Erratum: Minkowski Flux Vacua of Type II Supergravities **[Phys. Rev. Lett. 118, 011603 (2017)]**

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A few terms have been missed in an equation of [1]. The correct version leads to additional terms in Eq. (4) of our Letter: it becomes

$$\begin{aligned}
 e^{-2A}\tilde{\mathcal{R}}_4 = & -\left|e^{4A} *_\perp de^{-4A} - \varepsilon_p e^\phi F_{8-p}^{(0)}\right|^2 - (\text{flux})^2 \\
 & - \frac{1}{2} \left| *_\perp H|_\perp + \varepsilon_p e^\phi F_{6-p}|_\perp \right|^2 + e^{-2A} \partial(\dots) \\
 & - \frac{1}{2} \sum_{a_\parallel} \left| *_\perp (de^{a_\parallel})|_\perp - \varepsilon_p e^\phi (l_{a_\parallel} F_{8-p}^{(1)}) \right|^2 \\
 & - (\mathcal{R}_\parallel + \mathcal{R}_\parallel^\perp) + \frac{1}{2} |H^{(2)}|^2 + |H^{(3)}|^2.
 \end{aligned} \tag{1}$$

This expression now involves the H -flux components $H^{(2)}$ and $H^{(3)}$, having two and three indices along the internal flat directions parallel to the sources. Integrating Eq. (1) over the 6D compact manifold as in [1] leads to a 4D Ricci scalar $\tilde{\mathcal{R}}_4$ only given in terms of negative contributions, except for the last line of Eq. (1) of indefinite sign. We thus set this line to zero by imposing

$$H^{(2)} = H^{(3)} = 0, \tag{2}$$

on top of the geometric restrictions

$$\tilde{f}^{a_\parallel}_{b_\parallel c_\parallel} = \tilde{f}^{a_\parallel}_{b_\perp c_\parallel} = \tilde{f}^{a_\perp}_{b_\parallel c_\parallel} = \tilde{f}^{a_\perp}_{b_\perp c_\parallel} = 0, \tag{3}$$

$$\tilde{f}^{a_\perp}_{b_\perp c_\perp} = 0, \tag{4}$$

that imply $\mathcal{R}_\parallel = \mathcal{R}_\parallel^\perp = 0$.

The only restrictions made on the Minkowski solutions are Eqs. (2)–(4) above. The new one, (2), does not change anything in our result: the same class of Minkowski solutions is obtained, with the same method. Interestingly, all known supersymmetric Minkowski flux solutions on twisted tori with parallel sources verify the new requirement in Eq. (2), so our final discussion is not modified either.

It is worth noting that the last line of Eq. (1) is better motivated than the curvature terms alone: indeed, on nilmanifolds with constant H -flux at least, this field combination is T -duality invariant [1]. Similarly to Sec. IV of our Letter, this line may then admit an extension that includes nongeometric fluxes.

[1] D. Andriot and J. Blåbäck, Refining the boundaries of the classical de Sitter landscape, *J. High Energy Phys.* **03** (2017) 102.