$

Making use of these results, we finally obtain for  $P_k$

$$P_k \approx \frac{3}{d+1} \sin^2(2k\theta). \quad (77)$$

The maximum probability is given by  $P_{max} \approx \frac{3}{d+1}$  and the number of iterations to reach the maximum probability is  $k_{max} \approx \frac{\pi}{4\theta} \approx \frac{\pi}{4} \sqrt{\frac{N(d+1)}{3}}$ . Since the maximum probability decreases as  $d$  increases, it gets increasingly important to check in the end that the solution produced by the search is indeed correct. This can be done either classically or quantum mechanically. When  $d = 3$ , we have  $P_{max} \approx \frac{3}{4}$ , in agreement with the result presented in the previous section. Although the maximum probability decreases with  $d$ , the number  $k_{max}$  of iterations required to reach the maximum probability still scales as  $\sqrt{N}$  for  $d \geq 3$ . This indicates that the quantum walk search still provides quadratic speedup, the best one can gain from a quantum search algorithm [23–25], for  $d \geq 3$ .

So far, we have assumed that there is a single match. We now briefly comment on the situation where there are multiple matches [24]. If the number  $M$  of matches is small compared with the total number  $N$ , only a trivial modification is needed. Since the algebra involved is straightforward, we here state the result only. The maximum probability  $P_{max}$  remains the same as long as  $M \ll N$ , but the number  $k_{max}$  of steps to reach the maximum probability decreases with  $M$ . In fact, we obtain  $P_k \approx \frac{3}{d+1} \sin^2(2k\theta)$ , where  $\sin \theta = \sqrt{\frac{3M}{N(d+1)}}$ , and therefore  $k_{max} \approx \frac{\pi}{4} \sqrt{\frac{N(d+1)}{3M}}$ . The number  $k_{max}$  scales as  $1/\sqrt{M}$ . Classically, when there are  $M$  matches, one would expect to find a match  $M$  times faster than when there is a single match. One thus expects that classically the number  $k_{max}$  to scale as  $1/M$ . It is then natural to expect that this number would scale as  $1/\sqrt{M}$  for the quantum search which is quadratically faster than any classical search. If the number  $M$  of matches is not small compared with  $N$ , a different search strategy from that described so far should be employed because a match can be found with a reasonably high probability after only a small number of iterations. In general, the state of the qudits in the main register changes rapidly, and it seems difficult to find a general rule here. One simply needs to find the best strategy for a given situation.

#### IV. CONCLUSION

We have studied a quantum walk search on a star graph and an analogous oracle search problem. The quantum

walk is trying to find the set of vertices that reflect the particle with a particular phase, and the oracle search is trying to find the input values for a multivalued function that produce a particular output value. We have shown that the maximum probability  $P_{max}$  for the search on a star graph to find the correct vertex decreases with  $d$ , the number of different phases with which the particle is reflected from outer vertices. We have also shown that the number  $k_{max}$  of steps needed to reach the maximum probability increases with  $d$ . These two effects combine to increase the average number of steps required to be taken in order to find the correct vertex. Despite that, however, this number still scales as  $\sqrt{N}$  and the quantum walk search still gives the quadratic speedup over a classical search for any value of  $d$  as long as  $N \gg d$ .

We examined two different oracle searches. One mimicked the quantum walk, and for this search the probability of finding the inputs that produced a particular output decreased with  $d$ , just as for the quantum walk. However, we also showed that the oracle can be used in a different way so that the probability of finding the desired elements is close to one. This indicates that the use of the oracle is more powerful than the quantum walk search we studied.

Finally, we would like to note that it has been possible in some cases to improve the maximum probability of finding a marked element in a quantum walks search. The original quantum-walk search on the hypercube due to Shenvi, Kempe and Whaley found the marked vertex with a probability of less than  $1/2$  [26]. Recently, methods have been found to improve this probability [27]. Whether similar methods could be used to improve the search on the star graph is a subject for further research.

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