Erratum: Large- N_c limit reduces the number of independent few-body parity-violating low-energy constants in pionless effective field theory [Phys. Rev. C **93, 025502 (2016)]**

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(Received 20 March 2018; published 8 May 2018)

DOI: [10.1103/PhysRevC.97.059901](https://doi.org/10.1103/PhysRevC.97.059901)

There are instances of a missing $\sin^2 \theta_W$ where we discuss the isotensor terms. Reference [1] is also missing this factor. The corrections are as follows: Eq. (9),

$$\mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) \mathcal{I}_{ab} \tau_1^a \tau_2^b$$

scales as $O(N_c) \sin^2 \theta_W$, while Eqs. (18) and (19),

$$\mathbf{p}_+ \cdot (\vec{\sigma}_1 - \vec{\sigma}_2) \mathcal{I}_{ab} \tau_1^a \tau_2^b,$$

$$\mathbf{p}_+^2 \mathbf{p}_- \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2) \mathcal{I}_{ab} \tau_1^a \tau_2^b$$

contribute at $O(N_c^{-1}) \sin^2 \theta_W$. In both Eqs. (24) and (29) the $\tilde{\mathcal{G}}_5$ term scales as

$$\tilde{\mathcal{G}}_5 \sim N_c \sin^2 \theta_W,$$

while in Eq. (27) the terms $\mathcal{A}_5^\pm$ scale as follows:

$$\mathcal{A}_5^+ \sim N_c^{-1} \sin^2 \theta_W, \quad \mathcal{A}_5^- \sim N_c \sin^2 \theta_W.$$

Finally, the $C_{(\Delta I=2)}^{(1S_0-3P_0)}$ term in Eq. (33) scales as

$$C_{(\Delta I=2)}^{(1S_0-3P_0)} \sim N_c \sin^2 \theta_W.$$

Our conclusions about the impact of large- N_c scaling on observables remain unchanged. While factors of $\sin^2 \theta_W$ appear at the weak (matching) scale, the low-energy constants evolve in a nonperturbative (and currently undetermined) way down into hadronic scales. It is always the case that unknown multipliers can modify the naive ordering of large- N_c terms. The physics case for determining $C_{(\Delta I=2)}^{(1S_0-3P_0)}$ from experiment and from lattice QCD remains strong.

[1] D. R. Phillips, D. Samart, and C. Schat, *Phys. Rev. Lett.* **114**, 062301 (2015).